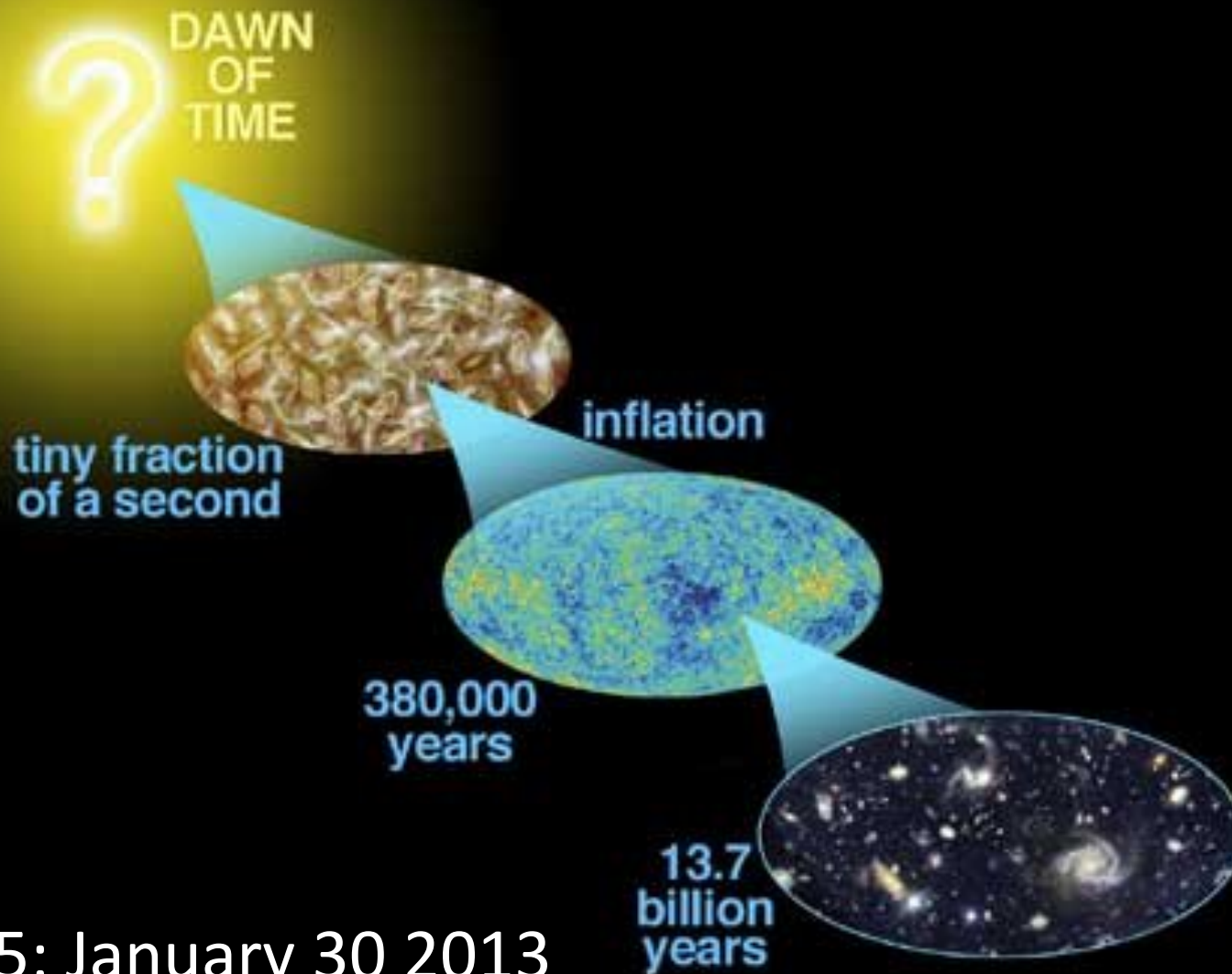


Cosmology W13



Lecture 5: January 30 2013

Strong Lensing

- Introduction to strong gravitational lensing
- strong lensing as a tool for cosmology:
cosmography from gravitational time delays

References

- General
 - Schneider, Kochanek and Wambsganss 2006
 - Bartelmann 2010, CQGra, 27, 23
- Strong lensing by galaxies
 - Treu ARA&A, 2010, 87, 48
- Lensing by clusters
 - Kneib & Natarajan 2011, AARev, 19, 47
- Cosmology from gravitational time delays
 - Suyu et al. 2013, ApJ, arXiv 1208.6010

A very useful reference

- <http://www.astro.uni-bonn.de/~peter/SaaSfee.html>

33rd Advanced Saas-Fee Course of the Swiss Society for Astrophysics and Astronomy

**GRAVITATIONAL LENSING:
STRONG, WEAK, AND MICRO**

**April 7-12, 2003
Les Diablerets, Switzerland**

SPEAKERS:
Strong Lensing, C. S. Kochanek (CIA)
Weak Lensing, P. Schneider (Bonn)
Micro Lensing, J. Wambsgans (Potsdam)

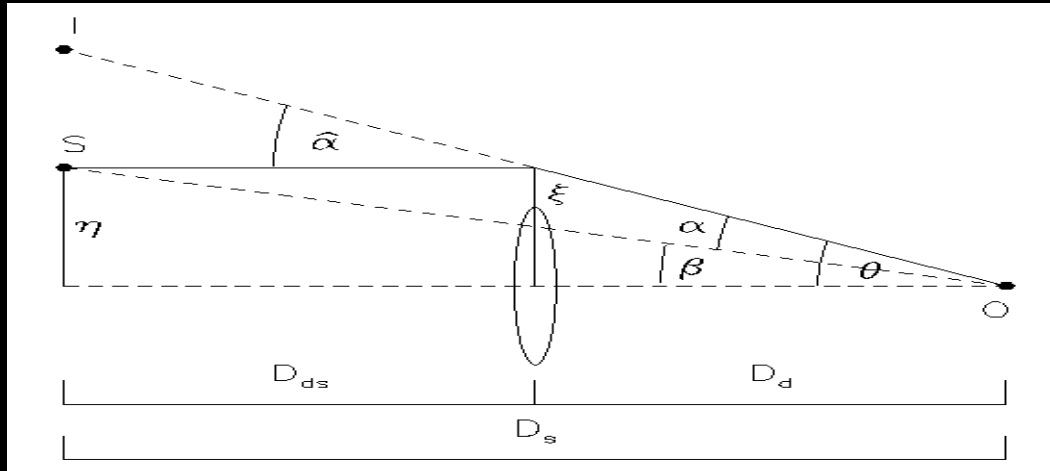
CONTACT:
web, <http://obswww.unige.ch/saas-fee>
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or gmeylan@stsci.edu

ORGANIZERS:
Georges Meylan (STSO)
Philippe Jetzer (Geneve)
Pierre North (Lausanne)

IMAGE CREDITS:
Left, Kurt Müller and <http://photo.zemmat.ch>
Right, R. McLeod (CIA, Cardiff) and F. Summers (STSO)

You are expected to study and know the content of the introductory chapter by Schneider

How does it work? Fermat's principle



Fermat distance

Shapiro delay

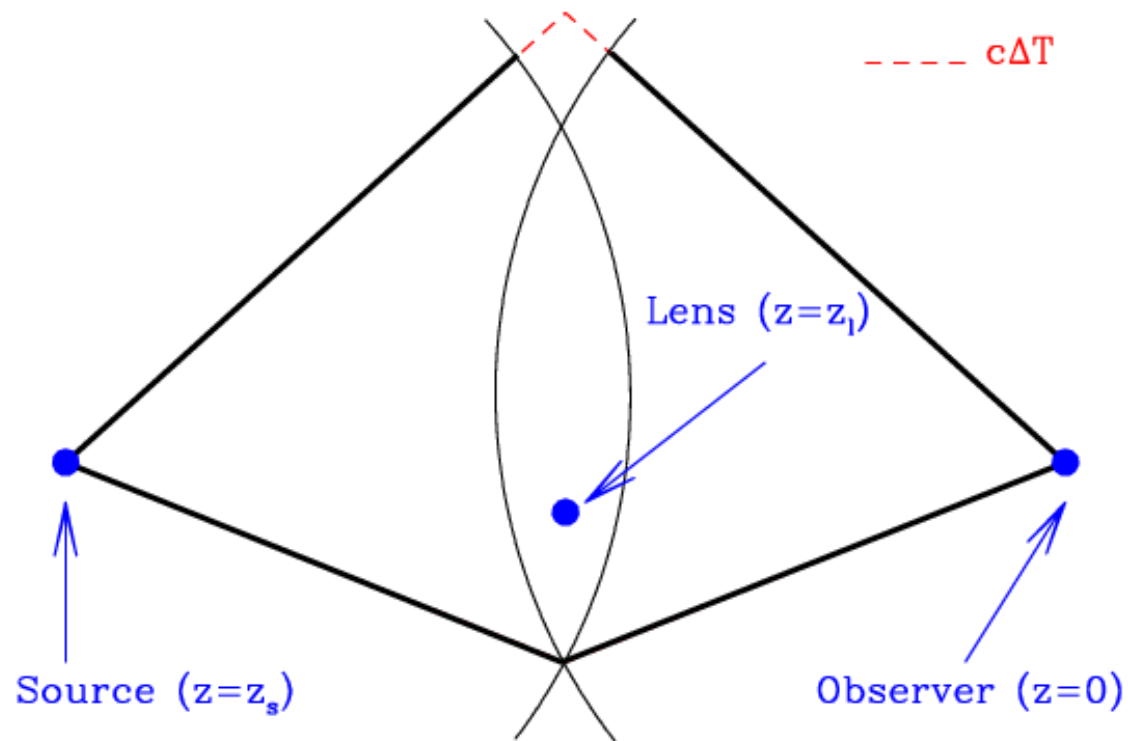
$$t(\vec{\theta}) = \frac{(1 + z_d)}{c} \frac{D_d D_s}{D_{ds}} \left[\frac{1}{2} (\vec{\theta} - \vec{\beta})^2 - \psi(\vec{\theta}) \right]$$

Excess time delay

geometric time delay

How does it work?

$$t(\vec{\theta}) = \frac{(1+z_d)}{c} \frac{D_d D_s}{D_{ds}} \left[\frac{1}{2} (\vec{\theta} - \vec{\beta})^2 - \psi(\vec{\theta}) \right]$$



Recap of useful formulae - notation

$$\Sigma(\vec{\xi}) = \int \rho(\vec{\xi}, z) dz$$

$$\vec{\hat{\alpha}}(\vec{\xi}) = \frac{4G}{c^2} \int \frac{(\vec{\xi} - \vec{\xi}') \Sigma(\vec{\xi}')}{|\vec{\xi} - \vec{\xi}'|^2} d^2 \xi'$$

$$\hat{\alpha}(\xi) = \frac{4GM(\xi)}{c^2 \xi}$$

$$M(\xi) = 2\pi \int_0^\xi \Sigma(\xi') \xi' d\xi'$$

$$\vec{\alpha} = \frac{D_{\text{ds}}}{D_s} \vec{\hat{\alpha}}$$

$$\vec{\beta} = \vec{\theta} - \vec{\alpha}(\vec{\theta})$$

$$\Sigma_{\text{cr}} = \frac{c^2}{4\pi G} \frac{D_s}{D_d D_{\text{ds}}} = 0.35 \text{ g cm}^{-2} \left(\frac{D}{1 \text{ Gpc}} \right)^{-1}$$

Recap of useful formulae - notation

$$\psi(\vec{\theta}) = \frac{D_{\text{ds}}}{D_{\text{d}}D_{\text{s}}} \frac{2}{c^2} \int \Phi(D_{\text{d}}\vec{\theta}, z) dz$$

$$\vec{\nabla}_{\theta}\psi = D_{\text{d}}\vec{\nabla}_{\xi}\psi = \frac{2}{c^2} \frac{D_{\text{ds}}}{D_{\text{s}}} \int \vec{\nabla}_{\perp}\Phi dz = \vec{\alpha}$$

$$\nabla_{\theta}^2\psi = \frac{2}{c^2} \frac{D_{\text{d}}D_{\text{ds}}}{D_{\text{s}}} \int \nabla_{\xi}^2\Phi dz = \frac{2}{c^2} \frac{D_{\text{d}}D_{\text{ds}}}{D_{\text{s}}} \cdot 4\pi G \Sigma = 2 \frac{\Sigma(\vec{\theta})}{\Sigma_{\text{cr}}} \equiv 2\kappa(\vec{\theta})$$

$$\vec{\alpha}(\vec{\theta}) = \vec{\nabla}\psi = \frac{1}{\pi} \int \kappa(\vec{\theta}') \frac{\vec{\theta} - \vec{\theta}'}{|\vec{\theta} - \vec{\theta}'|^2} d^2\theta'$$

$$\mathcal{A} \equiv \frac{\partial \vec{\beta}}{\partial \vec{\theta}} = \left(\delta_{ij} - \frac{\partial \alpha_i(\vec{\theta})}{\partial \theta_j} \right) = \left(\delta_{ij} - \frac{\partial^2 \psi(\vec{\theta})}{\partial \theta_i \partial \theta_j} \right) = \mathcal{M}^{-1}$$

$$\mu = \det M = \frac{1}{\det \mathcal{A}} = \frac{1}{(1 - \kappa)^2 - |\gamma|^2}$$

What can strong lensing do for us?

- Strong lensing measures very accurately:
 - Mass enclosed within Einstein Radius
 - Derivative of enclosed mass (sometimes)
 - Projected orientation of the mass
 - Projected ellipticity of the mass
- Strengths:
 - Mass independent of presence of tracers
 - Mass independent of dynamic state
 - Measures total mass
- Weaknesses:
 - Projection effects
 - Measures total mass

Example: SIS as a simple model for galaxy scale lenses

$$\rho(r) = \frac{\sigma_v^2}{2\pi G r^2}$$

$$\Sigma(\xi) = \int_{-\infty}^{\infty} dr_3 \rho\left(\sqrt{\xi^2 + r_3^2}\right) = \frac{\sigma_v^2}{2G} \xi^{-1}$$

$$\theta_E = 4\pi \left(\frac{\sigma_v}{c}\right)^2 \frac{D_{ds}}{D_s}$$

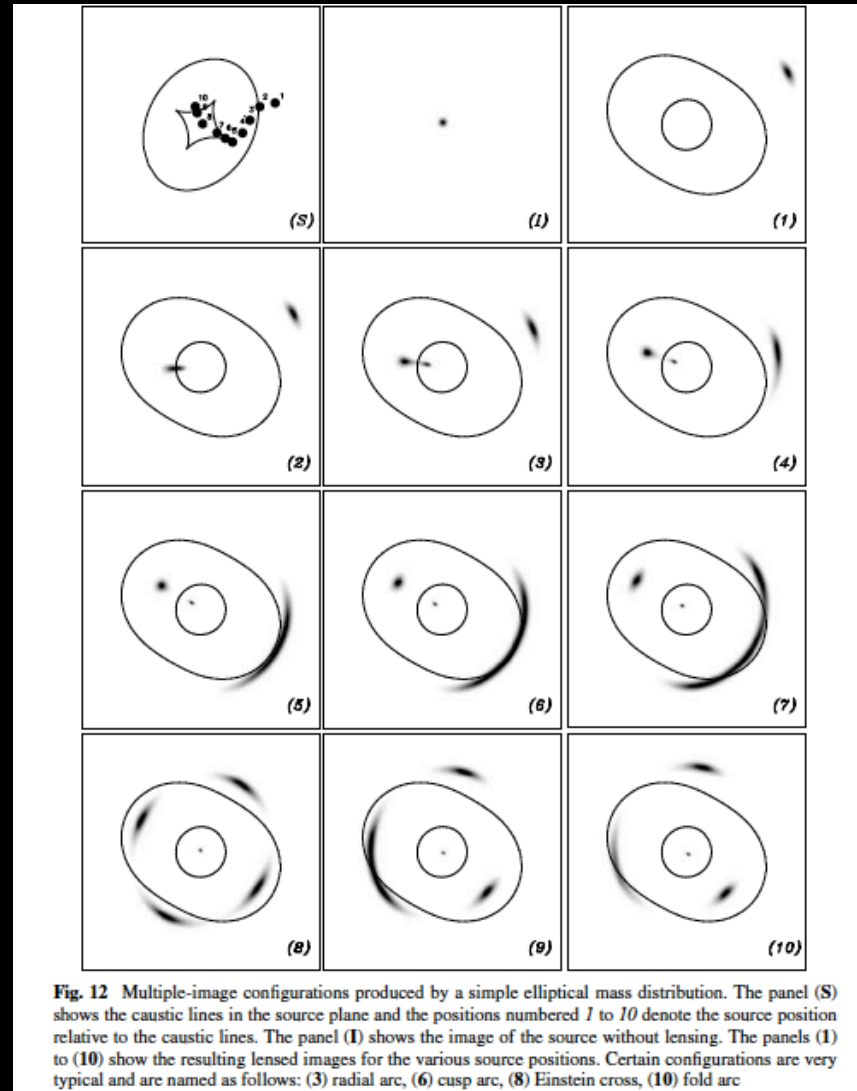
$$\kappa(\theta) = \frac{\theta_E}{2|\theta|}; \quad \bar{\kappa}(\theta) = \frac{\theta_E}{|\theta|}; \quad |\gamma|(\theta) = \frac{\theta_E}{2|\theta|}; \quad \alpha(\theta) = \theta_E \frac{\theta}{|\theta|}$$

$$x = \theta/\theta_E, \quad y = \beta/\theta_E$$

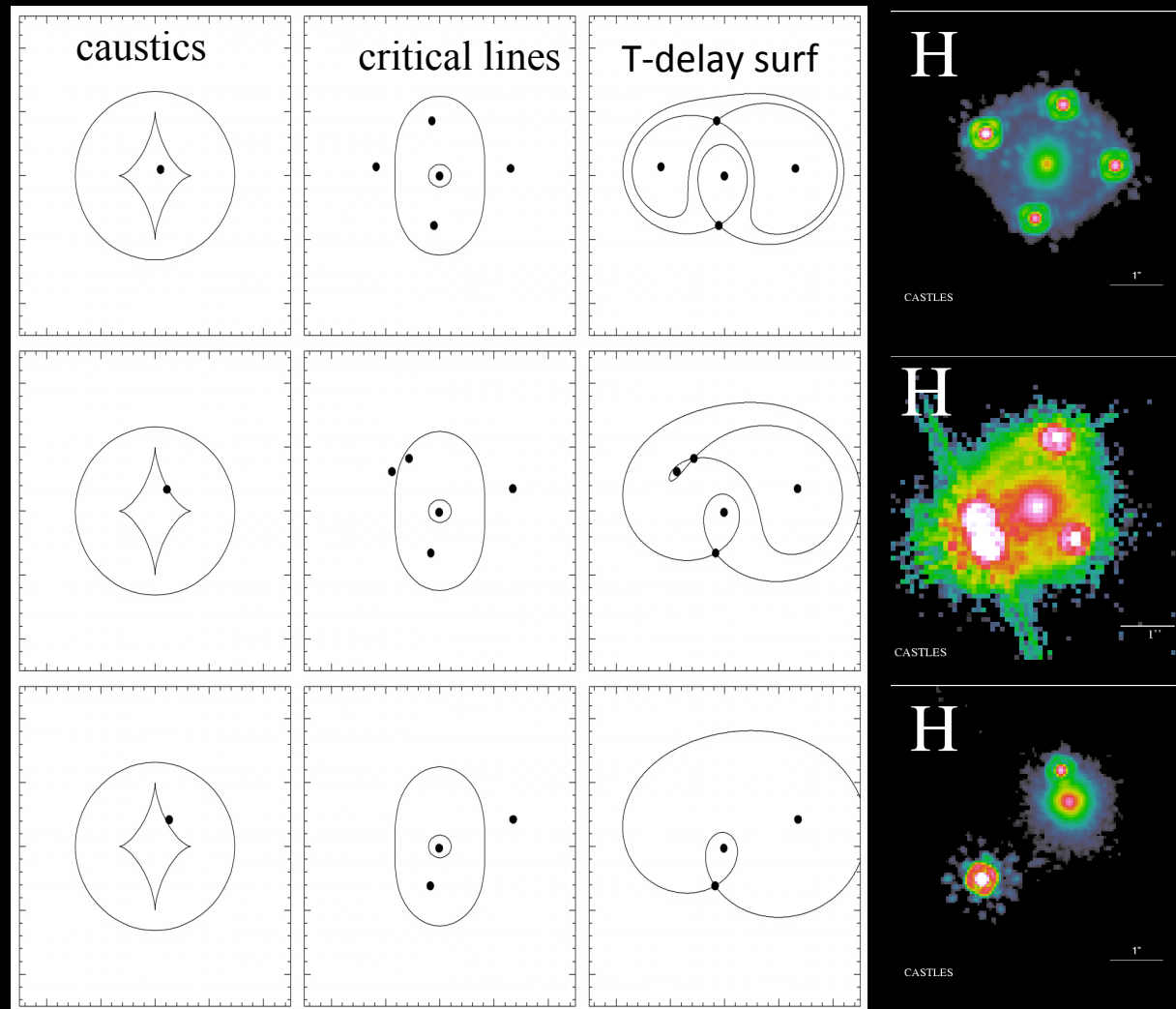
$$\beta = \theta - \theta_E \frac{\theta}{|\theta|}, \quad \text{or} \quad y = x - \frac{x}{|x|}$$

In practice you need elliptical models at least and those cannot be treated analytically

For simple elliptical mass distributions



Examples of galaxy-scale lenses



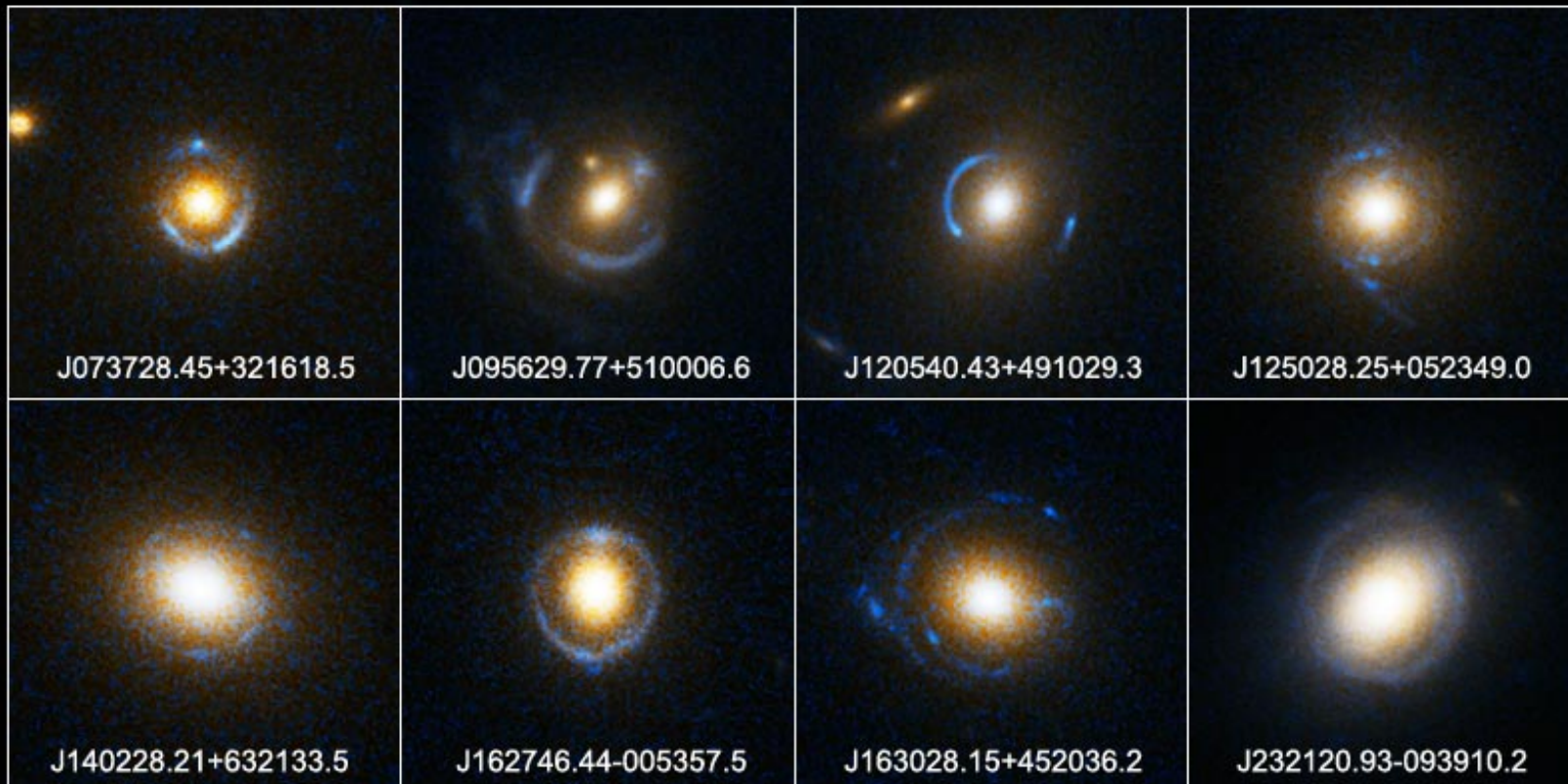
Saha & Williams 2003

Examples of galaxy-scale lenses

Classify the configurations

Einstein Ring Gravitational Lenses

Hubble Space Telescope • ACS



NASA, ESA, A. Bolton (Harvard-Smithsonian CfA), and the SLACS Team

STScI-PRC05-32

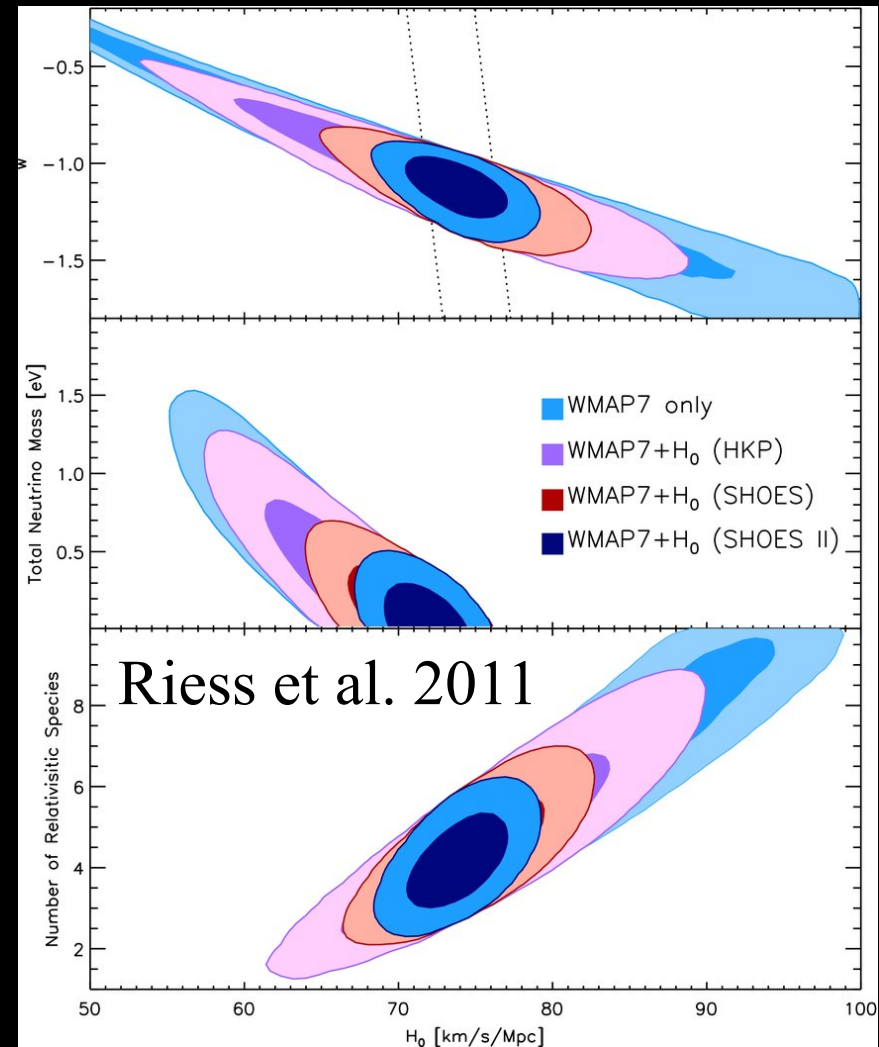
Cosmography with time-delays

Time delay distance in practice

$$\Delta t \propto D_{\Delta t}(z_s, z_d) \propto H_0^{-1} f(\Omega_m, w, \dots)$$

Steps:

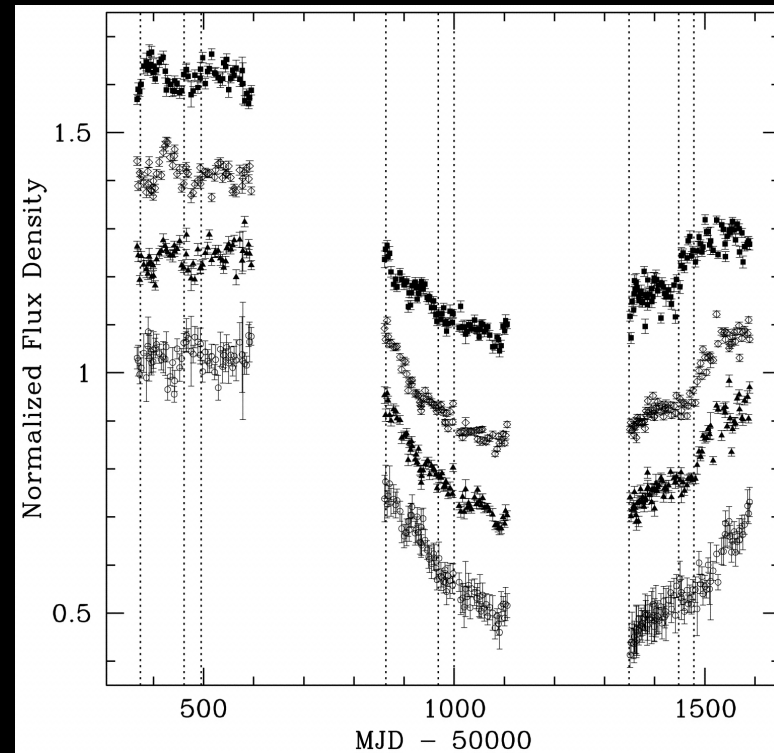
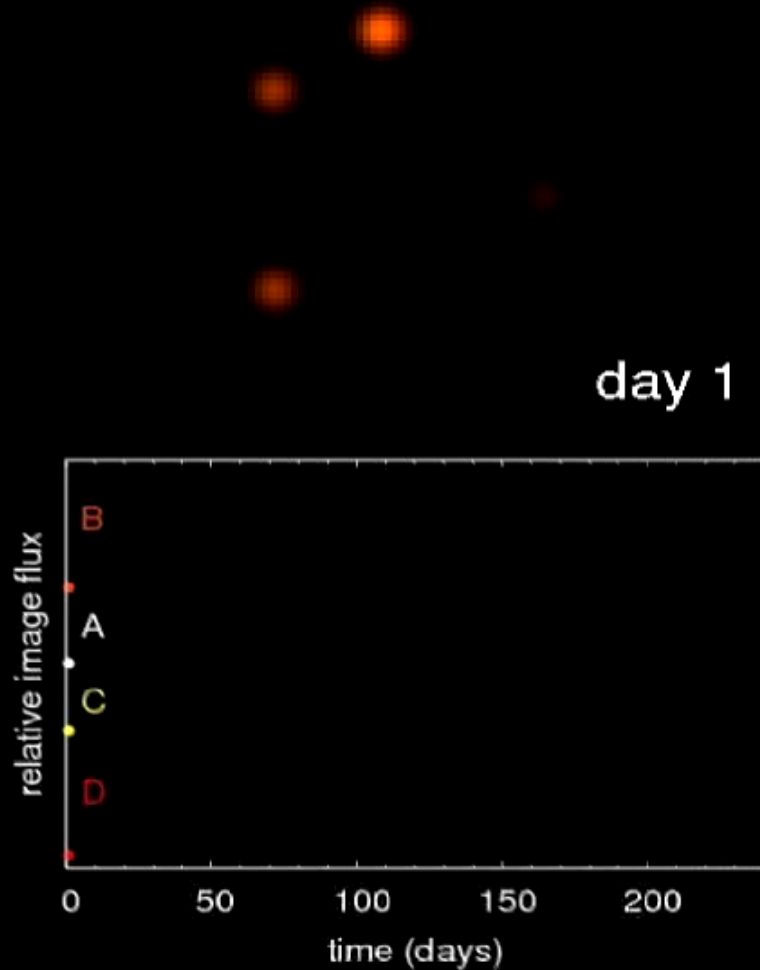
- Measure the time-delay between two images
- Measure and model the potential
- Infer the time-delay distance
- Convert it into cosmological parameters



Cosmography with strong lenses: the 4 problems solved

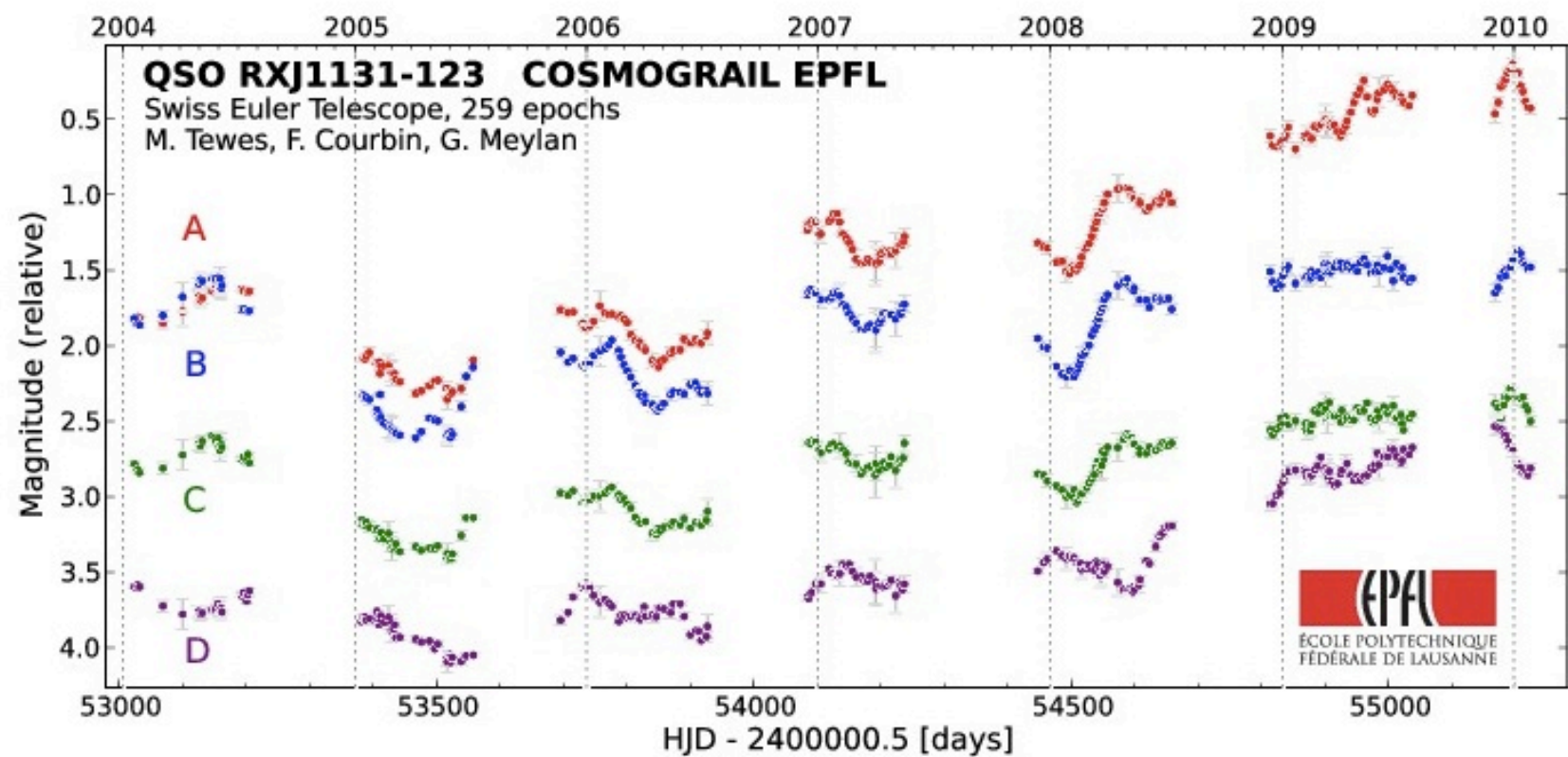
- Time delay – 2-3 %
 - Tenacious monitoring (e.g. Fassnacht et al. 2002); COSMOGRAIL (Meylan/Courbin)
- Astrometry – 10-20 mas
 - Hubble/VLA/(Adaptive Optics?)
- Lens potential (2-3%)
 - Stellar kinematics/Extended sources (Treu & Koopmans 2002; Suyu et al. 2009)
- Structure along the line of sight (2-3%) and the mass-sheet degeneracy,
 - Galaxy counts and numerical simulations (Suyu et al. 2009)
 - Stellar kinematics (Koopmans et al. 2003)

Gravitational Lens Time Delays



B1608+656 Variability in Radio Observations

Credits: S. H. Suyu, C. D. Fassnacht, NRAO/AUI/NSF

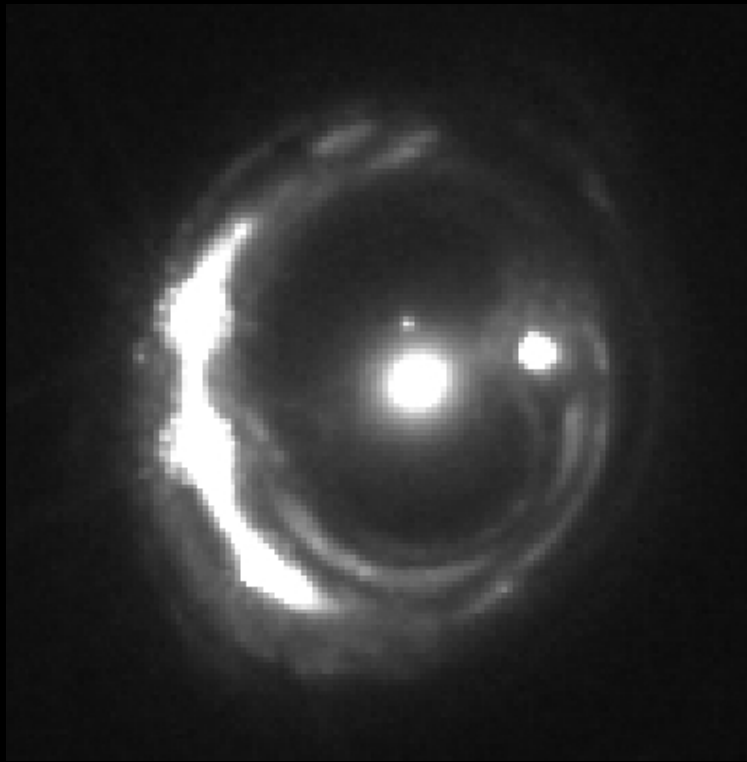


The slope- H_0 degeneracy

- Uncertainty of the mass density profile of the main lens is the single largest source of uncertainty.
- If mass density profile is a power law $\rho \sim r^{-\gamma'}$ then $H_0(\gamma') \sim H_0(2)(\gamma' - 1)$ (Wucknitz 2002)
- Can be broken using lensing (Suyu et al. 2010) and stellar kinematic data (Treu & Koopmans 2002)

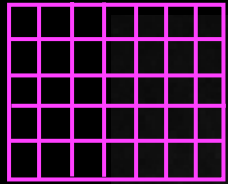
Lens Model

Observed Image

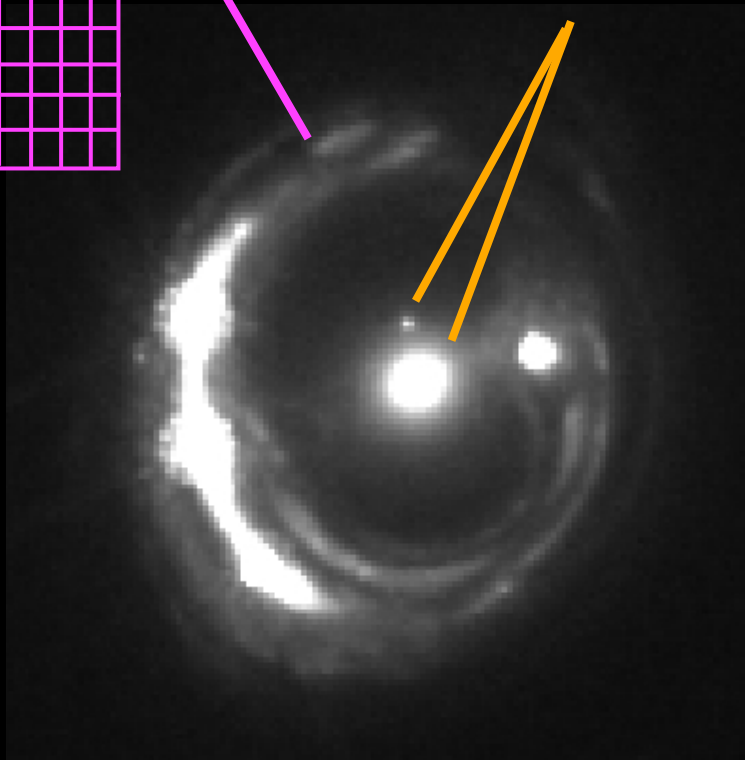


Lens Model

light distribution
of extended source



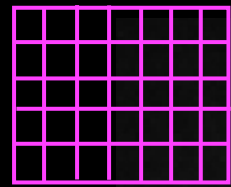
mass distribution
of lens



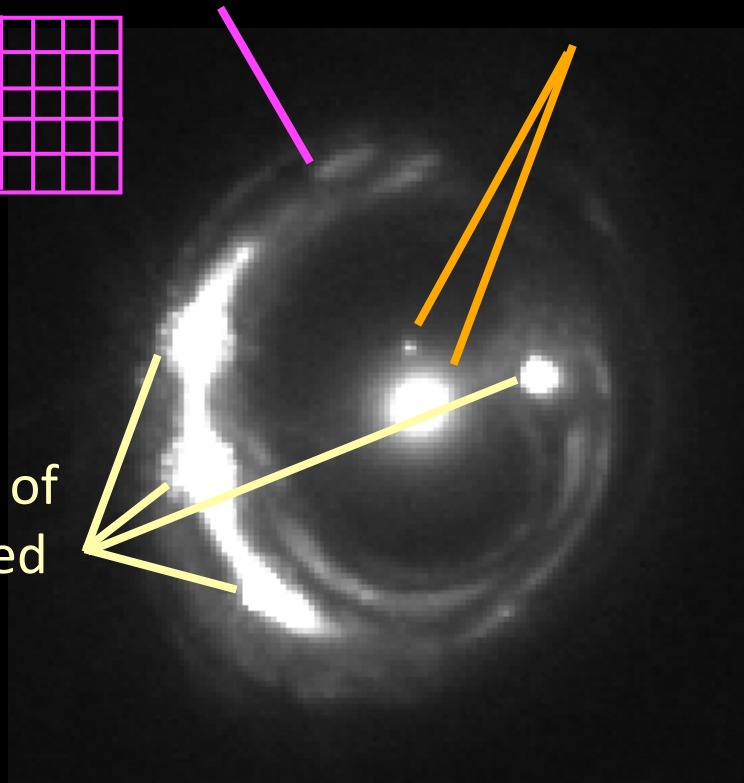
Lens Model

light distribution
of extended source

mass distribution
of lens

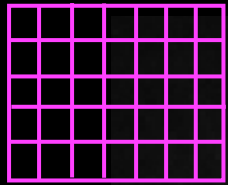


light of
lensed
AGN



Lens Model

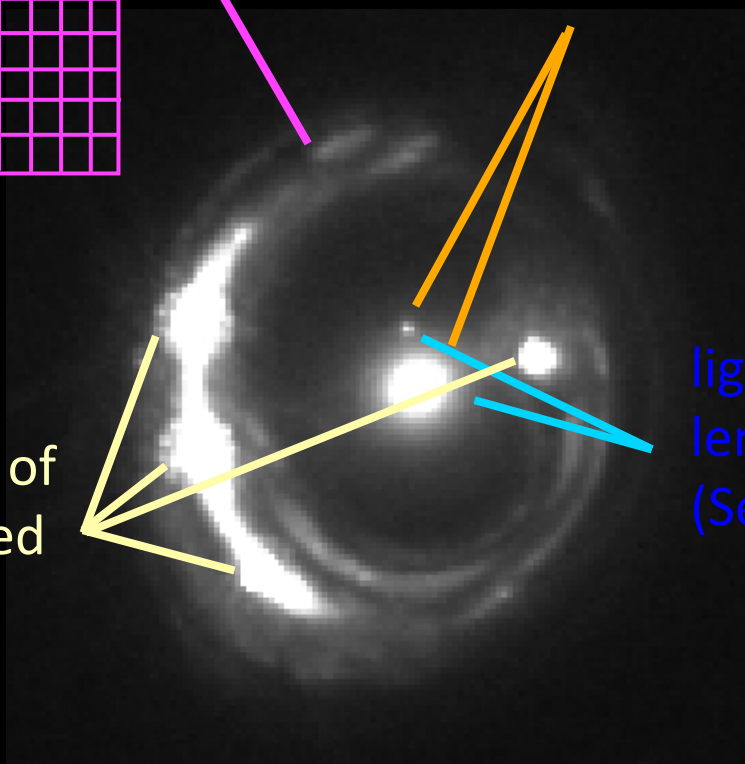
light distribution
of extended source



mass distribution
of lens

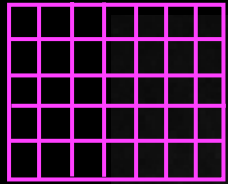
light of
lensed
AGN

light of
lens
(Sersic)

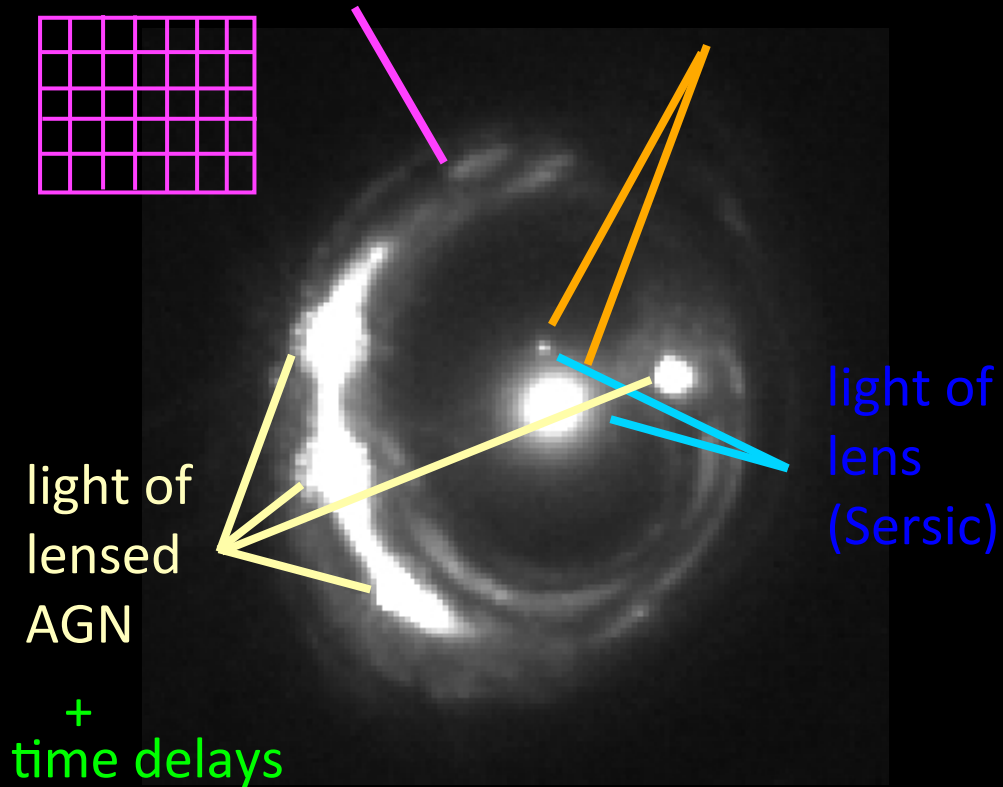


Lens Model

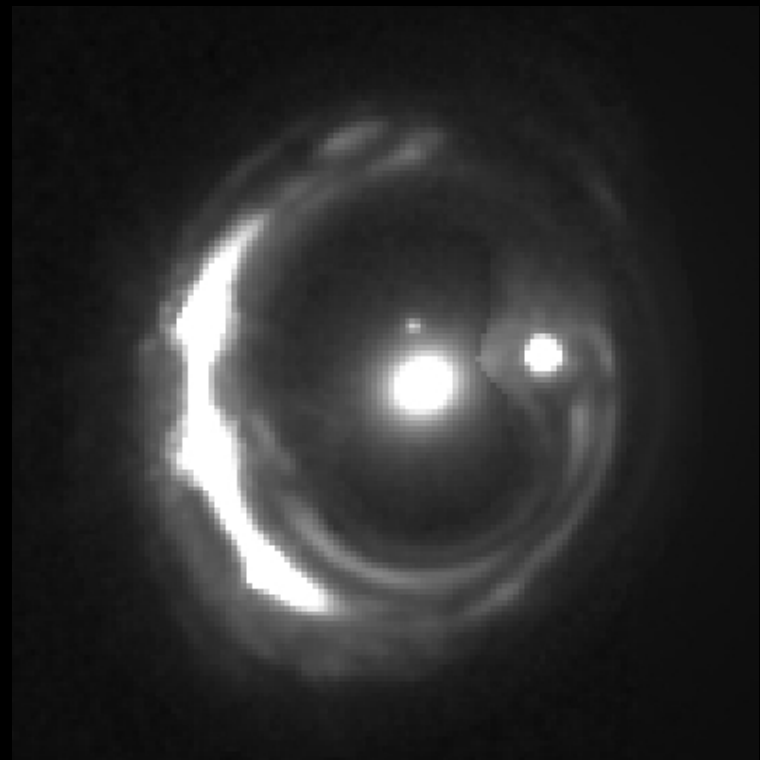
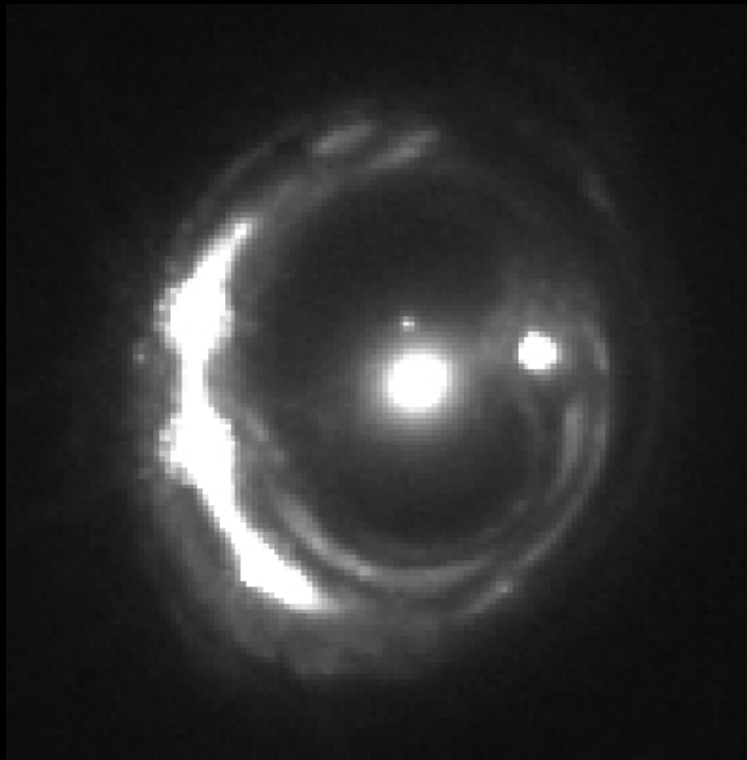
light distribution
of extended source



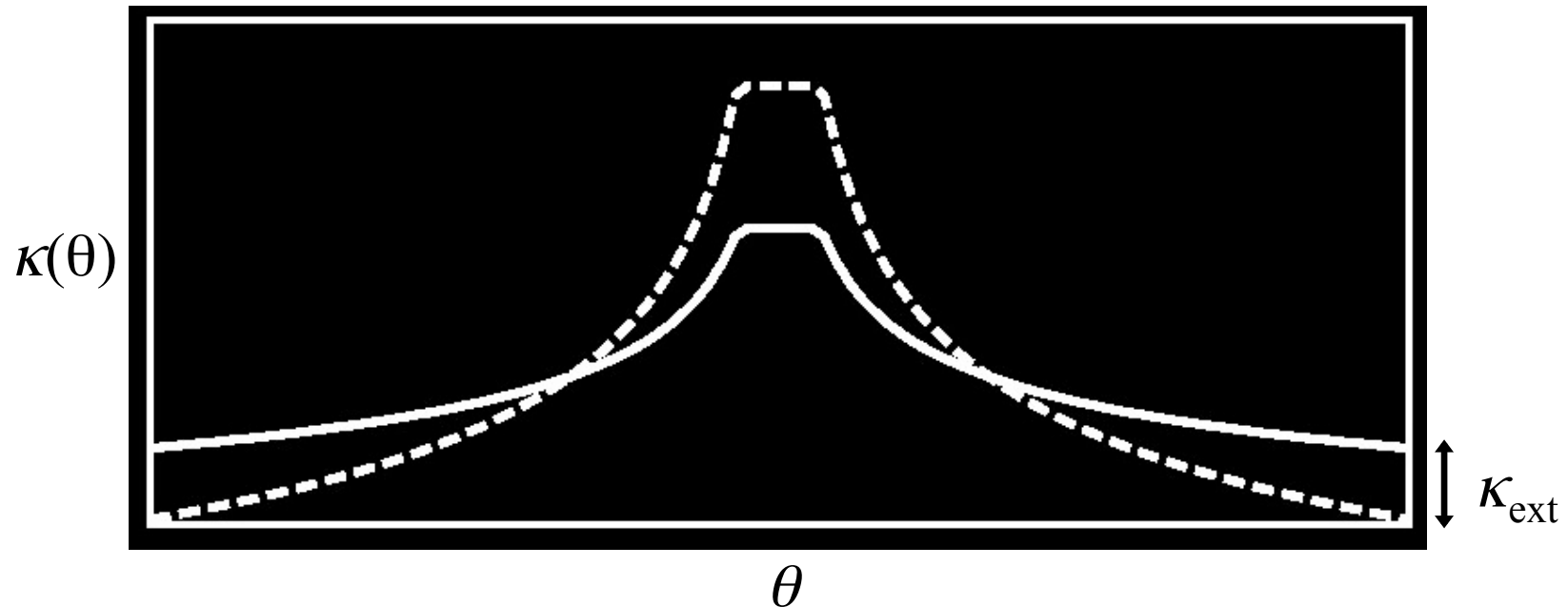
mass distribution
of lens



Lens Model



Mass-sheet degeneracy



lensing observables
do not change, but

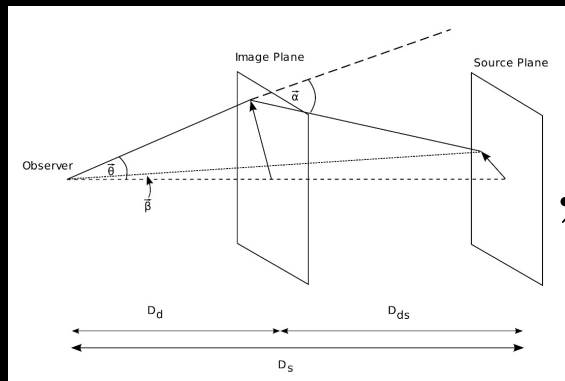
$$D_{\Delta t}^{\text{true}} = \frac{D_{\Delta t}^{\text{model}}}{1 - \kappa_{\text{ext}}}$$

Mass-sheet degeneracy: why?

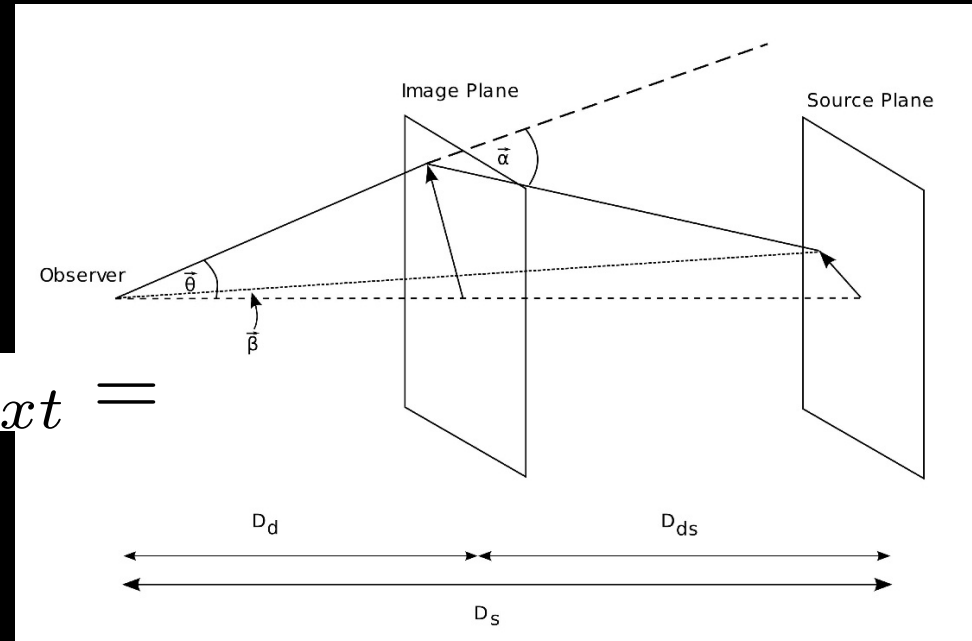
$$\kappa_\lambda = (1 - \lambda) + \lambda\kappa; \quad \beta_\lambda = \beta/\lambda$$

$$\Delta t_\lambda = \lambda\Delta t; \quad \mu_\lambda = \mu_\lambda^2$$

$$\lambda = 1 - \kappa_{ext}$$

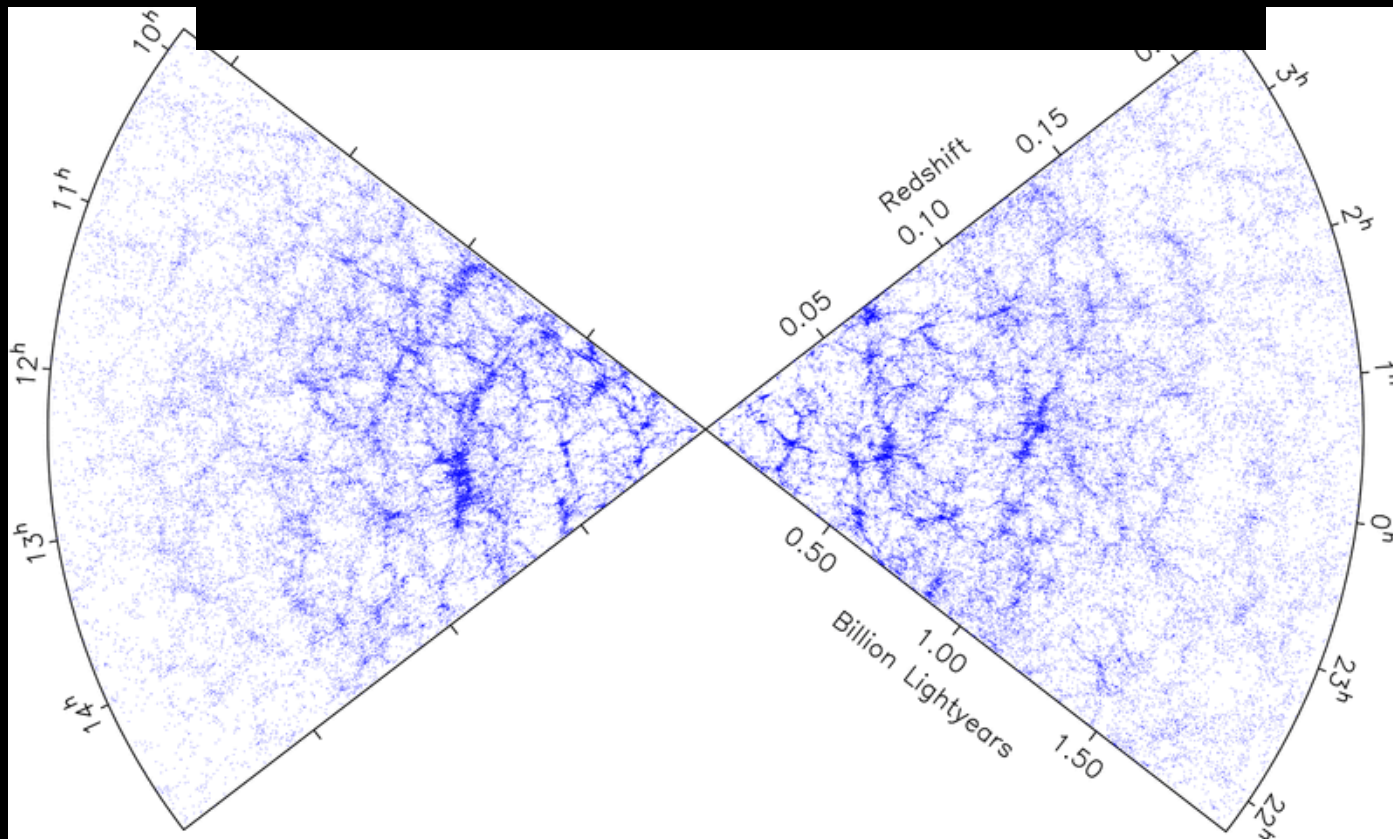


” \times ” $\kappa_{ext} =$



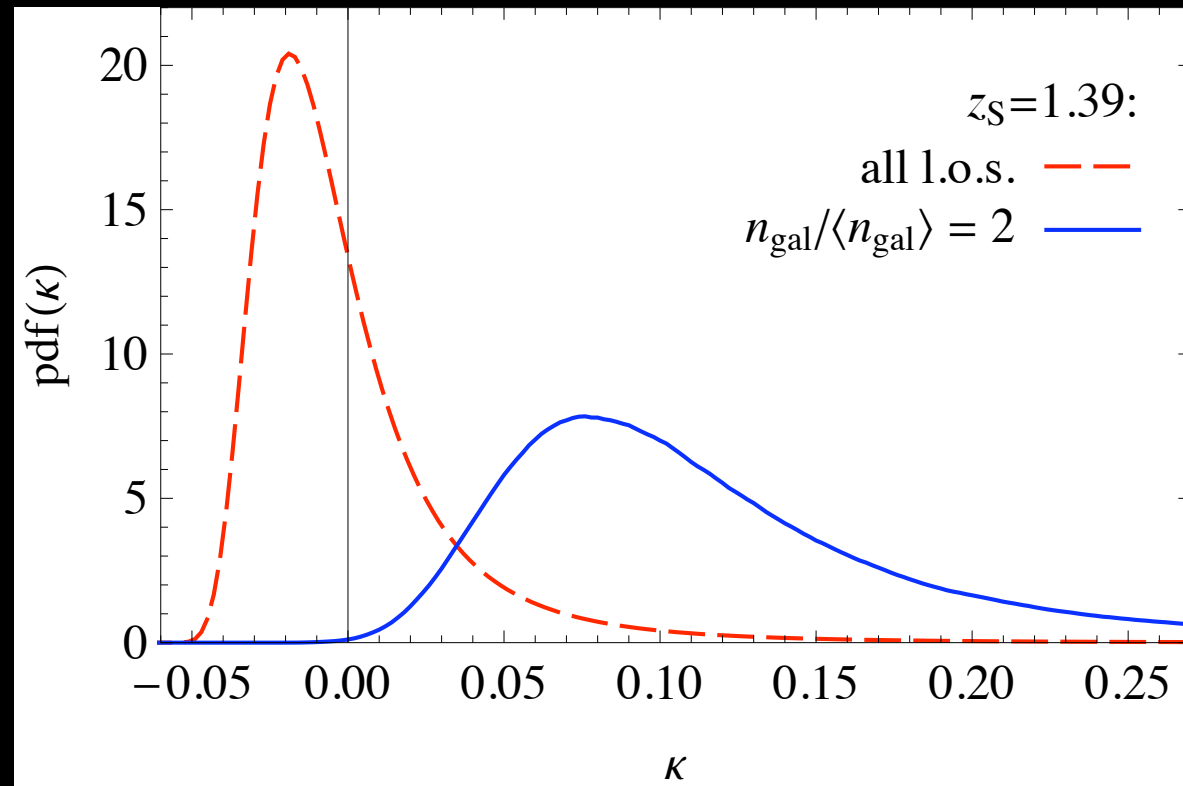
Line of sight vs mass sheet

$$t(\vec{\theta}) = \frac{(1+z_d) D_d D_s}{c D_{ds}} \left[\frac{1}{2} (\vec{\theta} - \vec{\beta})^2 - \psi(\vec{\theta}) \right]$$



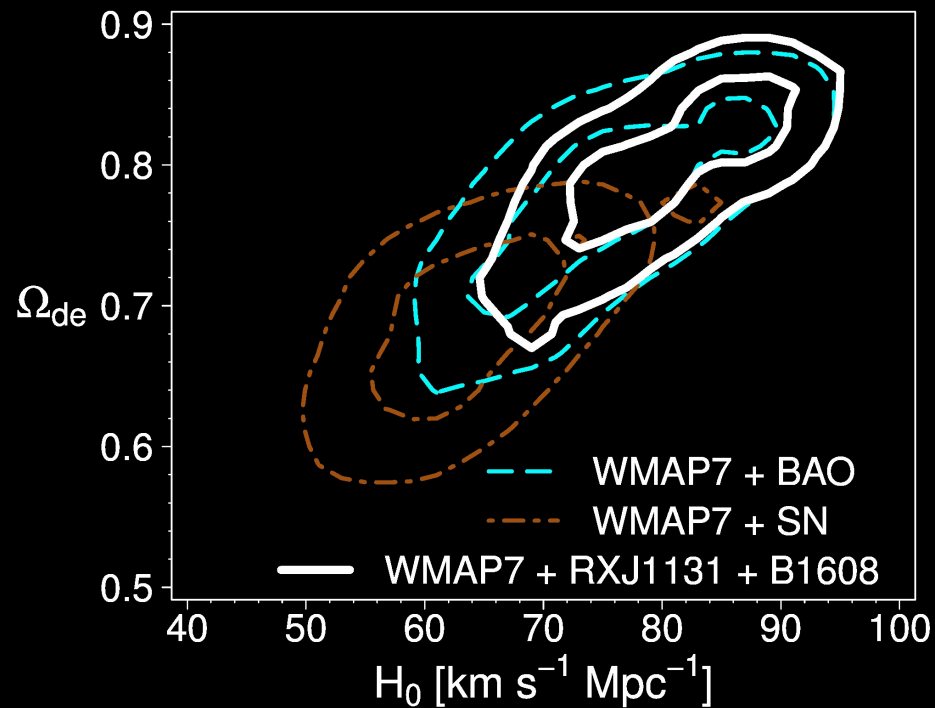
Structure along the line of sight

- $H_0(\kappa_{\text{ext}}) = H_0(0) * (1 - \kappa_{\text{ext}})$
- Ray-trace through simulations
- Choose light cones comparable to those of the lens
- Observables:
 - Number counts
 - Shear
- Transform line-of-sight structure into κ_{ext} !!

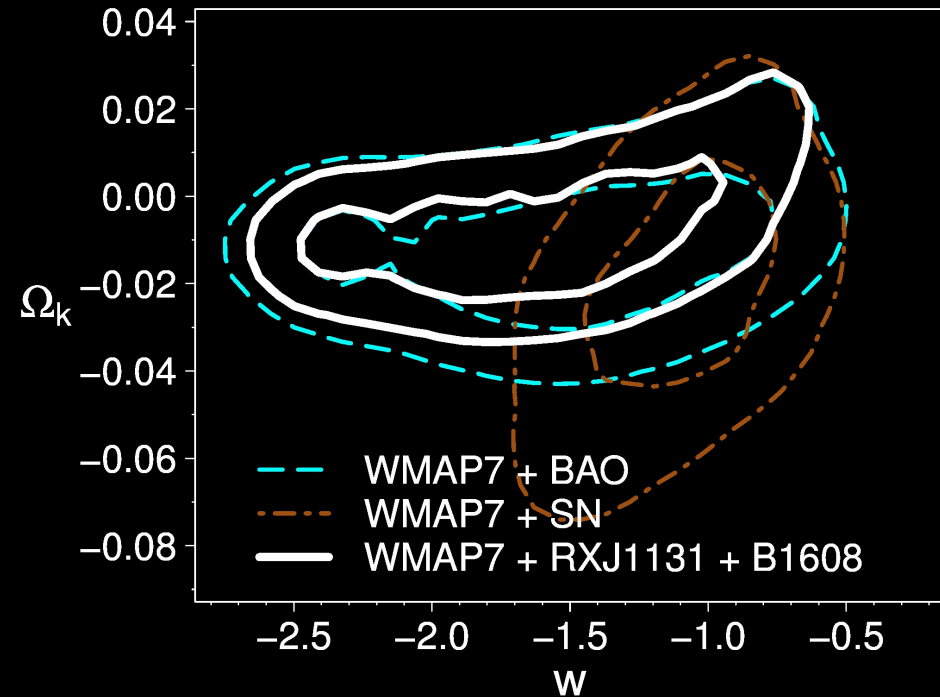


Time delays vs other probes

WMAP7 Λ CDM prior



[Suyu et al. 2012]



- contour orientations are different: complementarity b/w probes
- contour sizes are similar: lensing is a competitive probe

The end