

Math

$$\mathrm{d}f = a\,\mathrm{d}x + b\,\mathrm{d}y \rightarrow \left(\frac{\partial a}{\partial y}\right)_x = \left(\frac{\partial b}{\partial x}\right)_y \quad \text{[Maxwell Relations]}$$

$$\left(\frac{\partial x}{\partial y}\right)_z = -\left(\frac{\partial x}{\partial z}\right)_y \left(\frac{\partial z}{\partial y}\right)_x \quad \text{[Triple Product Rule]}$$

$$g = f - \sum I_j X_j; \quad I_j = \frac{\partial f}{\partial X_j} \quad \text{[Legendre Transformations]}$$

$$f(\lambda x) = \lambda^n f(x) \rightarrow f(x) = \tfrac{1}{n} \sum \left(\frac{\partial f}{\partial x_i}\right) x_i \quad \text{[Euler's Theorem]}$$

$$(\delta^n f) = \tfrac{1}{n!} \left(\frac{\partial^n f}{\partial \vec{X}^n}\right) (\delta \vec{X})^n$$

$$\langle G \rangle = \sum P_\nu G_\nu \quad \text{[Expected Value]}$$

$$I_s = \int_0^\infty f_{F.D.}(\epsilon) g(\epsilon) \mathrm{d}\epsilon \quad \text{[Sommerfeld]}$$
$$= \int_0^\mu g(\epsilon) \mathrm{d}\epsilon + \sum a_{2i} (k_B T)^{2i} \mathrm{d}_\epsilon^{2i-1} g(\epsilon)$$

$$a_{2i} = 2(1 - 2^{1-2i}) \zeta(2i)$$

$$\sum \exp(-\beta(\epsilon - \mu))^i = (1 - \exp[-\beta(\epsilon - \mu)])^{-1} \quad \text{[Sum } |x| < 1]$$

$$\zeta(s) = \sum s^{-i} = \frac{1}{\Gamma(s)} \int_0^\infty \frac{x^{s-1}}{e^x - 1} \mathrm{d}x \quad \text{[Riemann-Zeta Function]}$$

$$\Gamma(s) = \int_0^\infty x^{s-1} e^{-x} \mathrm{d}x \quad \text{[Gamma Function]}$$

Laws of Thermodynamics

$$\mathrm{d}E = \mathrm{d}Q + \mathrm{d}W \quad \text{[First Law]}$$

$$(\Delta S)_{adiabatic} \geq 0 \quad \text{[Second Law]}$$

Thermodynamic Variables

$$\{T\,\mathrm{d}S, -p\,\mathrm{d}V, \mu\,\mathrm{d}N, f\,\mathrm{d}L, E\,\mathrm{d}q, M\,\mathrm{d}H, \gamma\,\mathrm{d}A\} \quad \text{[Conjugate Pairs]}$$

$$T \equiv \left(\frac{\partial E}{\partial S}\right)_{\vec{X}} \geq 0 \quad \text{[Temperature]}$$

$$C_X \equiv T \left(\frac{\partial S}{\partial T}\right)_X \quad \text{[Heat Capacity]}$$

$$K_X \equiv -\tfrac{1}{V} \left(\frac{\partial V}{\partial p}\right)_{X,N} \quad \text{[Compressibility]}$$

Thermodynamic Potentials

$$E = T S - p V + \sum \mu_i N_i \quad \text{[Internal Energy]}$$

$$\mathrm{d}E = T\,\mathrm{d}S - p\,\mathrm{d}V + \sum \mu_i\,\mathrm{d}N_i$$

$$0 = S\,\mathrm{d}T - V\,\mathrm{d}p + \sum_i N_i\,\mathrm{d}\mu_i \quad \text{[Gibbs-Duhem Equation]}$$

$$A = E - T S \quad \text{[Helmholtz Free Energy]}$$

$$\mathrm{d}A = -S\,\mathrm{d}T - p\,\mathrm{d}V + \sum \mu_i\,\mathrm{d}N_i$$

$$H = E + p V \quad \text{[Enthalpy]}$$

$$\mathrm{d}H = T\,\mathrm{d}S + V\,\mathrm{d}p + \sum \mu_i\,\mathrm{d}N_i$$

$$G = E + p V - T S \quad \text{[Gibbs Free Energy]}$$

$$\mathrm{d}G = -S\,\mathrm{d}T + V\,\mathrm{d}p + \sum \mu_i\,\mathrm{d}N_i$$

Ensembles

$$Q(N, V, T) = \sum_E \exp(S/k_B) \exp(-\beta E) = \exp(-\beta A)$$

$$\Xi(\mu, V, T) = \sum_N Q(N, V, T) \exp(\beta \mu N) = \exp(\beta p V)$$

$$\Delta(N, p, T) = \sum_V Q(N, V, T) \exp(-\beta p V) = \exp(-\beta G)$$

$$\phi(V, E, \beta \mu) = \sum_N \Omega(N, V, E) \exp(\beta \mu N) = \exp(\beta H)$$

$$\Psi(V, T, \mu_1, N_2) = \sum_{N_1} Q(N_1, N_2, T, V) \exp(\beta \mu_1 N_1)$$
$$= \exp(\beta p V) \exp(-\beta \mu_2 N_2)$$

$$W(p, \gamma, T, N) = \sum_{V,A} Q(N, V, A, T) \exp(-\beta p V) \exp(\beta \gamma A)$$
$$= \exp(-\beta G)$$

Equilibriums

$$(\delta E)_{S, \vec{X}} \geq 0; \quad (\Delta E)_{S, \vec{X}} > 0 \quad \text{[Equilibrium]}$$

$$(\delta^2 E)_{S, \vec{X}} \geq 0; 0 \leq \left(\frac{\partial I_i}{\partial X_i}\right)_{X_i \neq j} \quad \text{[Stability]}$$

$$0 \leq \delta E = \sum_{j \neq i} (I^{(j)} - I^{(i)}) \delta X^{(j)} \quad \text{[Equality of Intensives]}$$

Entropy

$$\Omega = \frac{N!}{(N-m)!m!} \quad \text{[Number of Microstates]}$$

$$S = k_b \log(\Omega) \quad \text{[Entropy]}$$

$$S = -k_b \sum P_\nu \log(P_\nu) \quad \text{[Gibbs Entropy Formula]}$$

$$P_\nu = \tfrac{1}{\Omega}$$

$$P_\nu = Q^{-1} \exp[-\beta E_\nu]$$

Phase Space

$$\frac{\mathrm{d}p}{\mathrm{d}T} = \frac{\Delta s(T)}{\Delta v(t)} \quad \text{[Clausius-Clapeyron Equation]}$$

$$f = 2 + r - \nu \quad \text{[Gibbs Phase Rule]}$$

$$\frac{\partial p}{\partial V} = 0 = \frac{\partial^2 p}{\partial V^2} \quad \text{[Critical Point]}$$

$$I^{(1)} = I^{(2)} \quad \text{[Equality of Intensive Variables Across Transition]}$$

$$\delta \mu = 0 \text{ \& \;} \frac{\partial \mu}{\partial p} = v \quad \text{[Equal Area]}$$

Ideal Gas

$$Q = \tfrac{1}{N!} q^N \quad \text{[Partition Function]}$$

$$q_{mona} = q_{trans} q_{el} q_{nuc} \quad \text{[Single Monatomic Partition]}$$

$$q_{dia} = q_{mona} q_{rot} q_{vib} \quad \text{[Single Diatomic Partition]}$$

$$q_{trans} = \frac{V}{\Lambda^3}; \text{ where } \Lambda = \sqrt{\frac{m}{2\pi\beta\hbar^2}} \quad \text{[Translational]}$$

$$q_{el} = \exp[-\beta \epsilon_0] (g_0 + \sum g_i \exp[-\beta(\epsilon_i - \epsilon_0)]) \quad \text{[Electron]}$$

$$\text{(N.B. large spacing makes higher terms ignorable)}$$

$$q_{nuc} = (2S_{nuc} + 1) \exp[-\beta \epsilon_{0,nuc}] \quad \text{[Nucleus]}$$

$$\text{(N.B. } \epsilon_{0,nuc} = 0 \text{ is a general good approximation)}$$

$$q_{rot} = \frac{T}{\Theta_{rot}}; \text{ where } \Theta_{rot} = \frac{\hbar^2}{2Ik_B} \quad \text{[Rotational]}$$

$$\text{(N.B. Only when } \Theta_{rot} \ll T \text{ \& \;} I = \mu r^2)$$

$$q_{vib} = [2 \sinh(\beta \hbar \omega_{vib}/2)]^{-1} \quad \text{[Vibrational]}$$

$$\text{(N.B. If } \hbar \omega_{vib} \ll k_B T \text{ then equal to } (\beta \hbar \omega_{vib})^{-1})$$

Chemical Equilibrium

$$\text{For } [\nu_A A + \nu_B B \rightarrow \nu_C C + \nu_D D]$$

$$K_c(T) = \frac{\rho_C^{\nu_C} \rho_D^{\nu_D}}{\rho_A^{\nu_A} \rho_B^{\nu_B}} = \frac{(q_C/V)^{\nu_C} (q_D/V)^{\nu_D}}{(q_A/V)^{\nu_A} (q_B/V)^{\nu_B}}$$

$$K_p(T) = \frac{p_C^{\nu_C} p_D^{\nu_D}}{p_A^{\nu_A} p_B^{\nu_B}} = (k_B T)^{\nu_C + \nu_D - \nu_A - \nu_B} K_c(T)$$

Fermion/Boson Statistics

$$f(\epsilon)_{F.D./B.E.} = [\exp(\beta(\epsilon - \mu)) \pm 1]^{-1} \quad \text{[Statistics]}$$

$$\Xi_{F.D./B.E.} = \Pi[1 \pm \exp(-\beta(\epsilon_i - \mu))]^{\pm 1} \quad \text{[Grand Canonical]}$$

Density of State

$$D(\epsilon) = \frac{g}{\gamma} S_{d-1} \left(\frac{2\pi L}{\hbar \alpha^{1/\gamma}}\right)^d \epsilon^{d/\gamma-1} \quad \text{[d-space D.o.S. of } \epsilon = \alpha|p|^\gamma]$$

$$\langle E \rangle = \int \epsilon D(\epsilon) f(\epsilon) \mathrm{d}\epsilon \quad \text{[Expected Energy]}$$

$$\langle N \rangle = \int D(\epsilon) f(\epsilon) \mathrm{d}\epsilon \quad \text{[Expected Number]}$$

Virial Expansion

$$Q = \tfrac{1}{h^{3N}} \int \dots \int \exp[-\beta H] \mathrm{d}\mathbf{p}_1 \dots \mathrm{d}\mathbf{p}_N \mathrm{d}\mathbf{r}_1 \dots \mathrm{d}\mathbf{r}_N$$
$$= \tfrac{1}{\Lambda^{3N} N!} Z_N \quad \text{[Classical Partition Function]}$$

$$Z_N = \int \dots \int \exp[-\beta U_N] \mathrm{d}\mathbf{r}_1 \dots \mathrm{d}\mathbf{r}_N \quad \text{[Configuration Function]}$$

$$\Xi = \sum \frac{Z_N}{N!} s^N$$

$$s = \exp[\beta \mu] / \Lambda^3 \quad \text{[Activity]}$$

$$p = kT \sum b_i s^i \quad \text{[Activity Expansion]}$$

$$\beta p = \sum B_n(T) \rho^n \quad \text{[Virial Expansion]}$$

$$f_{ij} = \exp[-\beta u(r_{ij})] - 1 \quad \text{[Meyer's } f\text{-function]}$$

$$b_2 = \tfrac{1}{2!V} (Z_2 - Z_1^2) = \tfrac{1}{2!V} \int \mathrm{d}\mathbf{r}_1 \int \mathrm{d}\mathbf{r}_2 f_{12} = \tfrac{1}{2!V} [\cdot - \cdot]$$

$$b_3 = \tfrac{1}{3!V} (Z_3 - 3Z_2 Z_1 + 2Z_1^3)$$
$$= \tfrac{1}{3!V} \int \mathrm{d}\mathbf{r}_1 \int \mathrm{d}\mathbf{r}_2 \int \mathrm{d}\mathbf{r}_3 [f_{12} f_{13} f_{23} + 3f_{12} f_{23}]$$
$$= \tfrac{1}{3!V} (\Delta + 3\Lambda) = \tfrac{1}{6V} \Delta - \tfrac{1}{2V^2} [\cdot - \cdot]^2; \text{ since } \Lambda = V^{-1} [\cdot - \cdot]^2$$

$$B_n(T) = -\tfrac{(n-1)}{n!V} S'_n; \quad s'_{i,2} = -\tfrac{s_{i,1,3} = \Delta}{3\Box + 6\Diamond + \Box} \quad \text{[Virial Coefficients]}$$

$$p = \frac{Nk_B T}{V} - \frac{1}{3V} \left\langle \sum_{i < j} r_{ij} \frac{\partial}{\partial r_{ij}} u(r_{ij}) \right\rangle \quad \text{[Virial Pressure]}$$

Reduced Distribution Functions

$$\rho^{(n/N)} = \frac{N!}{(N-n)!} \int \mathrm{d}\mathbf{r}^{N-n} \frac{\exp[-\beta U(r^N)]}{\int \mathrm{d}\mathbf{r}^N \exp[-\beta U(r^N)]}$$

$$g(r) = \rho^{(2/N)} / \rho^2 \quad \text{[Pair Correlation Function]}$$

$$\text{(N.B. For ideal gas } g(r) = (1 - N^{-1}))$$

$$\left\langle \sum_{i < j} u(r_{ij}) \right\rangle = \frac{N(N-1)}{2} \frac{1}{Z_N} \int u(r_{12}) \exp[-\beta U] \mathrm{d}\mathbf{r}$$

$$= \frac{N}{2} \int \rho g(r) u(r) \mathrm{d}\mathbf{r}$$

$$\langle E \rangle = \frac{3Nk_B T}{2} + \frac{N\rho}{2} \int g(r) u(r) \mathrm{d}\mathbf{r}$$

$$\left\langle \sum u^{(3)}(r_{ij}, r_{jk}) \right\rangle = \frac{N\rho^2}{6} \int \mathrm{d}\mathbf{r}_{12} \int \mathrm{d}\mathbf{r}_{23} g^{(3)} u^{(3)}$$

$$\frac{\beta p}{\rho} = 1 - \frac{\beta \rho}{6} \int r g(r) \frac{\partial u(r)}{\partial r} \mathrm{d}\mathbf{r} \quad \text{[Virial Pressure]}$$

Random Equations

$$\left\langle \Gamma_i \frac{\partial H}{\partial \Gamma_j} \right\rangle = \delta_{ij} k_B T \quad \text{[Generalized Equipartition]}$$

$$\langle (\delta N)^2 \rangle = \langle N^2 \rangle - \langle N \rangle^2 = \frac{\partial^2 \log \Xi}{\partial (\beta \mu)^2}$$

Random Identities

$$\langle E \rangle = - \left(\frac{\partial \log(Q)}{\partial \beta} \right)_{N,V}$$

$$\left(\frac{\partial \beta A}{\partial \beta} \right)_{N,V} = E$$

$$\langle (\delta^2 E) \rangle = k_B T^2 C_v$$

$$\left(\frac{\partial \log(Q)}{\partial V} \right)_{T,N} = - \left(\frac{\partial \beta A}{\partial V} \right)_{T,N} = \beta p$$

$$\left(\frac{\partial \log(Q)}{\partial (-\beta)} \right)_{V,N} = \left(\frac{\partial (-\beta A)}{\partial (-\beta)} \right)_{V,N} = \langle E \rangle$$

$$\frac{\partial \mu}{\partial v} = v \frac{\partial p}{\partial v}$$

$$\beta = \left(\frac{\partial \log(\Omega)}{\partial E} \right)_{N,V}$$

