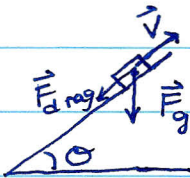


# Phys 105A Solution #2

(1)



The component of the gravitational force in the direction of the velocity (along the slope) is  $-mg \sin \theta$ . The drag force is  $\vec{F}_{\text{drag}} = -b\vec{v}$ . So

$$m \frac{dv}{dt} = -mg \sin \theta - bv \quad (*)$$

$$\Rightarrow \int_{v_0}^v \frac{d\tilde{v}}{g \sin \theta + b\tilde{v}/m} = - \int_0^t d\tilde{t} = -t$$

$$\frac{m}{b} \ln \left( g \sin \theta + \frac{b\tilde{v}}{m} \right) \Big|_{v_0}^v$$

Exponentiate:  $\frac{g \sin \theta + bv/m}{g \sin \theta + bv_0/m} = e^{-bt/m}$

Solving for  $v$ :

$$v(t) = \frac{mg}{b} \sin \theta (e^{-bt/m} - 1) + v_0 e^{-bt/m}$$

Note that we solved the equation of motion (\*) using separation of variables. You can also do it by defining  $u = v + \frac{mg}{b} \sin \theta$  and solving a simpler equation for  $u$ .

$$(2.8) (2) \quad m \frac{dv}{dt} = F(v) = -c v^{3/2}$$

$$\text{So } \frac{dv}{v^{3/2}} = -\frac{c}{m} dt$$

Integrating both sides:

$$\frac{c}{m} t = -\int_{v_0}^v \frac{dv}{v^{3/2}} = \frac{2}{v^{1/2}} \Big|_{v_0}^v = \frac{2}{\sqrt{v}} - \frac{2}{\sqrt{v_0}}$$

$$\text{Solving for } v(t): \quad \frac{4}{v} = \left( \frac{ct}{m} + \frac{2}{\sqrt{v_0}} \right)^2$$

$$\Rightarrow v(t) = \left( \frac{ct}{2m} + \frac{1}{\sqrt{v_0}} \right)^{-2}$$

It takes an  $\infty$  time for the mass to come to rest.

$$(2.11) \quad 3. (a) \quad m \frac{dv_y}{dt} = -mg - bv \quad \text{Set } v_{ter} = \frac{mg}{b}$$

This was solved in class (and is very similar to the solution 2.29-2.31 in the text).

The solution is:

$$v_y(t) = (v_0 + v_{ter}) e^{-bt/m} - v_{ter}$$

We get the position by integrating & setting  $y(0)=0$ :

$$y(t) = \frac{m}{b} (v_0 + v_{ter}) (1 - e^{-bt/m}) - v_{ter} t$$

(b) At the highest point,  $v_y=0$ . This occurs at a time  $\frac{v_0 + v_{ter}}{v_{ter}} = e^{bt/m}$

So

$$t_{top} = \frac{m}{b} \ln \left( 1 + \frac{v_0}{v_{ter}} \right)$$

and

$y_{max} = y(t_{top})$ . Using  $v_y(t_{top}) = 0$ , we have

$$y_{max} = \frac{m}{b} \left[ v_0 - v_{ter} \ln \left( 1 + \frac{v_0}{v_{ter}} \right) \right]$$

(c) When  $v_0/v_{ter} \ll 1$ , we can expand

$$\ln \left( 1 + \frac{v_0}{v_{ter}} \right) \approx \frac{v_0}{v_{ter}} - \frac{1}{2} \left( \frac{v_0}{v_{ter}} \right)^2$$

$$\text{So } y_{max} = \frac{m}{b} \left[ v_0 - v_0 + \frac{1}{2} \frac{v_0^2}{v_{ter}} \right] = \frac{v_0^2}{2g}$$

(2.31) 4.(a) For quadratic air resistance  $v_{ter} = \sqrt{\frac{mg}{c}}$

where  $c = \gamma D^2$ ,  $\gamma = .25 \text{ N s}^2/\text{m}^4$  (eq. 2.6).

Plugging in  $m = .6 \text{ Kg}$ ,  $g = 9.8 \text{ m/sec}^2$ ,  $D = .24 \text{ m}$   
we get

$$v_{ter} = 20.2 \text{ m/s}$$

(b) From eq. (2.58):  $y = \frac{v_{ter}^2}{g} \ln \left[ \cosh \left( \frac{gt}{v_{ter}} \right) \right]$

$$\text{So } \cosh \left( \frac{gt}{v_{ter}} \right) = e^{gy/v_{ter}^2}$$

$$t = \frac{v_{ter}}{g} \operatorname{arccosh} \left( e^{gy/v_{ter}^2} \right)$$

Setting  $y = 30 \text{ m}$ , and using the above value  
for  $v_{ter}$  yields

$$t = 2.78 \text{ sec}$$

From eq. (2.57):  $v = v_{ter} \tanh \left( \frac{gt}{v_{ter}} \right)$

Plugging in the values for  $g, t, v_{ter}$  yields

$$v = 17.7 \text{ m/sec}$$

In vacuum,  $t = \left( \frac{2y}{g} \right)^{1/2} = 2.47 \text{ sec}$  ~~in vac~~

and  $v = gt = (2gy)^{1/2} = 24.2 \text{ m/sec}$

(2.40) 5.

$$m \frac{dv}{dt} = -bv - cv^2$$

$$\frac{m dv}{bv + cv^2} = -dt \Rightarrow \int_{v_0}^v \frac{dv}{bv + cv^2} = -\frac{t}{m}$$

$$\Rightarrow \frac{1}{b} \ln \frac{v}{b + cv} \Big|_{v_0}^v = -\frac{t}{m}$$

$$\Rightarrow \ln \frac{v}{b + cv} - \ln \frac{v_0}{b + cv_0} = -\frac{bt}{m}$$

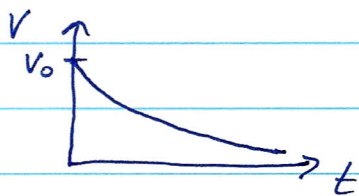
$$\Rightarrow \frac{v}{b + cv} = \frac{v_0}{b + cv_0} e^{-bt/m} = a e^{-bt/m}$$

where we defined  $a = \frac{v_0}{b + cv_0}$ . Now solve for  $v(t)$ :

$$v = (b + cv) a e^{-bt/m}$$

$$v(1 - cae^{-bt/m}) = ba e^{-bt/m}$$

$$\Rightarrow v(t) = \frac{ba e^{-bt/m}}{1 - cae^{-bt/m}}$$



$v(t)$  decays exponentially at late time. The linear term dominates since  $v$  is small.