

Curie Susceptibility

- Magnetic moment in general is proportional to spin

$$\mu = g\mu_B \mathbf{S} / \hbar$$

spin S quantum spin
 $S^2 = S(S+1)\hbar^2$

g-factor
(could be a tensor)
 $g \approx 2$ for pure spin
moment

Bohr magneton
 $\mu_B = \frac{e\hbar}{2m_e} = 9.3 \times 10^{-24} J/T$
 $= 0.671 K/T$

- Magnetic dipole interaction

$$\mathcal{H} = -\mu \cdot \mathbf{H}$$

Curie Susceptibility

- Thermodynamics $M = -\frac{\partial F}{\partial H}$ $F = -kT \ln Z$

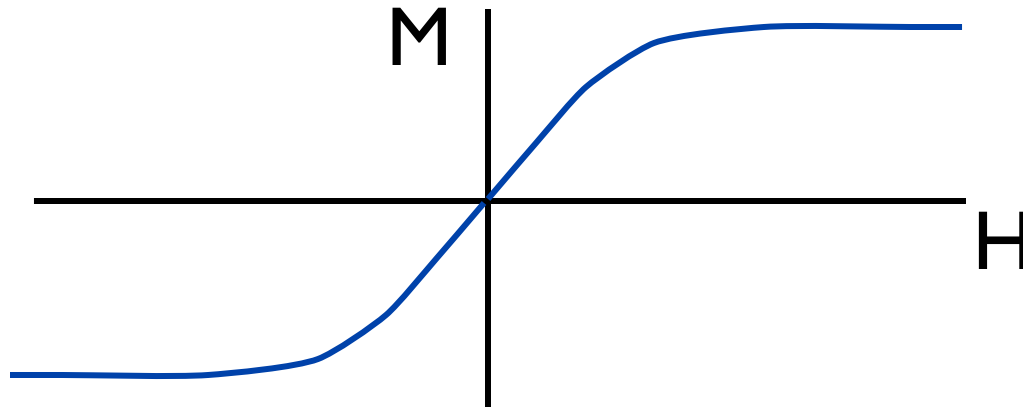
$$Z = \sum_n e^{-E_n/kT}$$

- For example, $S=1/2$ $E_{\uparrow/\downarrow} = \mp g\mu_B H/2$

$$Z = 2 \cosh \left(\frac{g\mu_B H}{2kT} \right) \quad M = \frac{g\mu_B N}{2} \tanh \left(\frac{g\mu_B H}{2kT} \right)$$

Curie Susceptibility

- Magnetization $M = \frac{g\mu_B N}{2} \tanh\left(\frac{g\mu_B H}{2kT}\right)$



- Susceptibility

$$\chi = \frac{(g\mu_B)^2 N}{4k_B T} = \frac{A}{T}$$

Curie Law:
indicates local moment

Curie Susceptibility

- Thermodynamics $M = -\frac{\partial F}{\partial H}$ $F = -kT \ln Z$

$$Z = \sum_n e^{-E_n/kT}$$

- General result

$$\chi = \frac{N(g\mu_B)^2}{3} \frac{S(S+1)}{kT}$$

plot χ^{-1} versus T to extract “effective moment”

Contrast with metals

Curie Law

$$\chi = \frac{N(g\mu_B)^2}{3} \frac{S(S+1)}{kT}$$

Pauli paramagnetism

$$\chi = V \frac{(g\mu_B)^2}{4} D(\epsilon_F)$$



basically $kT \rightarrow \epsilon_F$

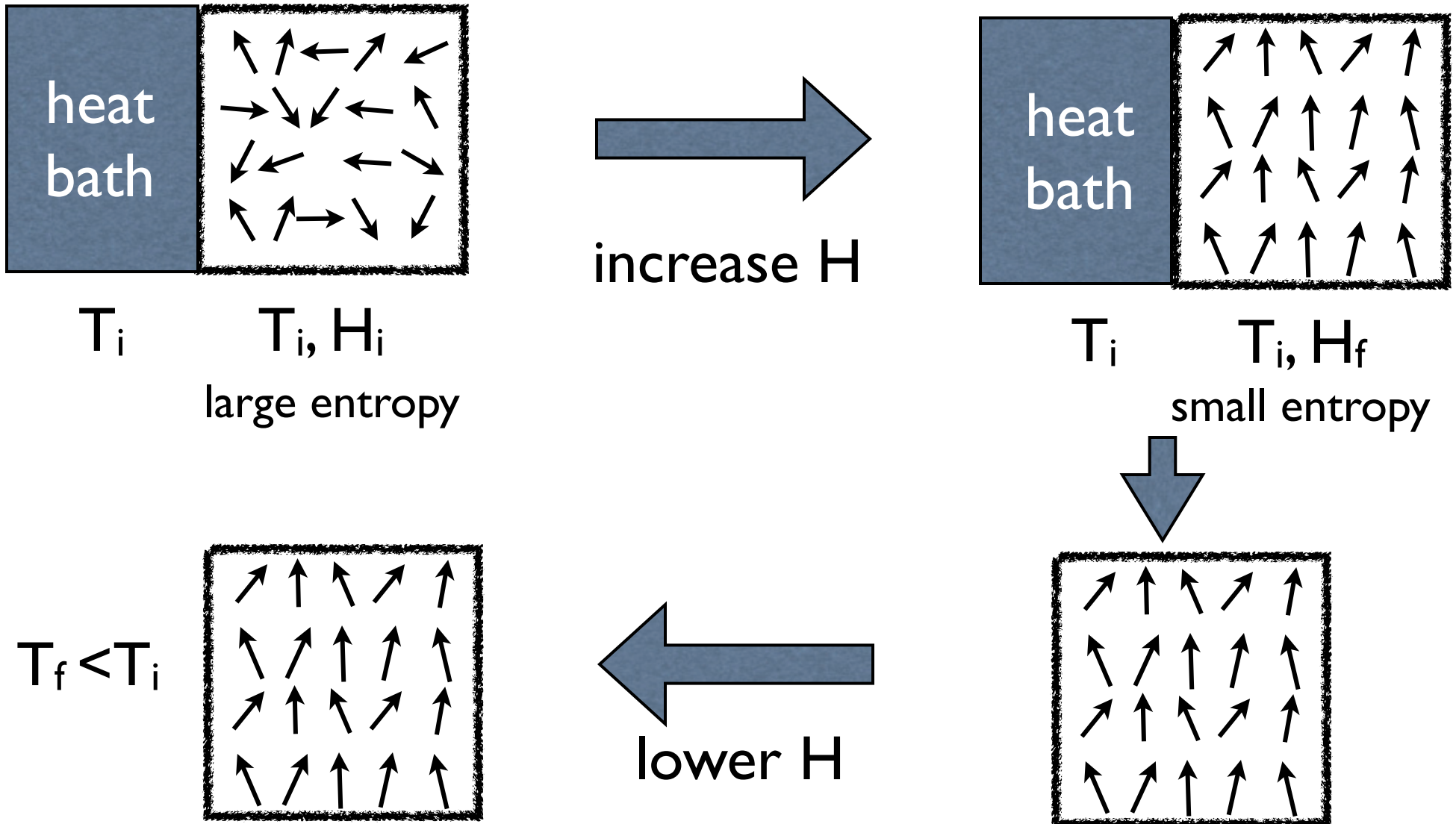
Does χ really diverge for local moments?

NO, because they interact

Magnetic cooling

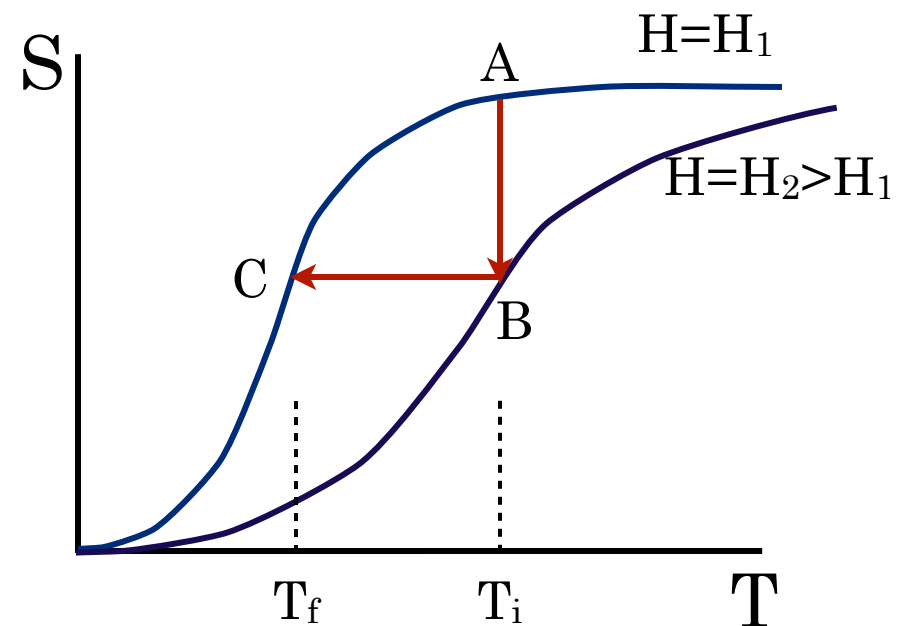
- The large susceptibility of free spins at low temperature means they are easily aligned by small magnetic fields
- This alignment corresponds to a drastic reduction of entropy. One can use this control over entropy to remove entropy from another system, thereby cooling it.

Magnetic Cooling



Magnetic Cooling

- A → B: isothermal step - raise field, lower entropy
- B → C: adiabatic step - lower field, same entropy: lower temperature
- For paramagnetic spins, $S = S(H/T)$
 - Hence $H_1/T_f = H_2/T_i$



Exchange

- How do spins interact?
- Magnetostatic dipole-dipole coupling

$$H_{d-d} = -\frac{\mu_0}{4\pi r^3} [3(\mathbf{m} \cdot \mathbf{r})(\mathbf{m}' \cdot \mathbf{r}) - \mathbf{m} \cdot \mathbf{m}']$$

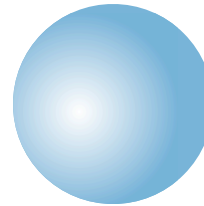
- This is rather weak, \approx 1K for even large spins
- Electrostatic interaction usually dominates, just as it does inside atoms
- Indirectly leads to spin coupling through Pauli principle

Exchange

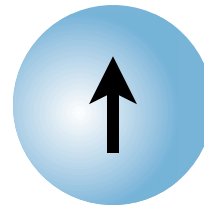
- Generally think of exchange interactions between spins as due to hopping of electrons from one orbital to another
- Many types of exchange, complicated by varieties of orbitals involved, including those on non-magnetic atoms
- We will just discuss the simplest case

Direct exchange

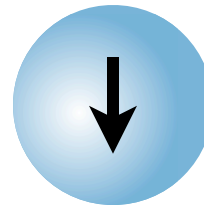
Assume individual
 $S=1/2$ ions



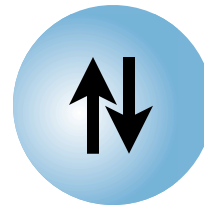
$$E=0$$



$$E=\varepsilon$$



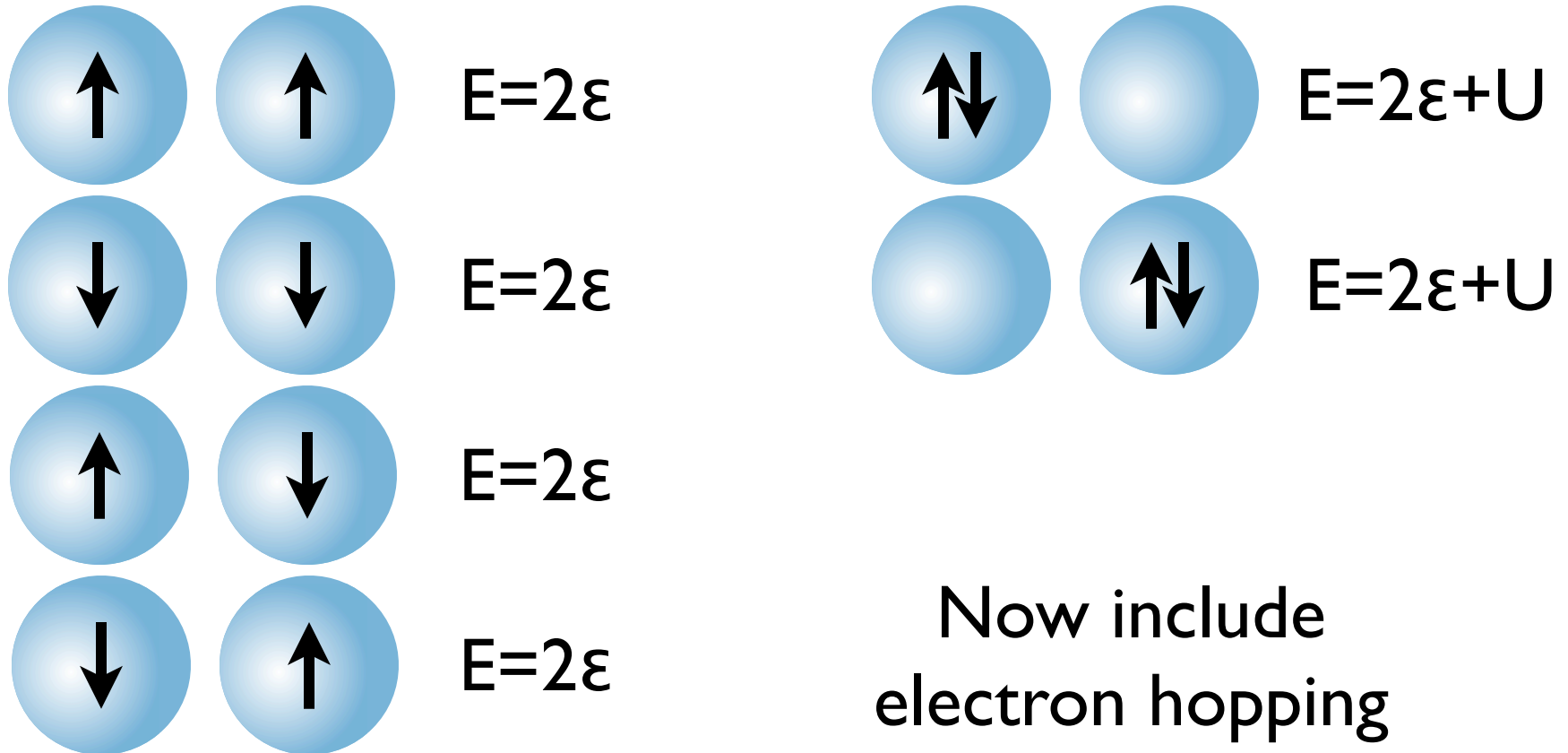
$$E=\varepsilon$$



$$E=2\varepsilon+U$$

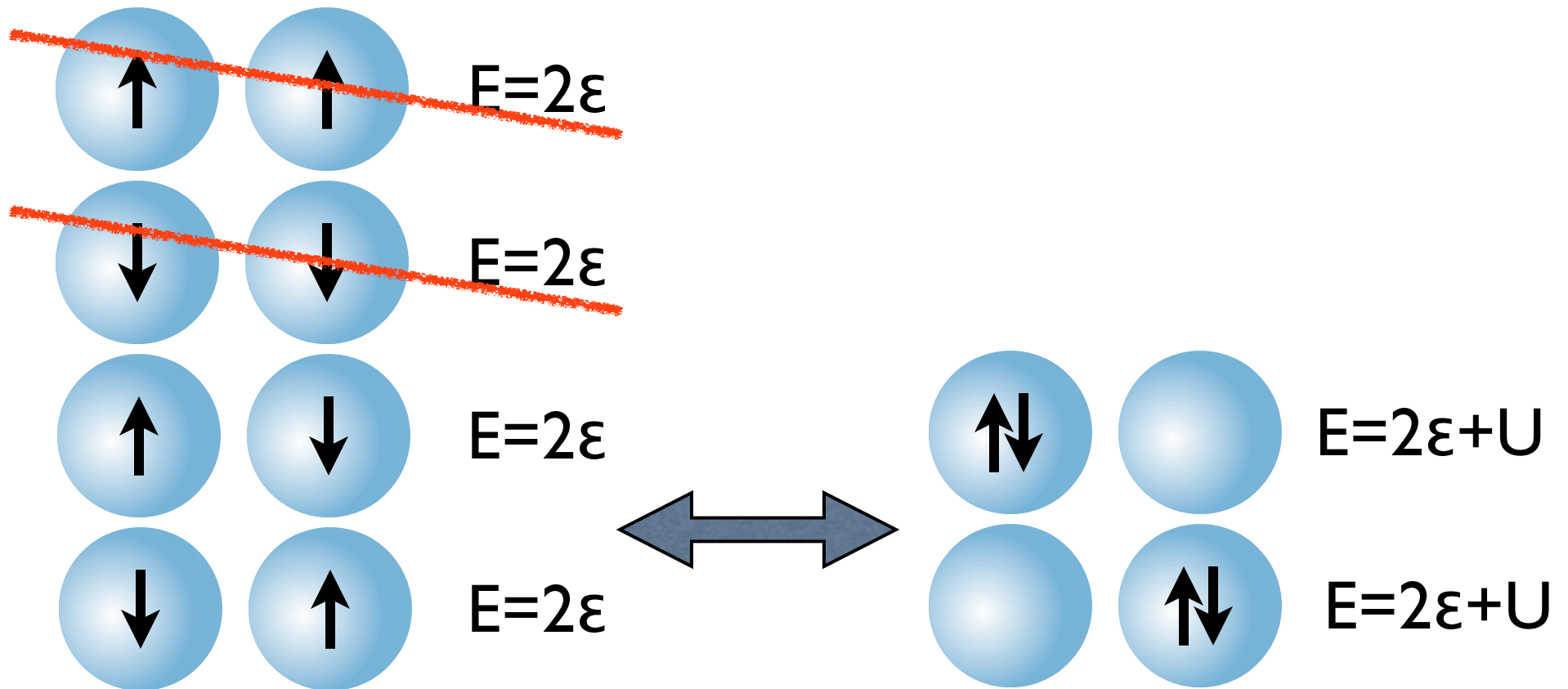
Direct exchange

- Two ions, 2 electrons. With no contact:



Direct exchange

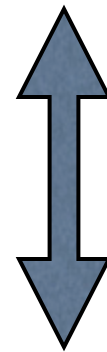
- Hopping? Prohibited by Pauli for parallel spins ($S=1$)



Direct exchange

- Since spin is conserved, only $S=0$ state can hop

$$\frac{1}{\sqrt{2}} \left(\begin{array}{c} \uparrow \\ \downarrow \end{array} \begin{array}{c} \downarrow \\ \uparrow \end{array} - \begin{array}{c} \downarrow \\ \uparrow \end{array} \begin{array}{c} \uparrow \\ \downarrow \end{array} \right) \quad \text{parity } P=-1$$

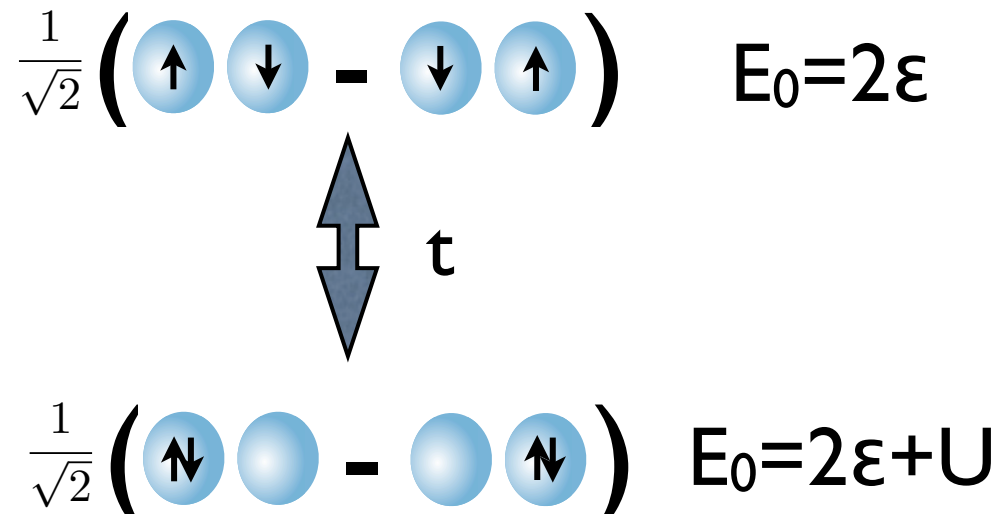


amplitude = t

$$\frac{1}{\sqrt{2}} \left(\begin{array}{c} \uparrow\downarrow \\ \bullet \end{array} \begin{array}{c} \bullet \\ \uparrow\downarrow \end{array} - \begin{array}{c} \bullet \\ \uparrow\downarrow \end{array} \begin{array}{c} \uparrow\downarrow \\ \bullet \end{array} \right) \quad \text{parity } P=-1$$

Direct exchange

- Since spin is conserved, only $S=0$ state can hop



$$H = \begin{pmatrix} 2\varepsilon & t \\ t & 2\varepsilon + U \end{pmatrix}$$

$$E_- \approx 2\varepsilon - \frac{t^2}{U}$$

singlet energy
lowered by $O(t^2/U)$

Direct exchange

- Write as an “effective Hamiltonian”:

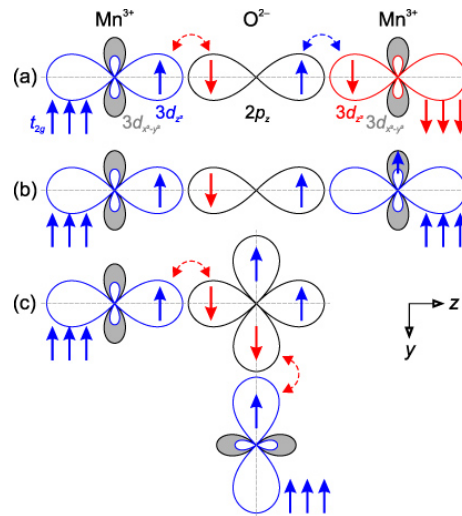
$$\begin{aligned} H_{\text{eff}} &= -\frac{t^2}{U} \hat{P}_{S=0} \\ &= -\frac{t^2}{U} \left[1 - \frac{1}{2} (\mathbf{S}_i + \mathbf{S}_j)^2 \right] \\ &= \text{const.} + J \mathbf{S}_i \cdot \mathbf{S}_j \end{aligned}$$

- With $J \sim t^2/U$. Note this is typically *antiferromagnetic* interaction: favors singlet/anti-parallel spins

J can be as large as 1000K, or as small as a few K

Other exchanges

- Super-exchange: exchange due to electrons hopping through an intermediate orbital



usually antiferromagnetic

this can be ferromagnetic

Other exchanges

- RKKY exchange:
 - interaction of spins mediated by delocalized metallic electrons
 - only in a metal, obviously
- Double exchange
 - more exotic exchange in some metals with some localized electrons