A FAST MEAN-REVERTING CORRECTION TO HESTON
STOCHASTIC VOLATILITY MODEL

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Abstract. We propose a multi-scale stochastic volatility model in which a fast mean-reverting factor of volatility is built on top of the Heston stochastic volatility model. A singular perturbative expansion is then used to obtain an approximation for European option prices. The resulting pricing formulas are semi-analytic, in the sense that they can be expressed as integrals. Difficulties associated with the numerical evaluation of these integrals are discussed, and techniques for avoiding these difficulties are provided. Overall, it is shown that computational complexity for our model is comparable to the case of a pure Heston model, but our correction brings significant flexibility in terms of fitting to the implied volatility surface. This is illustrated numerically and with option data.

Key words. Stochastic volatility, Heston model, fast mean-reversion, asymptotics, implied volatility smile/skew.

AMS subject classifications. 60F99, 91B70

1. Introduction. Since its publication in 1993, the Heston model [11] has received considerable attention from academics and practitioners alike. The Heston model belongs to a class of models known as stochastic volatility models. Such models relax the assumption of constant volatility in the stock price process, and instead, allow volatility to evolve stochastically through time. As a result, stochastic volatility models are able to capture some of the well-known features of the implied volatility surface, such as the volatility smile and skew (slope at the money). Among stochastic volatility models, the Heston model enjoys wide popularity because it provides an explicit, easy-to-compute, integral formula for calculating European option prices. In terms of the computational resources needed to calibrate a model to market data, the existence of such a formula makes the Heston model extremely efficient compared to models that rely on Monte Carlo techniques for computation and calibration.

Yet, despite its success, the Heston model has a number of documented shortcomings. For example, it has been statistically verified that the model misprices far in-the-money and out-of-the-money European options [6] [20]. In addition, the model is unable to simultaneously fit implied volatility levels across the full spectrum of option expirations available on the market [9]. In particular, the Heston model has difficulty fitting implied volatility levels for options with short expirations [10]. In fact, such problems are not limited to the Heston model. Any stochastic volatility model in which the volatility is modeled as a one-factor diffusion (as is the case in the Heston model) has trouble fitting implied volatility levels across all strikes and maturities [10].

One possible explanation for why such models are unable to fit the implied volatility surface is that a single factor of volatility, running on a single time scale, is simply not sufficient for describing the dynamics of the volatility process. Indeed, the existence of several stochastic volatility factors running on different time scales has been well-documented in literature that uses empirical return data [1] [2] [3] [5] [8] [12] [15]
Multi-scale stochastic volatility models represent a struggle between two opposing forces. On one hand, adding a second factor of volatility can greatly improve a model’s fit to the implied volatility surface of the market. On the other hand, adding a second factor of volatility often results in the loss of some, if not all, analytic tractability. Thus, in developing a multi-scale stochastic volatility model, one seeks to model market dynamics as accurately as possible, while at the same time retaining a certain level of analyticity. Because the Heston model provides explicit integral formulas for calculating European option prices, it is an ideal template on which to build a multi-scale model and accomplish this delicate balancing act.

In this paper, we show one way to bring the Heston model into the realm of multi-scale stochastic volatility models without sacrificing analytic tractability. Specifically, we add a fast mean-reverting component of volatility on top of the Cox–Ingersoll–Ross (CIR) process that drives the volatility in the Heston model. Using the multi-scale model, we perform a singular perturbation expansion, as outlined in [7], in order to obtain a correction to the Heston price of a European option. This correction is easy to implement, as it has an integral representation that is quite similar to that of the European option pricing formula produced by the Heston model.

The paper is organized as follows. In Section 2 we introduce the multi-scale stochastic volatility model and we derive the resulting pricing partial differential equation (PDE) and boundary value for the European option pricing problem. In Section 4 we use a singular perturbative expansion to derive a PDE for a correction to the Heston price of a European option and in Section 6 we obtain a solution for this PDE. A proof of the accuracy of the pricing approximation is provide in Section 5. In Section 7 we examine how the implied volatility surface, as obtained from the multi-scale model, compares with that of the Heston model, and in Section 8 we present an example of calibration to market data. In Appendix A we review the dynamics of the Heston Stochastic volatility model under the risk-neutral measure, and present the pricing formula for European options. An explicit formula for the correction is given in Appendix B and a key moment estimate is derived in Appendix C. The issues associated with numerically evaluating the integrals-representations of option prices obtained from the multi-scale model are explored in Appendix D.
Here, $W^x_t$, $W^y_t$ and $W^z_t$ are one-dimensional standard Brownian motions with the correlation structure

$$
d \langle W^x, W^y \rangle_t = \rho_{xy} dt, \quad (2.5)$$

$$
d \langle W^x, W^z \rangle_t = \rho_{xz} dt, \quad (2.6)$$

$$
d \langle W^y, W^z \rangle_t = \rho_{yz} dt, \quad (2.7)$$

where the correlation coefficients $\rho_{xy}$, $\rho_{xz}$ and $\rho_{yz}$ are constants satisfying $\rho_{xy}^2 < 1$, $\rho_{xz}^2 < 1$, $\rho_{yz}^2 < 1$, and $\rho_{xy}^2 + \rho_{xz}^2 + \rho_{yz}^2 - 2\rho_{xy}\rho_{xz}\rho_{yz} < 1$.

As it should be, in (2.1) the stock price discounted by the risk-free rate $r$ is a martingale under the pricing risk neutral measure. The volatility $\Sigma_t$ is driven by two processes $Y_t$ and $Z_t$, through the product $\sqrt{Z_t} f(Y_t)$. The process $Z_t$ is a Cox–Ingersoll–Ross (CIR) process with long-run mean $\theta$, rate of mean reversion $\kappa$, and “CIR-volatility” $\sigma$. We assume that $\kappa$, $\theta$ and $\sigma$ are positive, and that $2\kappa\theta > \sigma^2$, which ensures that $Z_t > 0$ at all times.

Note that given $Z_t$, the process $Y_t$ in (2.1) appears as an Ornstein–Uhlenbeck (OU) process evolving on the time scale $\epsilon / Z_t$, and with the invariant (or long-run) distribution $\mathcal{N}(m, \nu^2)$. This way of “modulating” the rate of mean reversion of the process $Y_t$ by $Z_t$ has also been used in [4] in the context of interest rate modeling.

Multiple time scales are incorporated in this model through the parameter $\epsilon > 0$, which is intended to be small, so that $Y_t$ is fast-reverting.

We do not specify the precise form of $f(y)$ which will not play an essential role in the asymptotic results derived in this paper. However, in order to ensure $\Sigma_t$ has the same behavior at zero and infinity as in the case of a pure Heston model, we assume there exist constants $c_1$ and $c_2$ such that $0 < c_1 \leq f(y) \leq c_2 < \infty$ for all $y \in \mathbb{R}$. Likewise, the particular choice of an OU-like process for $Y_t$ is not crucial in the analysis. The mean-reversion aspect (or ergodicity) is the important property. In fact, we could have chosen $Y_t$ to be a CIR-like process instead of an OU-like process without changing the nature of the correction to the Heston model presented in the paper.

We note that if one chooses $f(y) = 1$, the multi-scale model becomes $\epsilon$-independent and reduces to the pure Heston model expressed under the risk-neutral measure with stock price $X_t$ and square volatility $Z_t$:

$$
dX_t = r X_t dt + \sqrt{Z_t} X_t dW^x_t,
\quad dZ_t = \kappa (\theta - Z_t) dt + \sigma \sqrt{Z_t} dW^z_t.
\quad d \langle W^x, W^z \rangle_t = \rho_{xz} dt.
$$

Thus, the multi-scale model can be thought of as a Heston-like model with a fast-varying factor of volatility, $f(Y_t)$, build on top of the CIR process $Z_t$, which drives the volatility in the Heston Model.

We consider a European option expiring at time $T > t$ with payoff $h(X_T)$. As the dynamics of the stock in the multi-scale model are specified under the risk-neutral measure, the price of the option, denoted by $P_t$, can be expressed as an expectation of the option payoff, discounted at the risk-free rate:

$$
P_t = \mathbb{E} \left[ e^{-r(T-t)} h(X_T) \mid X_t, Y_t, Z_t \right] =: P^r(t, X_t, Y_t, Z_t),
$$

where we have used the Markov property of $(X_t, Y_t, Z_t)$, and defined the pricing func-
tion $P^\epsilon(t,x,y,z)$, the superscript $\epsilon$ denoting the dependence on the small parameter $\epsilon$. Using the Feynman–Kac formula, $P^\epsilon(t,x,y,z)$ satisfies the following PDE and boundary condition:

$$L^\epsilon P^\epsilon(t,x,y,z) = 0, \quad \text{(2.8)}$$
$$L^\epsilon = \frac{\partial}{\partial t} + L_{(X,Y,Z)} - r, \quad \text{(2.9)}$$
$$P^\epsilon(T,x,y,z) = h(x), \quad \text{(2.10)}$$

where the operator $L_{(X,Y,Z)}$ is the infinitesimal generator of the process $(X_t,Y_t,Z_t)$:

$$L_{(X,Y,Z)} = rx \frac{\partial}{\partial x} + \frac{1}{2} f^2(y)xx^2 \frac{\partial^2}{\partial x^2} + \rho xz \sigma f(y)zx \frac{\partial^2}{\partial x \partial z}$$
$$+ \kappa(\theta - z) \frac{\partial}{\partial z} + \frac{1}{2} \sigma^2 z \frac{\partial^2}{\partial z^2}$$
$$+ \frac{z}{\epsilon} \left( (m - y) \frac{\partial}{\partial y} + \nu \frac{\partial^2}{\partial y^2} \right)$$
$$+ \frac{z}{\sqrt{\epsilon}} \left( \rho yz \sqrt{2} \frac{\partial^2}{\partial y \partial z} + \rho xy \nu \sqrt{2} f(y)x \frac{\partial^2}{\partial x \partial y} \right).$$

It will be convenient to separate $L^\epsilon$ into groups of like-powers of $1/\sqrt{\epsilon}$. To this end, we define the operators $L_0, L_1$ and $L_2$ as follows:

$$L_0 := \nu^2 \frac{\partial^2}{\partial y^2} + (m - y) \frac{\partial}{\partial y}, \quad \text{(2.11)}$$
$$L_1 := \rho yz \sigma \sqrt{2} \frac{\partial^2}{\partial y \partial z} + \rho xy \nu \sqrt{2} f(y)x \frac{\partial^2}{\partial x \partial y}, \quad \text{(2.12)}$$
$$L_2 := \frac{\partial}{\partial t} + \frac{1}{2} f^2(y)xx^2 \frac{\partial^2}{\partial x^2} + r \left( x \frac{\partial}{\partial x} - \cdot \right)$$
$$+ \frac{1}{2} \sigma^2 z \frac{\partial^2}{\partial z^2} + \kappa(\theta - z) \frac{\partial}{\partial z} + \rho xz \sigma f(y)zx \frac{\partial^2}{\partial x \partial z}.$$

With these definitions, $L^\epsilon$ is expressed as:

$$L^\epsilon = \frac{z}{\epsilon} L_0 + \frac{z}{\sqrt{\epsilon}} L_1 + L_2. \quad \text{(2.14)}$$

Note that $L_0$ is the infinitesimal generator of an OU process with unit rate of mean-reversion, and $L_2$ is the pricing operator of the Heston model with volatility and correlation modulated by $f(y)$.

### 3. Asymptotic Analysis.

For a general function $f$, there is no analytic solution to the boundary value problem (2.8–2.10). Thus, we proceed with an asymptotic analysis as developed in [7]. Specifically, we perform a singular perturbation with respect to the small parameter $\epsilon$, expanding our solution in powers of $\sqrt{\epsilon}$:

$$P^\epsilon = P_0 + \sqrt{\epsilon} P_1 + \epsilon P_2 + \ldots. \quad \text{(3.1)}$$

We now plug (3.1) and (2.14) into (2.8) and (2.10), and collect terms of equal powers of $\sqrt{\epsilon}$. 


The Order $1/\epsilon$ Terms. Collecting terms of order $1/\epsilon$ we have the following PDE:

$$0 = zL_0 P_0.$$  \hfill (3.2)

We see from (2.11) that both terms in $L_0$ take derivatives with respect to $y$. In fact, $L_0$ is an infinitesimal generator and consequently zero is an eigenvalue with constant eigenfunctions. Thus, we seek $P_0$ of the form

$$P_0 = P_0(t, x, z),$$

so that (3.2) is satisfied.

The Order $1/\sqrt{\epsilon}$ Terms. Collecting terms of order $1/\sqrt{\epsilon}$ leads to the following PDE

$$0 = zL_0 P_1 + zL_1 P_0 = zL_0 P_1.$$  \hfill (3.3)

Note that we have used that $L_1 P_0 = 0$, since both terms in $L_1$ take derivatives with respect to $y$ and $P_0$ is independent of $y$. As above, we seek $P_1$ of the form

$$P_1 = P_1(t, x, z),$$

so that (3.3) is satisfied.

The Order 1 Terms. Matching terms of order 1 leads to the following PDE and boundary condition:

$$0 = zL_0 P_2 + zL_1 P_1 + L_2 P_0 = zL_0 P_2 + L_2 P_0$$  \hfill (3.4)

$$h(x) = P_0(T, x, z).$$  \hfill (3.5)

In deriving (3.4) we have used that $L_1 P_1 = 0$, since $L_1$ takes derivative with respect to $y$ and $P_1$ is independent of $y$.

Note that (3.4) is a Poisson equation in $y$ with respect to the infinitesimal generator $L_0$ and with source term $L_2 P_0$; in solving this equation, $(t, x, z)$ are fixed parameters. In order for this equation to admit solutions with reasonable growth at infinity, the source term must satisfy the centering condition:

$$0 = \langle L_2 P_0 \rangle = \langle L_2 \rangle P_0,$$  \hfill (3.6)

where we have used the notation

$$\langle g \rangle := \int g(y)\Phi(y)dy,$$  \hfill (3.7)

here $\Phi$ denotes the density of the invariant distribution of the process $Y_t$, which we remind the reader is $\mathcal{N}(m, \nu^2)$. Note that in (3.6) we have pulled $P_0(t, x, z)$ out of the linear $\langle \cdot \rangle$ operator since it does not depend on $y$.

Note that the PDE (3.4) and the boundary condition (3.5) jointly define a boundary value problem that $P_0(t, x, z)$ must satisfy.
Using equation (3.4) and the centering condition (3.6) we deduce:

\[
P_2 = -\frac{1}{z} \mathcal{L}^{-1}_0 (\mathcal{L}_2 - \langle \mathcal{L}_2 \rangle) P_0.
\]  

(3.8)

**The Order \(\sqrt{\varepsilon}\) Terms.** Collecting terms of order \(\sqrt{\varepsilon}\), we obtain the following PDE and boundary condition:

\[
0 = z\mathcal{L}_0 P_3 + z\mathcal{L}_1 P_2 + \mathcal{L}_2 P_1,
\]  

(3.9)

\[
0 = P_1(T, x, z).
\]  

(3.10)

We note that \(P_3(t, x, y, z)\) solves the Poisson equation (3.9) in \(y\) with respect to \(\mathcal{L}_0\). Thus, we impose the corresponding centering condition on the source \(z\mathcal{L}_1 P_2 + \mathcal{L}_2 P_1\), leading to

\[
\langle \mathcal{L}_2 \rangle P_1 = -\langle z\mathcal{L}_1 P_2 \rangle.
\]  

(3.11)

Plugging \(P_2\), given by (3.8), into equation (3.11) gives:

\[
\langle \mathcal{L}_2 \rangle P_1 = AP_0,
\]  

(3.12)

\[
A := \left\langle z\mathcal{L}_1 \frac{1}{z} \mathcal{L}^{-1}_0 (\mathcal{L}_2 - \langle \mathcal{L}_2 \rangle) \right\rangle.
\]  

(3.13)

Note that the PDE (3.12) and the zero boundary condition (3.10) define a boundary value problem that \(P_1(t, x, z)\) must satisfy.

**Summary of the Key Results.** We summarize the key results of our asymptotic analysis. We have written the expansion (3.1) for the solution of the PDE problem (2.8–2.10). Along the way, he have chosen solutions for \(P_0\) and \(P_1\) which are of the form \(P_0 = P_0(t, x, z)\) and \(P_1 = P_1(t, x, z)\). These choices lead us to conclude that \(P_0(t, x, z)\) and \(P_1(t, x, z)\) must satisfy the following boundary value problems

\[
\langle \mathcal{L}_2 \rangle P_0 = 0,
\]  

(3.14)

\[
P_0(T, x, z) = h(x),
\]  

(3.15)

and

\[
\langle \mathcal{L}_2 \rangle P_1(t, x, z) = AP_0(t, x, z),
\]  

(3.16)

\[
P_1(T, x, z) = 0,
\]  

(3.17)

where

\[
\langle \mathcal{L}_2 \rangle = \frac{\partial}{\partial t} + \frac{1}{2} \langle f^2 \rangle z x^2 \frac{\partial^2}{\partial x^2} + r \left( x \frac{\partial}{\partial x} - \right)
\]

\[
+ \frac{1}{2} \sigma^2 z \frac{\partial^2}{\partial z^2} + \kappa (\theta - z) \frac{\partial}{\partial z} + \rho z \sigma \langle f \rangle z x \frac{\partial^2}{\partial x \partial z},
\]  

(3.18)

and \(A\) is given by (3.13). Recall that the bracket notation is defined in (3.7).

**4. Formulas for \(P_0(t, x, z)\) and \(P_1(t, x, z)\).** In this section we use the results of our asymptotic calculations to find explicit solutions for \(P_0(t, x, z)\) and \(P_1(t, x, z)\).
4.1. Formula for $P_0(t, x, z)$. Recall that $P_0(t, x, z)$ satisfies a boundary value problem defined by equations (3.14) and (3.15).

Without loss of generality, we normalize $f$ so that $\langle f^2 \rangle = 1$. Thus, we rewrite $\langle L_2 \rangle$ given by (3.18) as follows:

$$\langle L_2 \rangle = \frac{\partial}{\partial t} + \frac{1}{2} \sigma^2 z^2 \frac{\partial^2}{\partial x^2} + r \left( x \frac{\partial}{\partial x} - \cdot \right) + \frac{1}{2} \kappa (\theta - z) \frac{\partial}{\partial z} + \rho \sigma x z \frac{\partial^2}{\partial x \partial z}, \quad (4.1)$$

$$:= L_H,$$

$$\rho := \rho_{xz} \langle f \rangle. \quad (4.2)$$

We note that $\rho^2 \leq 1$ since $\langle f \rangle^2 \leq \langle f^2 \rangle = 1$. So, $\rho$ can be thought of as an effective correlation between the Brownian motions in the Heston model obtained in the limit $\epsilon \to 0$, where $\langle L_2 \rangle = L_H$, the pricing operator for European options as calculated in the Heston model. Thus, we see that $P_0(t, x, z) := P_H(t, x, z)$ is the classical solution for the price of a European option as calculated in the Heston model with effective correlation $\rho = \rho_{xz} \langle f \rangle$.

The derivation of pricing formulas for the Heston model is given in Appendix A. Here, we simply state the main result:

$$P_H(t, x, z) = e^{-r \tau} \frac{1}{2 \pi} \int e^{-ikq} \hat{G}(\tau, k, z) \hat{h}(k) dk,$$  \(4.3\)

$$\tau(t) = T - t,$$  \(4.4\)

$$q(t, x) = r(T - t) + \log x,$$  \(4.5\)

$$\hat{h}(k) = \int e^{ikq} h(e^q) dq,$$  \(4.6\)

$$\hat{G}(\tau, k, z) = e^{C(\tau, k) + z D(\tau, k)},$$  \(4.7\)

$$C(\tau, k) = \frac{\kappa \theta}{\sigma^2} \left( (\kappa + \rho i k \sigma + d(k)) \tau - 2 \log \left( \frac{1 - g(k)e^{\tau d(k)}}{1 - g(k)} \right) \right),$$  \(4.8\)

$$D(\tau, k) = \frac{\kappa + \rho i k \sigma + d(k)}{\sigma^2} \left( \frac{1 - e^{\tau d(k)}}{1 - g(k)e^{\tau d(k)}} \right),$$  \(4.9\)

$$d(k) = \sqrt{\sigma^2(k^2 - i k) + (\kappa + \rho i k \sigma)^2},$$  \(4.10\)

$$g(k) = \frac{\kappa + \rho i k \sigma + d(k)}{\kappa + \rho i k \sigma - d(k)}.$$  \(4.11\)

4.2. Formula for $P_1(t, x, z)$. Recall that $P_1(t, x, z)$ satisfies a boundary value problem defined by equations (3.16) and (3.17). In order to find a solution for $P_1(t, x, z)$ we must first identify the operator $A$. To this end, we introduce two functions, $\phi(y)$ and $\psi(y)$, which solve the following Poisson equations in $y$ with respect to the operator $L_0$:

$$L_0 \phi = \frac{1}{2} \left( f^2 - \langle f^2 \rangle \right),$$  \(4.12\)

$$L_0 \psi = f - \langle f \rangle.$$  \(4.13\)
From equation (3.13) we have:
\[
A = \left\langle z \mathcal{L}_1 \frac{1}{z} \mathcal{L}_0^{-1} (\mathcal{L}_2 - \langle \mathcal{L}_2 \rangle) \right\rangle \\
= \left\langle z \mathcal{L}_1 \frac{1}{z} \mathcal{L}_0^{-1} \left( f^2 - \langle f^2 \rangle \right) x^2 \frac{\partial^2}{\partial x^2} \right\rangle \\
+ \left\langle z \mathcal{L}_1 \frac{1}{z} \mathcal{L}_0^{-1} \rho_{xz} \sigma z (f - \langle f \rangle) x \frac{\partial^2}{\partial x \partial z} \right\rangle \\
= z \left\langle \mathcal{L}_1 \phi(y) x^2 \frac{\partial^2}{\partial x^2} \right\rangle + \rho_{xz} \sigma z \left\langle \mathcal{L}_1 \psi(y) x \frac{\partial^2}{\partial x \partial z} \right\rangle.
\]

Using the definition (2.12) of \( \mathcal{L}_1 \), one deduces the following expression for \( A \):
\[
A = V_1 z x^2 \frac{\partial^3}{\partial z \partial x^2} + V_2 z x^2 \frac{\partial^3}{\partial z^2 \partial x} \\
+ V_3 z x \frac{\partial}{\partial x} \left( x \frac{\partial^2}{\partial x^2} \right) + V_4 z \frac{\partial}{\partial z} \left( x \frac{\partial}{\partial x} \right)^2,
\]
(4.14)
\[
V_1 = \rho_{yz} \sigma \sqrt{2} \langle \phi' \rangle, \\
V_2 = \rho_{xx} \rho_{yz} \sigma^2 \nu \sqrt{2} \langle \psi' \rangle, \\
V_3 = \rho_{xy} \nu \sqrt{2} \langle f \phi' \rangle, \\
V_4 = \rho_{xy} \rho_{xz} \sigma \nu \sqrt{2} \langle f \psi' \rangle.
\]
(4.15-4.18)

Note that we have introduced four group parameters, \( V_i, i = 1 \ldots 4 \), which are constants, and can be obtained by calibrating our model to the market as will be done in Section 7.

Now that we have expressions for \( A, P_H, \) and \( L_H \), we are in a position to solve for \( P_1(t, x, z) \), which is the solution to the boundary value problem defined by equations (3.16) and (3.17). We leave the details of the calculation to Appendix B. Here, we
simply present the main result.

\[ P_1(t, x, z) = \frac{e^{-r\tau}}{2\pi} \int_{\mathbb{R}} e^{-ikq} \left( \kappa \theta \hat{f}_0(\tau, k) + z \hat{f}_1(\tau, k) \right) \times \hat{G}(\tau, k, z) \hat{h}(k) \text{d}k, \]  
\[ \tau(t) = T - t, \]
\[ q(t, x) = r(T - t) + \log x, \]
\[ \hat{h}(k) = \int e^{ikq} h(e^q) \text{d}q, \]
\[ \hat{G}(\tau, k, z) = e^{C(\tau, k) + zD(\tau, k)}, \]
\[ \hat{f}_0(\tau, k) = \int_0^\tau \hat{f}_1(t, k) \text{d}t, \]
\[ \hat{f}_1(\tau, k) = \int_0^\tau b(s, k) e^{A(\tau, k, s)} \text{d}s, \]
\[ C(\tau, k) = \frac{\kappa \theta}{\sigma^2} \left( (\kappa + \rho ik \sigma + d(k)) \tau - 2\log \left( \frac{1 - g(k) e^{r d(k)}}{1 - g(k)} \right) \right), \]
\[ D(\tau, k) = \frac{\kappa + \rho ik \sigma + d(k)}{\sigma^2} \left( \frac{1 - e^{r d(k)}}{1 - g(k) e^{r d(k)}} \right), \]
\[ A(\tau, k, s) = (\kappa + \rho \sigma ik + d(k)) \frac{1 - g(k)}{d(k) g(k)} \log \left( \frac{g(k) e^{r d(k)} - 1}{g(k) e^{s d(k)} - 1} \right) + d(k) (\tau - s), \]
\[ d(k) = \sqrt{\sigma^2 (k^2 - ik) + (\kappa + \rho ik \sigma)^2}, \]
\[ g(k) = \frac{\kappa + \rho ik \sigma + d(k)}{\kappa + \rho ik \sigma - d(k)}, \]
\[ b(\tau, k) = - \left( V_1 D(\tau, k) (-k^2 + ik) + V_2 D^2(\tau, k) (-ik) + V_3 (ik^3 + k^2) + V_4 D(\tau, k) (-k^2) \right). \]

5. Accuracy of the Approximation. In this section, we prove that the approximation \( P^\epsilon \sim P_0 + \sqrt{\epsilon} P_1 \), where \( P_0 \) and \( P_1 \) are defined in the previous sections, is accurate to order \( \epsilon^\alpha \) for any given \( \alpha < 1 \). Specifically, for a European option with a smooth bounded payoff, \( h(x) \), we will show:

\[ |P^\epsilon(t, x, y, z) - (P_0(t, x, z) + \sqrt{\epsilon} P_1(t, x, z))| \leq C \epsilon^\alpha, \]  

(5.1)

where \( C \) is a constant which depends on \((y, z)\), but is independent of \( \epsilon \).

We start by defining the remainder term \( R^\epsilon(t, x, y, z) \):

\[ R^\epsilon = (P_0 + \sqrt{\epsilon} P_1 + \epsilon P_2 + \sqrt{\epsilon} \epsilon P_3) - P^\epsilon. \]  

(5.2)
Recalling that
\begin{align*}
0 &= \mathcal{L}^\varepsilon P^\varepsilon, \\
0 &= z\mathcal{L}_0 P_0, \\
0 &= z\mathcal{L}_0 P_1 + z\mathcal{L}_1 P_0, \\
0 &= z\mathcal{L}_0 P_2 + z\mathcal{L}_1 P_1 + \mathcal{L}_2 P_0, \\
0 &= z\mathcal{L}_0 P_3 + z\mathcal{L}_1 P_2 + \mathcal{L}_2 P_1,
\end{align*}
and applying $\mathcal{L}^\varepsilon$ to $R^\varepsilon$, we obtain that $R^\varepsilon$ must satisfy the following PDE:
\begin{align*}
\mathcal{L}^\varepsilon R^\varepsilon &= \mathcal{L}^\varepsilon \left( P_0 + \sqrt{\varepsilon} P_1 + \varepsilon P_2 + \varepsilon \sqrt{\varepsilon} P_3 \right) - \mathcal{L}^\varepsilon P^\varepsilon \\
&= \left( \frac{z}{\varepsilon} \mathcal{L}_0 + \frac{z}{\sqrt{\varepsilon}} \mathcal{L}_1 + \mathcal{L}_2 \right) \left( P_0 + \sqrt{\varepsilon} P_1 + \varepsilon P_2 + \varepsilon \sqrt{\varepsilon} P_3 \right) \\
&= \varepsilon \left( z\mathcal{L}_1 P_3 + \mathcal{L}_2 P_2 + \sqrt{\varepsilon} \mathcal{L}_2 P_3 \right) \\
&= \varepsilon F^\varepsilon, \quad (5.3)
\end{align*}
where we have defined the $\varepsilon$-dependent source term $F^\varepsilon(t,x,y,z)$. Recalling that
\begin{align*}
P^\varepsilon(T,x,y,z) &= h(x), \\
P_0(T,x,z) &= h(x), \\
P_1(T,x,z) &= 0,
\end{align*}
we deduce from (5.2) that
\begin{align*}
R^\varepsilon(T,x,y,z) &= \varepsilon P_2(T,x,y,z) + \varepsilon \sqrt{\varepsilon} P_3(T,x,y,z) \\
&= \varepsilon G^\varepsilon(x,y,z), \quad (5.5)
\end{align*}
where we have defined the $\varepsilon$-dependent boundary term $G^\varepsilon(x,y,z)$.

Using the expression (2.9) for $\mathcal{L}^\varepsilon$ we find that $R^\varepsilon(t,x,y,z)$ satisfies the following boundary value problem with source:
\begin{align*}
\left( \frac{\partial}{\partial t} + \mathcal{L}_{X,Y,Z} - r \right) R^\varepsilon &= \varepsilon F^\varepsilon, \quad (5.7)
R^\varepsilon(T,x,y,z) &= \varepsilon G^\varepsilon(x,y,z). \quad (5.8)
\end{align*}
Therefore $R^\varepsilon$ admits the following probabilistic representation:
\begin{align*}
R^\varepsilon(t,x,y,z) &= \varepsilon \mathbb{E} \left[ e^{-r(T-t)} G^\varepsilon(X_T,Y_T,Z_T) \\
&- \int_t^T e^{-r(s-t)} F^\varepsilon(s,X_s,Y_s,Z_s) ds \bigg| X_t = x, Y_t = y, Z_t = z \right]. \quad (5.9)
\end{align*}
A careful but tedious analysis shows that $G^\varepsilon$ and $F^\varepsilon$ given by (5.4) and (5.6), can be expressed as combinations of a finite number of derivatives of $P_0$ with respect to $t,x$ and $z$. Using the Fourier representation formula (4.3) for $P_0 = P_H$, the boundedness and smoothness of the payoff $h(x)$, and growth estimates derived in Chapter 5 of [7]
for the solutions of the Poisson equations (4.12) and (4.13) with $f(y)$ bounded, one deduces the following uniform bounds in $\epsilon$:

\[
|G^\epsilon(x, y, z)| \leq C[1 + \log(1 + |y|)(1 + z)], \\
|F^\epsilon(t, x, y, z)| \leq C[1 + \log(1 + |y|)(1 + z^2)].
\]

(5.10) (5.11)

Using (5.9), the bounds (5.10) and (5.11), Cauchy-Schwarz inequality, and moments of the $\epsilon$-independent CIR process $Z_t$ (see for instance [14]), one obtains:

\[
|R^\epsilon(t, x, y, z)| \leq \epsilon C' \left( 1 + \mathbb{E}_{t,y,z} |Y_T| + \int_t^T \mathbb{E}_{t,y,z} |Y_s| ds \right),
\]

(5.12)

where $\mathbb{E}_{t,y,z}$ denotes the expectation starting at time $t$ from $Y_t = y$ and $Z_t = z$ under the dynamics (2.3)–(2.4). Under this dynamics, starting at time zero from $y$, we have

\[
Y_t = m + (y - m)e^{-\frac{1}{2} \int_0^t Z_s ds} + \frac{\nu \sqrt{2}}{\sqrt{\epsilon}} \int_0^t e^{-\frac{1}{2} \int_0^u Z_s du} \sqrt{Z_s} dW_y.
\]

(5.13)

Using the bound established in Appendix C we have that for any given $\alpha < 1$ there is a constant $C$ such that.

\[
\mathbb{E}|Y_t| \leq C \epsilon^{\alpha-1},
\]

(5.14)

and the error estimate (5.1) follows.

6. The Multi-Scale Implied Volatility Surface. In this section, we explore how the implied volatility surface produced by our multi-scale model compares to that produced by the Heston model. To begin, we remind the reader that an approximation to the price of a European option in the multi-scale model can be obtained through the formula

\[
P^\epsilon \sim P_0 + \sqrt{\epsilon} P_1 \\
= P_H + P_1^F, \\
P_1^F := \sqrt{\epsilon} P_1,
\]

where we have absorbed the $\sqrt{\epsilon}$ into the definition of $P_1^F$ and used $P_0 = P_H$, the Heston price. Form the formulas for the correction $P_1$, given in Section 4.2, it can be seen that $P_1$ is linear in $V_i$, $i = 1, \ldots, 4$. Therefore by setting

\[
V_i^\epsilon = \sqrt{\epsilon} V_i \quad i = 1 \ldots 4,
\]

the small correction $P_1^F$ is given by the same formulas as $P_1$ with the $V_i$ replaced by the $V_i^\epsilon$.

It is important to note that, although adding a fast mean-reverting factor of volatility on top of the Heston model introduces five new parameters ($\nu$, $m$, $\epsilon$, $\rho_{xy}$, $\rho_{yz}$) plus an unknown function $f$ to the dynamics of the stock (see (2.2) and (2.3)), neither knowledge of the values of these five parameters, nor the specific form of the function $f$ is required to price options using our approximation. The effect of adding a fast mean-reverting factor of volatility on top of the Heston model is entirely captured by the four group parameters $V_i^\epsilon$, which are constants that can be obtained by calibrating the multi-scale model to option prices on the market.
By setting $V^x_i = 0$ for $i = 1, \cdots, 4$, we see that $P^x_1 = 0$, $P^x = P_H$, and the resulting implied volatility surface, obtained by inverting Black-Scholes formula, corresponds to the implied volatility surface produced by the Heston model. If we then vary a single $V^x_i$ while holding $V^x_j = 0$ for $j \neq i$, we can see exactly how the multi-scale implied volatility surface changes as a function of each of the $V^x_i$. The results of this procedure are plotted in Figure 6.1.

![Figure 6.1](image)

**Fig. 6.1.** Implied volatility curves are plotted as a function of the strike price for European calls in the multi-scale model. In this example the initial stock price is $x = 100$. The Heston parameters are set to $\theta = 0.04$, $\gamma = 0.24$, $\kappa = 3.4$, $\sigma = 0.39$, $\rho_{xz} = -0.64$ and $r = 0.0$. In subfigure (a) we vary only $V^x_1$, fixing $V^x_i = 0$ for $i \neq 1$. Likewise, in subfigures (b), (c) and (d) we vary only $V^x_2$, $V^x_3$ and $V^x_4$ respectively, fixing all other $V^x_i = 0$. We remind the reader that, in all four plots, $V^x_i = 0$ corresponds to the implied volatility curve of the Heston model.

Because they are on the order of $\sqrt{\epsilon}$, typical values of the $V^x_i$ are quite small. However, in order to highlight their effect on the implied volatility surface, the range of values plotted for the $V^x_i$ in Figure 6.1 was intentionally chosen to be large. It is clear from Figure 6.1 that each $V^x_i$ has a distinct effect on the implied volatility surface. Thus, the multi-scale model provides considerable flexibility when it comes to calibrating the model to the implied volatility surface produced by options on the market.
7. Calibration. Denote by \( \Theta \) and \( \Phi \) the vectors of unobservable parameters in the Heston and Multicale approximation models respectively.

\[
\Theta = (\kappa, \rho, \sigma, \theta, z), \\
\Phi = (\kappa, \rho, \sigma, \theta, z, V_1^\tau, V_2^\tau, V_3^\tau, V_4^\tau).
\]

Let \( \sigma(T_i, K_{j(i)}) \) be the implied volatility of a call option on the market with maturity date \( T_i \) and strike price \( K_{j(i)} \). Note that, for each maturity date, \( T_i \), the set of available strikes, \( \{K_{j(i)}\} \), varies. Let \( \sigma_H(T_i, K_{j(i)}, \Theta) \) be the implied volatility of a call option with maturity date \( T_i \) and strike price \( K_{j(i)} \) as calculated in the Heston model using parameters \( \Theta \). And let \( \sigma_M(T_i, K_{j(i)}, \Phi) \) be the implied volatility of call option with maturity date \( T_i \) and strike price \( K_{j(i)} \) as calculated in the multi-scale approximation using parameters \( \Phi \).

We formulate the calibration problem as a constrained, nonlinear, least-squares optimization. Define the objective functions as

\[
\Delta_2^H(\Theta) = \sum_i \sum_{j(i)} (\sigma(T_i, K_{j(i)}) - \sigma_H(T_i, K_{j(i)}, \Theta))^2, \\
\Delta_2^M(\Phi) = \sum_i \sum_{j(i)} (\sigma(T_i, K_{j(i)}) - \sigma_M(T_i, K_{j(i)}, \Phi))^2.
\]

We consider \( \Theta^\ast \) and \( \Phi^\ast \) to be optimal if they satisfy

\[
\Delta_2^H(\Theta^\ast) = \min_{\Theta} \Delta_2^H(\Theta), \\
\Delta_2^M(\Phi^\ast) = \min_{\Phi} \Delta_2^M(\Phi).
\]

It is well-known that that the objective functions, \( \Delta_2^H \) and \( \Delta_2^M \), may exhibit a number of local minima. Therefore, if one uses a local gradient method to find \( \Theta^\ast \) and \( \Phi^\ast \) (as we do in this paper), there is a danger of ending up in a local minima, rather than the global minimum. Therefore, it becomes important to make a good initial guess for \( \Theta \) and \( \Phi \). Luckily, people familiar with the Heston model will know what constitutes a good initial guess for \( \Theta \). In this paper, we calibrate the Heston model first to find \( \Theta^\ast \). Then, for the multi-scale model we make an initial guess \( \Phi = (\Theta^\ast, 0, 0, 0, 0) \) (i.e. we set the \( V_1^\tau, V_2^\tau, V_3^\tau, V_4^\tau \) = 0 and use \( \Theta^\ast \) for the rest of the parameters of \( \Phi \)). This is a logical calibration procedure because the \( V_1^\tau, \) being of order \( \sqrt{\tau} \), are intended to be small parameters.

The data we consider consists of call options on the S&P500 index (SPX) taken from May 17, 2006. We limit our data set to options with maturities greater than 45 days, and with open interest greater than 100. We use the yield on the nominal 3-month, constant maturity, U.S. Government treasury bill as the risk-free interest rate. And we use a dividend yield on the S&P 500 index taken directly from the Standard & Poor’s website (www.standardandpoors.com). In Figures 7.1 through 7.7 we plot the implied volatilities of call options on the market, as well as the calibrated implied volatility curves for the Heston and multi-scale models. We would like to emphasize that, although the plots are presented maturity by maturity, they are the result of a single calibration procedure that uses the entire data set.

From Figures 7.1 through 7.7 it is apparent to the naked eye that the multi-scale model represents a vast improvement over the Heston model—especially, for call options with the shortest maturities. In order to quantify this result we define marginal
residual sum of squares

\[ \bar{\Delta}^2_H(T_i) = \frac{1}{N(T_i)} \sum_{j(i)} \left( \sigma(T_i, K_{j(i)}) - \sigma_H(T_i, K_{j(i)}, \Theta^*) \right)^2, \]

\[ \bar{\Delta}^2_M(T_i) = \frac{1}{N(T_i)} \sum_{j(i)} \left( \sigma(T_i, K_{j(i)}) - \sigma_M(T_i, K_{j(i)}, \Phi^*) \right)^2, \]

where \( N(T_i) \) is the number of different calls in the data set that expire at time \( T_i \) (i.e. \( N(T_i) = \#\{ K_{j(i)} \} \)). A comparison of \( \bar{\Delta}^2_H(T_i) \) and \( \bar{\Delta}^2_M(T_i) \) is given in Table 7.1. The table confirms what is apparent to the naked eye—namely, that the multi-scale model fits the market data significantly better than the Heston model for the two shortest maturities, as well as the longest maturity.

<table>
<thead>
<tr>
<th>( T_i - t ) (days)</th>
<th>( \bar{\Delta}^2_H(T_i) )</th>
<th>( \bar{\Delta}^2_M(T_i) )</th>
<th>( \bar{\Delta}^2_H(T_i)/\bar{\Delta}^2_M(T_i) )</th>
</tr>
</thead>
<tbody>
<tr>
<td>65</td>
<td>( 2.94 \times 10^{-5} )</td>
<td>( 7.91 \times 10^{-6} )</td>
<td>3.71</td>
</tr>
<tr>
<td>121</td>
<td>( 1.02 \times 10^{-5} )</td>
<td>( 3.72 \times 10^{-6} )</td>
<td>2.73</td>
</tr>
<tr>
<td>212</td>
<td>( 0.41 \times 10^{-5} )</td>
<td>( 8.11 \times 10^{-6} )</td>
<td>0.51</td>
</tr>
<tr>
<td>303</td>
<td>( 0.40 \times 10^{-5} )</td>
<td>( 3.51 \times 10^{-6} )</td>
<td>1.12</td>
</tr>
<tr>
<td>394</td>
<td>( 0.74 \times 10^{-5} )</td>
<td>( 5.17 \times 10^{-6} )</td>
<td>1.42</td>
</tr>
<tr>
<td>583</td>
<td>( 1.13 \times 10^{-5} )</td>
<td>( 9.28 \times 10^{-6} )</td>
<td>1.22</td>
</tr>
<tr>
<td>947</td>
<td>( 0.33 \times 10^{-5} )</td>
<td>( 1.47 \times 10^{-6} )</td>
<td>2.25</td>
</tr>
</tbody>
</table>

Table 7.1 Residual sum of squares for the Heston and the Multi-Scale models at several maturities.

Fig. 7.1. SPX Implied Volatilities from May 17, 2006

Appendix A. Heston Stochastic Volatility Model. There are a number of
excellent resources where one can read about the Heston stochastic volatility model—so many, in fact, that a detailed review of the model would seem superfluous. However, in order to establish some notation, we will briefly review the dynamics of the Heston model here, as well as show our preferred method for solving the corresponding European option pricing problem. The notes from this section closely follow [19]. The reader should be aware that a number of the equations developed in this section are
Let $X_t$ be the price of a stock. And denote by $r$ the risk-free rate of interest. Then, under the risk-neutral probability measure, $\mathbb{P}$, the Heston model takes the following
form:

\[ dX_t = r X_t dt + \sqrt{Z_t} X_t dW^x_t , \]
\[ dZ_t = \kappa (\theta - Z_t) dt + \sigma \sqrt{Z_t} dW^z_t , \]
\[ d\langle W^x , W^z \rangle_t = \rho dt. \]
Here, $W^x_t$ and $W^z_t$ are one-dimensional standard Brownian motions with correlation $\rho$, such that $|\rho| \leq 1$. The process, $Z_t$, is the stochastic variance of the stock. And, $\kappa$, $\theta$ and $\sigma$ are positive constants satisfying $2\kappa\theta > \sigma^2$; this ensures that $Z_t$ remains positive for all $t$.

We denote by $P_H$ the price of a European option, as calculated under the Heston framework. As we are already under the risk-neutral measure, we can express $P_H$ as an expectation of the option payoff, $h(X_T)$, discounted at the risk-free rate.

\[
P_H(t, x, z) = \mathbb{E} \left[ e^{-r(T-t)} h(X_T) \middle| X_t = x, Z_t = z \right].
\]

Using the Feynman-Kac formula, we find that $P_H(t, x, z)$ must satisfy the following PDE and boundary condition:

\[
\mathcal{L}_H P_H(t, x, z) = 0,
\]

\[
P_H(T, x, z) = h(x),
\]

\[
\mathcal{L}_H = \frac{\partial}{\partial t} - r + rx \frac{\partial}{\partial x} + \frac{1}{2} \sigma^2 \frac{\partial^2}{\partial x^2} + \kappa (\theta - z) \frac{\partial}{\partial z} + \rho \sigma z \frac{\partial^2}{\partial x \partial z}.
\]

In order to find a solution for $P_H(t, x, z)$, it will be convenient to transform variables as follows:

\[
\tau(t) = T - t,
\]

\[
q(t, x) = r(T - t) + \log x,
\]

\[
P_H(t, x, z) = P'_H(\tau(t), q(t, x), z) e^{-r\tau(t)}.
\]

This transformation leads us to the following PDE and boundary condition for $P'_H(\tau, q, z)$:

\[
\mathcal{L}'_H P'_H(\tau, q, z) = 0,
\]

\[
\mathcal{L}'_H = -\frac{\partial}{\partial \tau} + \frac{1}{2} \sigma^2 \left( \frac{\partial^2}{\partial q^2} - \frac{\partial}{\partial q} \right) + \rho \sigma z \frac{\partial^2}{\partial q \partial z} + \frac{1}{2} \sigma^2 
\]

\[
+ \kappa (\theta - z),
\]

\[
P'_H(0, q, z) = h(e^q).
\]

We will find a solution for $P'_H$ through the method of Green’s functions. Denote by $\delta(q)$ the Dirac delta function, and let $G(\tau, q, z)$, the Green’s function, be the solution to the following boundary value problem:

\[
\mathcal{L}'_H G(\tau, q, z) = 0, \quad \mathcal{L}'_H G(0, q, z) = \delta(q).
\]

Then,

\[
P'_H(\tau, q, z) = \int_{\mathbb{R}} G(\tau, q - p, z) h(e^p) dp.
\]
Now, let $\hat{P}_H(\tau, k, z)$, $\hat{G}(\tau, k, z)$ and $\hat{h}(k)$ be the Fourier transforms of $P'_H(\tau, q, z)$, $G(\tau, q, z)$ and $h(e^q)$ respectively.

\[
\hat{P}_H(\tau, k, z) = \int_{\mathbb{R}} \mathrm{e}^{ikq} P'_H(\tau, q, z) dq,
\]
\[
\hat{G}(\tau, k, z) = \int_{\mathbb{R}} \mathrm{e}^{ikq} G(\tau, q, z) dq,
\]
\[
\hat{h}(k) = \int_{\mathbb{R}} \mathrm{e}^{ikq} h(e^q) dq.
\]

Then, using the convolution property of Fourier transforms we have:

\[
P'_H(\tau, q, z) = \frac{1}{2\pi} \int_{\mathbb{R}} \mathrm{e}^{-ikq} \hat{P}_H(\tau, k, z) dk
\]
\[
= \frac{1}{2\pi} \int_{\mathbb{R}} \mathrm{e}^{-ikq} \hat{G}(\tau, k, z) \hat{h}(k) dk.
\]

Multiplying equations (A.5) and (A.6) by $\mathrm{e}^{ikq'}$ and integrating over $\mathbb{R}$ in $q'$, we find that $\hat{G}(\tau, k, z)$ satisfies the following boundary value problem:

\[
\hat{L}_H \hat{G}(\tau, k, z) = 0,
\]  
(A.7)

\[
\hat{L}_H = -\frac{\partial}{\partial \tau} + \frac{1}{2} z (-k^2 + ik) + \frac{1}{2} \sigma^2 z \frac{\partial^2}{\partial z^2} + (\kappa \theta - (\kappa + \rho \sigma ik) z) \frac{\partial}{\partial z},
\]

\[
\hat{G}(0, k, z) = 1.
\]  
(A.8)

Now, an ansatz: suppose $\hat{G}(\tau, k, z)$ can be written as follows:

\[
\hat{G}(\tau, k, z) = \mathrm{e}^{C(\tau, k) + zD(\tau, k)}.
\]  
(A.9)

Substituting (A.9) into (A.7) and (A.8), and collecting terms of like-powers of $z$, we find that $C(\tau, k)$ and $D(\tau, k)$ must satisfy the following ODE's

\[
\frac{dC}{d\tau}(\tau, k) = \kappa \theta D(\tau, k),
\]  
(A.10)

\[
C(0, k) = 0,
\]  
(A.11)

\[
\frac{dD}{d\tau}(\tau, k) = \frac{1}{2} \sigma^2 D^2(\tau, k) - (\kappa + \rho \sigma ik) D(\tau, k) + \frac{1}{2} (-k^2 + ik),
\]  
(A.12)

\[
D(0, k) = 0.
\]  
(A.13)

Equations (A.10), (A.11), (A.12) and (A.13) can be solved analytically. Their solutions, as well as the final solution to the European option pricing problem in the Heston framework, are given in (4.3–4.11).

Appendix B. Detailed solution for $P_1(t, x, z)$. In this section, we show how to solve for $P_1(t, x, z)$, which is the solution to the boundary value problem defined by equations (3.10) and (3.11). For convenience, we repeat these equations here with...
the notation $\mathcal{L}_H = \langle L \rangle$ and $P_H = P_0$:

$$\mathcal{L}_H P_1(t, x, z) = A P_H(t, x, z), \quad (B.1)$$

$$P_1(T, x, z) = 0. \quad (B.2)$$

We remind the reader that $A$ is given by equation (4.14), $L$ is given by equation (4.1), and $P_H(t, x, z)$ is given by equation (4.3). It will be convenient in our analysis to make the following variable transformation:

$$P_1(t, x, z) = P'_1(\tau(t), q(t), z) e^{-r \tau}, \quad (B.3)$$

$$\tau(t) = T - t, \quad q(t) = r(T - t) + \log x.$$ 

We now substitute equations (4.3), (4.14) and (B.3) into equations (B.1) and (B.2), which leads us to the following PDE and boundary condition for $P'_1(\tau, q, z)$:

$$\mathcal{L}'_H P'_1(\tau, q, z) = A' \frac{1}{2\pi} \int e^{-ikq} \hat{G}(\tau, k, z) \hat{h}(k) dk, \quad (B.4)$$

$$L'_{H} = -\frac{\partial}{\partial \tau} + \frac{1}{2} z \left( \frac{\partial^2}{\partial q^2} - \frac{\partial}{\partial q} \right) + \rho \sigma z \frac{\partial^2}{\partial q \partial z} + \frac{1}{2} \sigma^2 z \frac{\partial^2}{\partial z^2} + \kappa(\theta - z), \quad (B.5)$$

$$A' = V_1 z \frac{\partial}{\partial z} \left( \frac{\partial^2}{\partial q^2} - \frac{\partial}{\partial q} \right) + V_2 z \frac{\partial^3}{\partial z^2 \partial q} + V_3 z \left( \frac{\partial^3}{\partial q^3} - \frac{\partial^2}{\partial q^2} \right) + V_4 z \frac{\partial^3}{\partial z \partial q^2}.$$ 

$$P'_1(0, q, z) = 0.$$ 

Now, let $\hat{P}_1(\tau, k, z)$ be the fourier transform of $P'_1(\tau, q, z)$

$$\hat{P}_1(\tau, k, z) = \int_{\mathbb{R}} e^{ikq} P'_1(\tau, q, z) dq.$$ 

Then,

$$P'_1(\tau, q, z) = \frac{1}{2\pi} \int_{\mathbb{R}} e^{-ikq} \hat{P}_1(\tau, k, z) dk. \quad (B.6)$$

Multiplying equations (B.4) and (B.5) by $e^{ikq'}$ and integrating in $q'$ over $\mathbb{R}$, we find
that $\hat{P}_1(\tau, k, z)$ satisfies the following boundary value problem:

$$\hat{\mathcal{L}}_H \hat{P}_1(\tau, k, z) = \hat{A} \hat{G}(\tau, k, z) \hat{h}(k),$$

(B.7)

$$\hat{\mathcal{L}}_H = -\frac{\partial}{\partial \tau} + \frac{1}{2} z \left( -k^2 + ik \right) + \frac{1}{2} \sigma^2 \frac{\partial^2}{\partial z^2}$$

$$+ (\kappa \theta - (\kappa + \rho \sigma i k) z) \frac{\partial}{\partial z},$$

$$\hat{A} = z \left( V_1 \frac{\partial}{\partial z} \left( -k^2 + ik \right) + V_2 \frac{\partial^2}{\partial z^2} (-ik)$$

$$+ V_3 (ik^3 + k^2) + V_4 \frac{\partial}{\partial z} \left( -k^2 \right) \right),$$

$$\hat{P}_1(0, k, z) = 0.$$  

(B.8)

Now, an ansatz: we suppose that $\hat{P}_1(\tau, k, z)$ can be written as

$$\hat{P}_1(\tau, k, z) = \left( \kappa \theta \hat{f}_0(\tau, k) + z\hat{f}_1(\tau, k) \right) \hat{G}(\tau, k, z) \hat{h}(k).$$

(B.9)

We substitute (B.9) into (B.7) and (B.8). After a good deal of algebra (and in particular, making use of (A.10) and (A.12)), we find that $\hat{f}_0(\tau, k)$ and $\hat{f}_1(\tau, k)$ satisfy the following system of ODE’s:

$$\frac{d \hat{f}_1}{d\tau}(\tau, k) = a(\tau, k) \hat{f}_1(\tau, k) + b(\tau, k),$$

(B.10)

$$\hat{f}_1(0, k) = 0,$$  

(B.11)

$$\frac{d \hat{f}_0}{d\tau}(\tau, k) = \hat{f}_1(\tau, k),$$

(B.12)

$$\hat{f}_0(0, k) = 0,$$  

(B.13)

$$a(\tau, k) = \sigma^2 D(\tau, k) - (\kappa + \rho \sigma i k),$$

$$b(\tau, k) = - (V_1 D(\tau, k) (-k^2 + ik) + V_2 D^2(\tau, k) (-ik)$$

$$+ V_3 (ik^3 + k^2) + V_4 D(\tau, k) (-k^2)),$$

where $D(\tau, k)$ is given by equation (4.9).

Equations (B.10–B.13) can be solved analytically (to the extent that their solutions can be written down in integral form). The solutions for $\hat{f}_0(\tau, k)$ and $\hat{f}_1(\tau, k)$, along with the final solution for $P_1(t, x, z)$, are given by (4.19–4.23).

Appendix C. Moment Estimate for $Y_t$. In this section we derive a moment estimate for $Y_t$, whose dynamics under the pricing measure are given by equations (2.3, 2.4, 2.7). Specifically, we show that for all $\alpha < 1$ there exists a constant, $C$ (which depends on $\alpha$), such that $E|Y_t| \leq C e^{\alpha t}$. First, we define the continuous, strictly increasing, non-negative process, $\beta_t$, as

$$\beta_t := \int_0^t Z_s ds.$$  

(C.1)
Next, we decompose $W_t^y$ as
\[ W_t^y = \rho y_t^z W_t^z + \sqrt{1 - \rho^2} W_t^\perp, \tag{C.2} \]
where $W_t^\perp$ is a standard Brownian motion which is independent of $W_t^z$. Using equations (5.13), (C.1), and (C.2) we derive
\[ |Y_t| \leq C_1 + \frac{C_2}{\sqrt{\epsilon}} \left( \left| \int_0^t e^{-\frac{1}{2}(\beta_t - \beta_s)} \sqrt{Z_s} dW_s^\perp \right| + \left| \int_0^t e^{-\frac{1}{2}(\beta_t - \beta_s)} \sqrt{Z_s} dW_s^\parallel \right| \right), \tag{C.3} \]
where $C_1$ and $C_2$ are constants. For the second stochastic integral, we have:
\[
\frac{1}{\epsilon} \mathbb{E} \left[ \left( \int_0^t e^{-\frac{1}{2}(\beta_t - \beta_s)} \sqrt{Z_s} dW_s^\parallel \right)^2 \right] \leq \frac{1}{\epsilon} \mathbb{E} \left[ e^{-2\beta_t/\epsilon} \mathbb{E} \left[ \left( \int_0^t e^{\beta_s/\epsilon} \sqrt{Z_s} dW_s^\parallel \right)^2 | Z \right] \right] \\
= \frac{1}{\epsilon} \mathbb{E} \left[ e^{-2\beta_t/\epsilon} \mathbb{E} \left[ \int_0^t e^{2\beta_s/\epsilon} Z_s ds | Z \right] \right] \\
= \frac{1}{\epsilon} \mathbb{E} \left[ e^{-2\beta_t/\epsilon} \mathbb{E} \left[ \int_0^t e^{2\beta_s/\epsilon} d\beta | Z \right] \right] \\
= \frac{1}{\epsilon} \mathbb{E} \left[ e^{-2\beta_t/\epsilon} \frac{\epsilon}{2} (e^{2\beta_t/\epsilon} - 1) \right] \\
= \frac{1}{2} \left[ 1 - e^{-2\beta_t/\epsilon} \right] \leq \frac{1}{2}.
\]
Then, by the Cauchy-Schwarz inequality, we see that
\[ \frac{1}{\sqrt{\epsilon}} \mathbb{E} \left[ \int_0^t e^{-\frac{1}{2}(\beta_t - \beta_s)} \sqrt{Z_s} dW_s^\parallel \right] \leq \frac{1}{\sqrt{2}}. \tag{C.4} \]
We now consider the first moment of the first stochastic integral in (C.3):
\[ A := \frac{1}{\sqrt{\epsilon}} \mathbb{E} \left[ \int_0^t e^{-\frac{1}{2}(\beta_t - \beta_s)} \sqrt{Z_s} dW_s^\parallel \right]. \tag{C.5} \]
Naively, one might try to use the Cauchy-Schwarz inequality in the following manner
\[
A \leq \frac{1}{\sqrt{\epsilon}} \sqrt{\mathbb{E} \left[ e^{-2\beta_s/\epsilon} \right]} \sqrt{\mathbb{E} \left[ \int_0^t e^{2\beta_s/\epsilon} Z_s ds \right]} \\
= \frac{\sqrt{\epsilon}}{2} \sqrt{\mathbb{E} \left[ e^{-2\beta_s/\epsilon} \right]} \sqrt{\mathbb{E} \left[ e^{2\beta_s/\epsilon} \right]} - 1.
\]
However, $\mathbb{E} \left[ e^{2\beta_s/\epsilon} \right] = \infty$ for $\epsilon$ small enough. Seeking a more refined approach of bounding $A$, we note that
\[ \frac{1}{\sqrt{\epsilon}} \int_0^t e^{-\frac{1}{2}(\beta_t - \beta_s)} \sqrt{Z_s} dW_s^\parallel = \frac{1}{\sigma \sqrt{\epsilon}} (Z_t - z) e^{-\beta_t/\epsilon} - \frac{\kappa}{\sigma \sqrt{\epsilon}} \int_0^t e^{-\frac{1}{2}(\beta_t - \beta_s)} (\theta - Z_s) ds \\
+ \frac{1}{\sigma \epsilon^{3/2}} \int_0^t e^{-\frac{1}{2}(\beta_t - \beta_s)} Z_s (Z_t - Z_s) ds,
\]
which can be derived by replacing $t$ by $s$ in equation (2.4), multiplying by $e^{\beta_t/\epsilon}$, integrating from 0 to $t$, and using $Z_s^2 = Z_tZ_s - Z_s(Z_t - Z_s)$ and \[ \int_0^t e^{-\frac{t}{t}(\beta_t - \beta_s)} Z_s ds = \epsilon(1 - e^{-\beta_t/\epsilon}). \] Therefore, we have

\[
A \leq \frac{1}{\sigma \sqrt{\epsilon}} \mathbb{E} \left[ |Z_t - z|e^{-\beta_t/\epsilon} \right] + \frac{\kappa}{\sigma \sqrt{\epsilon}} \mathbb{E} \left[ \int_0^t e^{-\frac{t}{t}(\beta_t - \beta_s)} (\theta - Z_s) ds \right] \\
+ \frac{1}{\sigma \epsilon^{3/2}} \mathbb{E} \left[ \int_0^t e^{-\frac{t}{t}(\beta_t - \beta_s)} Z_s(Z_t - Z_s) ds \right].
\]

(C.6)

At this point, we need the moment generating function of $(Z_t, \beta_t)$ which can be derived by replacing $Z_t = \epsilon(1 - e^{-\beta_t/\epsilon})$. For instance, in [14]:

\[
\mathbb{E} \left[ e^{-\lambda_Z t - \mu \beta t} \right] = e^{-\kappa \phi_{\lambda, \mu}(t) - \psi_{\lambda, \mu}(t)},
\]

where

\[
\phi_{\lambda, \mu}(t) = \frac{-2}{\sigma^2} \log \left[ \frac{2 \gamma e^{(\gamma + \kappa)t/2}}{\lambda \sigma^2 (e^{\gamma t} - 1) + \gamma - \kappa + e^{\gamma t}(\gamma + \kappa)} \right],
\]

\[
\psi_{\lambda, \mu}(t) = \lambda \left( \frac{\gamma + \kappa + e^{\gamma t}(\gamma - \kappa)}{\lambda \sigma^2 (e^{\gamma t} - 1) + \gamma - \kappa + e^{\gamma t}(\gamma + \kappa)} \right) + 2 \mu (e^{\gamma t} - 1),
\]

\[
\gamma = \sqrt{\kappa^2 + 2\sigma^2 \mu}.
\]

For the first term in (C.6), using Cauchy-Schwarz, we have

\[
\frac{1}{\sigma \sqrt{\epsilon}} \mathbb{E} \left[ |Z_t - z|e^{-\beta_t/\epsilon} \right] \leq \frac{1}{\sigma \epsilon \sqrt{\epsilon}} \mathbb{E} [Z_t - z] \mathbb{E} [Z_t - z] \mathbb{E} [e^{-2\beta_t/\epsilon}].
\]

(C.7)

From (C.7), one can verify that

\[
\mathbb{E} \left[ |Z_t - z|^2 \right] \leq C_3,
\]

\[
\mathbb{E} \left[ e^{-2\beta_t/\epsilon} \right] = e^{-\kappa \phi_{0,2}(t) - \psi_{0,2}(t)} \sim e^{C_4},
\]

where $C_3$ and $C_4$ are constants, and therefore, for $\epsilon \leq 1$,

\[
\frac{1}{\sigma \sqrt{\epsilon}} \mathbb{E} \left[ |Z_t - z|e^{-\beta_t/\epsilon} \right] \leq C_5,
\]

(C.8)

for some constant $C_5$.

We now turn our attention to the second term in equation (C.6). We have

\[
\frac{\kappa}{\sigma \sqrt{\epsilon}} \mathbb{E} \left[ \int_0^t e^{-\frac{t}{t}(\beta_t - \beta_s)} (\theta - Z_s) ds \right] \\
\leq \frac{\kappa}{\sigma \sqrt{\epsilon}} \mathbb{E} \left[ \int_0^t e^{-\frac{t}{t}(\beta_t - \beta_s)} Z_s ds \right] + \frac{\kappa \theta}{\sigma \sqrt{\epsilon}} \mathbb{E} \left[ \int_0^t e^{-\frac{t}{t}(\beta_t - \beta_s)} ds \right] \\
= \frac{\kappa}{\sigma \sqrt{\epsilon}} \mathbb{E} \left[ \epsilon \left( 1 - e^{-\beta_t/\epsilon} \right) \right] + \frac{\kappa \theta}{\sigma \sqrt{\epsilon}} \int_0^t \mathbb{E} \left[ e^{-\frac{t}{t}(\beta_t - \beta_s)} ds \right] \\
\leq C_6 + \frac{\kappa \theta}{\sigma \sqrt{\epsilon}} \int_0^t \mathbb{E} \left[ e^{-\frac{t}{t}(\beta_t - \beta_s)} \right] ds,
\]

where $C_6$ is a constant.
for some constant $C_0$. To bound the remaining integral, we calculate

$$\mathbb{E} \left[ e^{-\frac{1}{t}(\beta_1-\beta_2)} \right] = \mathbb{E} \left[ e^{-\frac{1}{t}(\beta_1-\beta_2)} | Z_s \right]$$

$$= \mathbb{E} \left[ e^{-\kappa \theta \psi_{0,1}\gamma(t-s) - Z_0 - \psi_{0,1}\gamma(t-s)} \right]$$

$$= e^{-\kappa \theta \psi_{0,1}\gamma(t-s) - Z_0 - \psi_{0,1}\gamma(s)}$$

(C.9)

Using equation (C.9), one can show that for any $\alpha < 1,$

$$\int_0^t \mathbb{E} \left[ e^{-\frac{1}{t}(\beta_1-\beta_2)} \right] ds = \int_0^{t-\epsilon^\alpha} \mathbb{E} \left[ e^{-\frac{1}{t}(\beta_1-\beta_2)} \right] ds + \int_{t-\epsilon^\alpha}^t \mathbb{E} \left[ e^{-\frac{1}{t}(\beta_1-\beta_2)} \right] ds$$

$$\leq C_7 e^{-C_8 \epsilon^\alpha} + C_9 \epsilon \alpha,$$

for some constants $C_7, C_8$ and $C_9$. This implies that for $\alpha \in (\frac{1}{2}, 1)$ there exists a constant $C_10$ such that

$$\frac{\kappa \theta}{\sigma \sqrt{t}} \int_0^t \mathbb{E} \left[ e^{-\frac{1}{t}(\beta_1-\beta_2)} \right] ds \leq C_10.$$  (C.10)

Having established a uniform bound on the first two terms in (C.6), we turn our attention toward the third and final term. For $\alpha < 1$, we have

$$\frac{1}{\sigma e^{3/2}} \mathbb{E} \left[ \int_0^t e^{-\frac{1}{t}(\beta_1-\beta_2)} \psi_{t} (Z_t, Z_s) ds \right] \leq \frac{1}{\sigma e^{3/2}} \mathbb{E} \left[ \int_0^t e^{-\frac{1}{t}(\beta_1-\beta_2)} \psi_{t} (Z_t, Z_s) | Z_t - Z_s | ds \right],$$

and we decompose the time integral from 0 to $(t - \epsilon^\alpha)$ and from $(t - \epsilon^\alpha)$ to $t$. For the integral from 0 to $(t - \epsilon^\alpha)$ we compute

$$\frac{1}{\sigma e^{3/2}} \mathbb{E} \left[ \int_0^{t-\epsilon^\alpha} e^{-\frac{1}{t}(\beta_1-\beta_2)} \psi_{t} (Z_t, Z_s) | Z_t - Z_s | ds \right]$$

$$\leq \frac{1}{\sigma e^{3/2}} \sqrt{t \mathbb{E} \left[ \left( \sup_{0 \leq s \leq t} | Z_s | \right)^2 \right]} \sqrt{\int_0^{t-\epsilon^\alpha} \mathbb{E} \left[ e^{-\frac{1}{t}(\beta_1-\beta_2)} \right] ds}$$

$$\leq \frac{1}{\epsilon^{3/2}} C_{11} e^{-C_{12} \epsilon^{\alpha - 1}} \leq C_{13},$$  (C.11)

for some constants $C_{11}, C_{12}$ and $C_{13}$. For the integral from $(t - \epsilon^\alpha)$ to $t$, we have

$$\frac{1}{\sigma e^{3/2}} \mathbb{E} \left[ \int_0^{t} e^{-\frac{1}{t}(\beta_1-\beta_2)} \psi_{t} (Z_t, Z_s) | Z_t - Z_s | ds \right]$$

$$\leq \frac{1}{\sigma e^{3/2}} \mathbb{E} \left[ \sup_{t - \epsilon^\alpha \leq s \leq t} | Z_t - Z_s | \int_{t-\epsilon^\alpha}^t e^{-\frac{1}{t}(\beta_1-\beta_2)} Z_s ds \right]$$

$$= \frac{1}{\sigma e^{3/2}} \mathbb{E} \left[ \sup_{t - \epsilon^\alpha \leq s \leq t} | Z_t - Z_s | \epsilon (1 - e^{\frac{1}{t}(\beta_1-\beta_2)} \epsilon^\alpha) \right]$$

$$\leq \frac{1}{\sigma \epsilon^{3/2}} \mathbb{E} \left[ \sup_{t - \epsilon^\alpha \leq s \leq t} | Z_t - Z_s | \right] \leq C_{14} e^{\left( \frac{1}{2} \alpha - \frac{1}{2} \right)} \leq C_{14} \epsilon^{\alpha - 1},$$  (C.12)
for some constant $C_{14}$. Putting (C.8–C.12) together, we obtain that for $\frac{1}{2} < \alpha < 1$, the quantity $A$ defined by (C.5) is bounded by $C \epsilon^{\alpha-1}$ for some constant $C$ which depends on $\alpha$, but not on $\epsilon$. Combining with (C.4) gives the desired bound $E|Y_t| \leq C \epsilon^{\alpha-1}$ where $C$ can be chosen independent of $t$ for $t \leq T$.

**Appendix D. Numerical Computation of Option Prices.** The formulas (4.3) and (4.19) for $P_H(t, x, z)$ and $P_1(t, x, z)$ cannot be evaluated analytically. Therefore, in order for these formulas to be useful, an efficient and reliable numerical integration scheme is needed. Unfortunately, numerical evaluation of the integral in (4.3) is notoriously difficult. And, the double and triple integrals that appear in (4.19) are no easier to compute. In this section, we point out some of the difficulties associated with numerically evaluating these expressions, and show how these difficulties can be addressed. We begin by establishing some notation.

$$P^*(t, x, z) \sim P_H(t, x, z) + \sqrt{\epsilon} P_1(t, x, z),$$

$$= \frac{e^{-rt}}{2\pi} \int_{\mathbb{R}} e^{-ikq} \left( 1 + \sqrt{\epsilon} \left( \kappa \theta \hat{f}_0(\tau, k) + z \hat{f}_1(\tau, k) \right) \right) \hat{G}(\tau, k, z) \hat{h}(k) dk,$$

$$= \frac{e^{-rt}}{2\pi} \left( P_{0,0}(t, x, z) + \kappa \sqrt{\epsilon} P_{1,0}(t, x, z) + z \sqrt{\epsilon} P_{1,1}(t, x, z) \right),$$

where we have defined

$$P_{0,0}(t, x, z) := \int_{\mathbb{R}} e^{-ikq} \hat{G}(\tau, k, z) \hat{h}(k) dk,$$  
(D.1)

$$P_{1,0}(t, x, z) := \int_{\mathbb{R}} e^{-ikq} \hat{f}_0(\tau, k) \hat{G}(\tau, k, z) \hat{h}(k) dk,$$  
(D.2)

$$P_{1,1}(t, x, z) := \int_{\mathbb{R}} e^{-ikq} \hat{f}_1(\tau, k) \hat{G}(\tau, k, z) \hat{h}(k) dk.$$  
(D.3)

As they are written, (D.1), (D.2) and (D.3) are general enough to accommodate any European option. However, in order to make progress, we now specify an option payoff. We will limit ourselves to the case of an European call, which has payoff $h(x) = (x - K)^+$. Extension to other European options is straightforward.

We remind the reader that $\hat{h}(k)$ is the Fourier transform of the option payoff, expressed as a function of $q = r(T-t) + \log(x)$. For the case of the European call, we have:

$$\hat{h}(k) = \int_{\mathbb{R}} e^{ikq} (e^q - K)^+ dq = \frac{K^{1+ik}}{ik - k^2}.$$  
(D.4)

We note that (D.4) will not converge unless the imaginary part of $k$ is greater than one. Thus, we decompose $k$ into its real and imaginary parts, and impose the following condition on the imaginary part of $k$.

$$k = k_r + ik_i,$$

$$k_i > 1.$$  
(D.5)

When we integrate over $k$ in (D.1), (D.2) and (D.3), we hold $k_i > 1$ fixed, and integrate $k_r$ over $\mathbb{R}$. 


Numerical Evaluation of $P_{0,0}(t, x, z)$. We rewrite (D.1) here, explicitly using expressions (4.7) and (D.4), for $\hat{G}(\tau, k, z)$ and $\hat{h}(k)$ respectively.

$$P_{0,0}(t, x, z) = \int_{\mathbb{R}} e^{-ikq} e^{C(\tau, k) + zD(\tau, k)} \frac{K^{1+ik}}{ik - k^2} dk, \quad (D.6)$$

In order for any numerical integration scheme to work, we must verify the continuity of the integrand in (D.6). First, by (D.5), the poles at $k = 0$ and $k = i$ are avoided. The only other worrisome term in the integrand of (D.6) is $e^{C(\tau, k)}$, which may be discontinuous due to the presence of the log in $C(\tau, k)$. We recall that any $\theta \in \mathbb{C}$ can be represented in polar notation as $\zeta = r \exp(i\theta)$, where $\theta \in [-\pi, \pi]$. In this notation, $\log \zeta = \log r + i\theta$. Now, suppose we have a map $\zeta(k_r) : \mathbb{R} \to \mathbb{C}$. We see that whenever $\zeta(k_r)$ crosses the negative real axis, $\log \zeta(k_r)$ will be discontinuous (due to $\theta$ jumping from $-\pi$ to $\pi$ or from $\pi$ to $-\pi$). Thus, in order for $\log \zeta(k_r)$ to be continuous, we must ensure that $\zeta(k_r)$ does not cross the negative real axis.

We now return our attention to $C(\tau, k)$. We note that $C(\tau, k)$ has two algebraically equivalent representations, (4.8), and the following representation:

$$C(\tau, k) = \frac{\kappa \theta}{\sigma^2} ((\kappa + 2ik\sigma - d(k)) \tau - 2 \log \zeta(\tau, k)), \quad (D.7)$$
$$\zeta(\tau, k) := \frac{e^{-\tau d(k) / g(k)} - 1}{1 / g(k) - 1}. \quad (D.8)$$

It turns out that, under most reasonable conditions, $\zeta(\tau, k)$ does not cross the negative real axis [16]. As such, as one integrates over $k_r$, no discontinuities will arise from the log $\zeta(\tau, k)$ which appears in (D.7). Therefore, if we use expression (D.7) when evaluating (D.6), the integrand will be continuous.

Numerical Evaluation of $P_{1,1}(t, x, z)$ and $P_{1,0}(t, x, z)$. The integrands in (D.3) and (D.2) are identical to that of (D.1), except for the additional factor of $\hat{f}_1(\tau, k)$. Using equation (4.21) for $\hat{f}_1(\tau, k)$ we have the following expression for $P_{1,1}(t, x, z)$:

$$P_{1,1}(t, x, z) = \int_{\mathbb{R}} e^{-ikq} \left( \int_{0}^{\tau} b(s, k) e^{A(\tau, k, s)} ds \right) e^{C(\tau, k) + zD(\tau, k)} \frac{K^{1+ik}}{ik - k^2} dk, \quad (D.9)$$

Similarly:

$$P_{1,0}(t, x, z) = \int_{0}^{\tau} \int_{\mathbb{R}} e^{-ikq} b(s, k) e^{A(\tau, k, s) + C(\tau, k) + zD(\tau, k)} \frac{K^{1+ik}}{ik - k^2} dk, ds. \quad (D.10)$$

We already know, from our analysis of $P_{0,0}(t, x, z)$, how to deal with the log in $C(\tau, k)$. It turns out that the log in $A(\tau, k, s)$ can be dealt with in a similar manner. Consider the following representation for $A(\tau, k, s)$, which is algebraically equivalent to
expression (4.22):
\[
A(\tau,k,s) = (\kappa + \rho i k \sigma + d(k)) \left( \frac{1 - g(k)}{d(k)g(k)} \right) \\
\times (d(k)(\tau - s) + \log \zeta(\tau,k) - \log \zeta(s,k)) \\
+ d(k)(\tau - s),
\]
(D.11)

where \( \zeta(\tau,k) \) is defined in (D.8). As expressed in (D.11), \( A(\tau,k,s) \) is, under most reasonable conditions, a continuous function of \( k_r \). Thus, if we use (D.11) when numerically evaluating (D.9) and (D.10), their integrands will be continuous.

**Transforming the Domain of Integration.** Aside from using equations (D.7) and (D.11) for \( C(\tau,k) \) and \( A(\tau,k,s) \), there are a few other tricks we can use to facilitate the numerical evaluation of (D.6), (D.10), and (D.9). Denote by \( I_0(k) \) and \( I_1(k,s) \) the integrands appearing in (D.6), (D.9), and (D.10).

\[
P_{0,0} = \int_{\mathbb{R}} I_0(k)dk_r, \\
P_{1,1} = \int_{0}^{\tau} \int_{\mathbb{R}} I_1(k,s)dk_rds, \\
P_{1,0} = \int_{0}^{\tau} \int_{0}^{t} \int_{\mathbb{R}} I_1(k,s)dk_rdsdt.
\]

First, we note that the real and imaginary parts of \( I_0(k) \) and \( I_1(k,s) \) are even and odd functions of \( k_r \) respectively. As such, instead of integrating in \( k_r \) over \( \mathbb{R} \), we can integrate in \( k_r \) over \( \mathbb{R}^+ \), drop the imaginary part, and multiply the result by 2.

Second, numerically integrating in \( k_r \) over \( \mathbb{R}^+ \) requires that one arbitrarily truncate the integral at some \( k_{\text{cutoff}} \). Rather than doing this, we can make the following variable transformation, suggested by [13]:
\[
k_r = -\log \frac{u}{C_\infty}, \\
C_\infty := \sqrt{\frac{1 - \rho^2}{\sigma}} (z + \kappa \theta \tau).
\]
(D.12)

Then, for some arbitrary \( I(k) \) we have
\[
\int_{0}^{\infty} I(k)dk_r = \int_{0}^{1} I \left( \frac{-\log u}{C_\infty} + ik_i \right) \frac{1}{uC_\infty}du.
\]

Thus, we avoid having to establish a cutoff value, \( k_{\text{cutoff}} \) (and avoid the error that comes along with doing so).

Finally, evaluating (D.10) requires that one integrates over the triangular region paramaterized by \( 0 \leq s \leq t \leq \tau \). Unfortunately, most numerical integration packages only facilitate integration over a rectangular region. We can overcome this difficulty by performing the following transformation of variables:
\[
s = tv, \\
ds = tdv.
\]
Then, for some arbitrary $I(s)$ we have

$$
\int_0^\tau \int_0^t I(s) ds \, dt = \int_0^\tau \int_0^1 I(tv) dv \, dt.
$$

(D.13)

Pulling everything together we obtain:

$$
P_{0,0} = 2Re \int_0^1 \left( -\log u \frac{1}{C_\infty} + ik_1 \right) \frac{1}{uC_\infty} du,
$$

$$
P_{1,1} = 2Re \int_0^\tau \int_0^1 \int_0^1 I_1 \left( -\log u \frac{1}{C_\infty} + ik_1, s \right) \frac{1}{uC_\infty} du \, ds,
$$

$$
P_{1,0} = 2Re \int_0^\tau \int_0^1 \int_0^1 I_1 \left( -\log u \frac{1}{C_\infty} + ik_1, tv \right) \frac{t}{uC_\infty} du \, dv \, dt,
$$

where $C_\infty$ is given by (D.12). These three changes allow one to efficiently and accurately numerically evaluate (D.6), (D.9) and (D.10).

Numerical tests show that for strikes ranging from 0.5 to 1.5 the spot price, and for expirations ranging from 3 months to 3 years, it takes roughly 100 times longer to calculate a volatility surface using the multiscale model than it does to calculate the same surface using the Heston model.

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