

Magnetic order

- In a crystal with a periodic lattice of spins, exchange interactions typically induce an *ordered* state at low temperature
- For example, the ferromagnetic Heisenberg model:

$$H = -J \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j \quad J > 0$$

- Wants every spin parallel to its neighbor, so they choose a global axis

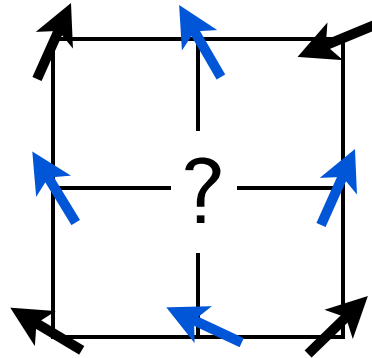
$$\langle \mathbf{S}_i \rangle = \mathbf{m}$$

Mean field theory

- When $kT \gtrsim J$, spins will fluctuate thermally, and \mathbf{m} will be reduced.
- We can study this with *mean field theory*

$$\begin{aligned} H &= -J \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j \\ &\rightarrow -J \sum_{\langle ij \rangle} [\langle \mathbf{S}_i \rangle \cdot \mathbf{S}_j + \mathbf{S}_i \cdot \langle \mathbf{S}_j \rangle - \langle \mathbf{S}_i \rangle \cdot \langle \mathbf{S}_j \rangle] \\ &= -zJ \sum_i \mathbf{m} \cdot \mathbf{S}_i + \text{const.} \end{aligned}$$

Mean field theory



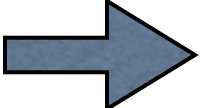
$$g\mu_B \mathbf{H}_{\text{eff}} = J \langle \uparrow + \uparrow + \downarrow + \uparrow \rangle$$

- This reduces the problem to independent spins in an *effective* “exchange field”
- Note: this exchange field can be a thousand times larger than physical laboratory fields!

Mean field theory

- Define $h = z J m$ ($= g \mu_B H_{\text{eff}}$)
- Then we know for a single spin

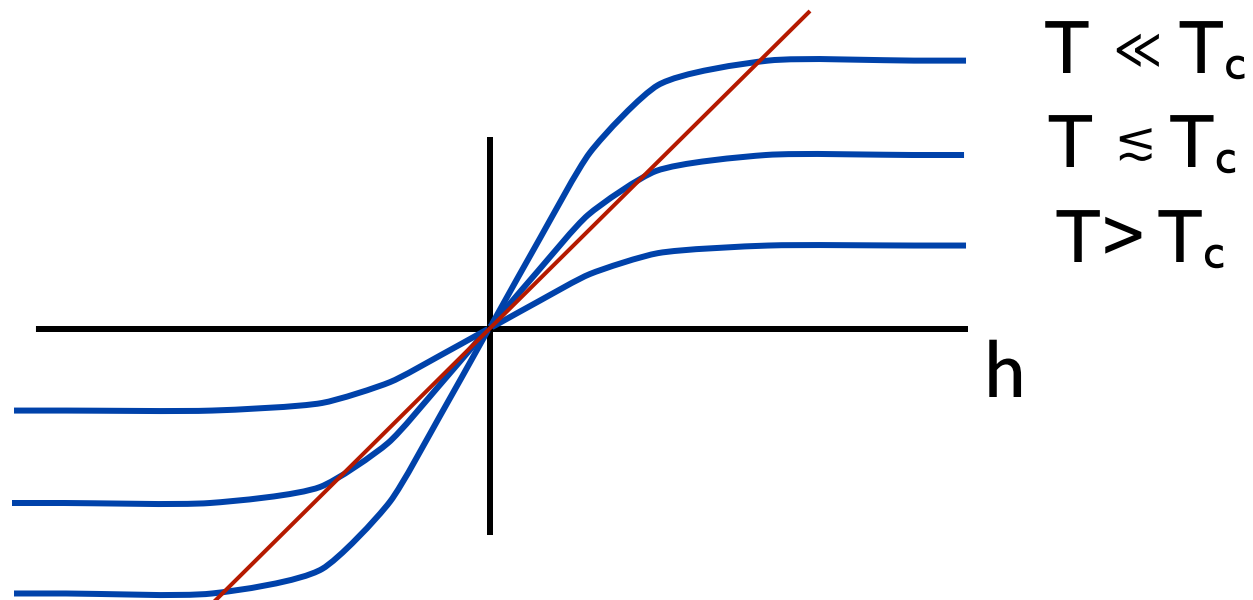
$$|\langle \mathbf{S}_i \rangle| = m = S B_S(\beta h S)$$

 $m = S B_S(\beta z J S m)$

- For example for $S=1/2$

$$m = \frac{1}{2} \tanh \left[\frac{z J m}{2kT} \right]$$

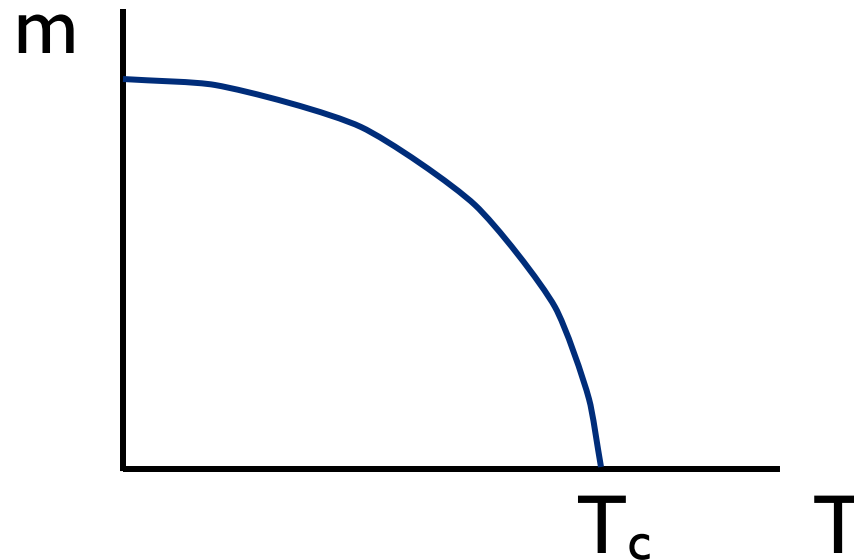
Mean field theory



- Non-zero solution for m appears only for $T < T_c$
- equality of slopes implies $kT_c = zJ/4$ ($= \frac{zJS(S+1)}{3}$)

Mean field theory

- *Zero field* magnetization:



- T_c is called the “Curie point” or critical temperature

Susceptibility

- We may guess that the susceptibility gets large on approaching the Curie point, since the material almost forms a magnetization with no field at all.
- This is indeed true.
- Within MFT, just shift $h \rightarrow h + g \mu_B H$

Susceptibility

$$m = \frac{1}{2} \tanh \left[\frac{zJm + g\mu_B H}{2kT} \right] \approx \frac{zJm + g\mu_B H}{4kT}$$

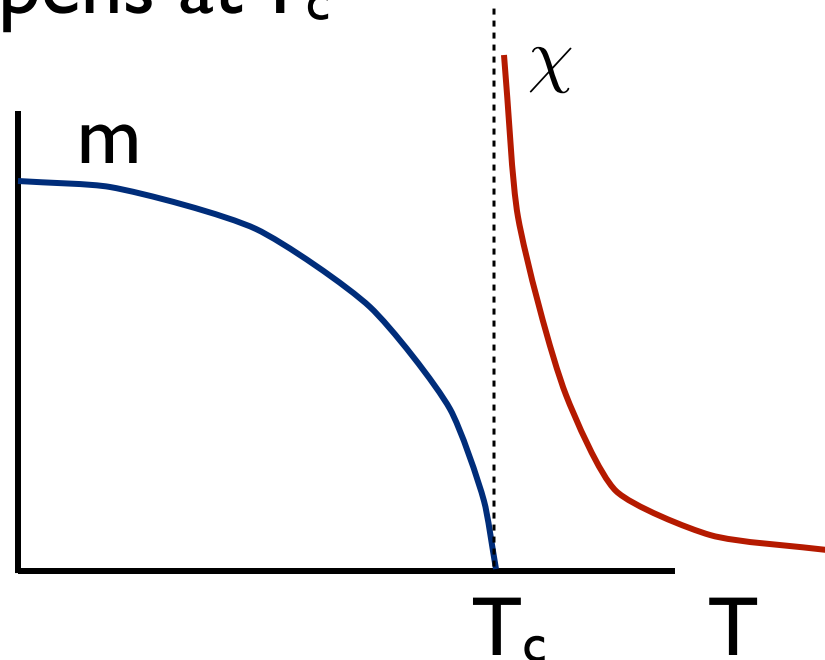
$$M/N = g\mu_B m \approx \frac{g\mu_B}{1 - \frac{zJ}{4kT}} \frac{g\mu_B H}{4kT} = \frac{(g\mu_B)^2 H}{4kT - zJ}$$

$$\chi = \frac{1}{N} \frac{\partial M}{\partial H} = \frac{A}{T - T_c} \quad \text{“Curie-Weiss law”}$$

Curie law is modified by shift of T by mean field T_c

Phase transition

- A lot happens at T_c



- Both $m(T)$ and $\chi(T)$ are *non-analytic* at T_c
- This is actually a sign of a *phase transition*