# Classical scaling for d>4

• Order parameter

$$m \sim b^{-(d-2)/2} \mathcal{M}(\pm 1, u|r|^{(d-4)/2})$$

 m vanishes for r>0 and again is singular for r<0 (m ~ u<sup>-1/2</sup>)

$$m \sim b^{-(d-2)/2} [u|r|^{(d-4)/2}]^{-1/2} \sim |r|^{1/2}$$
  
 $\beta = 1/2$ 

#### Back to Hertz critical point is "trivial" ?

 $S = \int \frac{d^d \mathbf{k} d\omega_n}{(2\pi)^{d+1}} \Big\{ (k^2 + \frac{|\omega_n|}{k^a} + r) |\Phi_{\mathbf{k},\omega_n}|^2 \Big\} + u \int \frac{d^{3d} \mathbf{k}_i d^3 \omega_{n,i}}{(2\pi)^{3d+3}} \Phi_{k_1,\omega_{n1}} \Phi_{k_2,\omega_{n2}} \Phi_{k_3,\omega_{n3}} \Phi_{-k_1-k_2-k_3,-\omega_{n1}-\omega_{n2}-\omega_{n3}} \Big\} + u \int \frac{d^{3d} \mathbf{k}_i d^3 \omega_{n,i}}{(2\pi)^{3d+3}} \Phi_{k_1,\omega_{n1}} \Phi_{k_2,\omega_{n2}} \Phi_{k_3,\omega_{n3}} \Phi_{-k_1-k_2-k_3,-\omega_{n1}-\omega_{n2}-\omega_{n3}} \Big\}$ 

- Additional ingredient for QCP: Temperature scaling:
  - relative to renormalized low energy scale, temperature *increases* under RG

$$k_B T \to b^z k_B T$$

• Also seen from action  $S = \int_0^p d\tau \cdots$ 



- Two relevant perturbations of QCP
  - r: deviation from critical point at T=0
  - T: temperature

$$r \to b^2 r$$
  $k_B T \to b^z k_B T$ 

# Quantum critical scaling

• Example: energy density

 $\varepsilon \sim b^{-(d+z)} \mathcal{E}(r b^2, k_B T b^z, u b^{4-d-z})$ 

• Let's sit *at* the QCP (r=0) and raise temperature

$$\varepsilon \sim b^{-(d+z)} \mathcal{E}(0, k_B T b^z, u b^{4-d-z})$$

$$\sim (k_B T)^{\frac{d+z}{z}} \tilde{\mathcal{E}}(u(k_B T)^{\frac{d+z-4}{z}}) \sim (k_B T)^{\frac{d+z}{z}}$$

• Specific heat

 $c_v \sim \partial \varepsilon / \partial T \sim T^{d/z} \sim T^{3/2}$  for 3d AF

# Quantum critical scaling

• Thermal expansion coefficient

$$\alpha = \frac{1}{V} \left. \frac{\partial V}{\partial T} \right|_p = -\frac{1}{V} \left. \frac{\partial S}{\partial p} \right|_T$$

- We can deduce entropy scaling from specific heat  $S \sim \int_0^T dT' \, \frac{C(T')}{T'} \sim T^{3/2}$
- Hence

$$S \sim T^{3/2} \mathcal{S}(rT^{-2/z})$$

For a pressure tuned transition then r ~ p

$$lpha \sim {\partial S \over \partial r} \sim T^{1/2}$$
 (it is usually linear in a metal)

### $Ce_{1-x}La_{x}Ru_{2}Si_{2}$

 This seems to be one of the rare examples where Hertz theory works



S. Kambe *et al*, JPSJ **65**, 3294 (1996)

Fit is to a (slightly) more sophisticated theory which includes  $r \neq 0$ 

### CeNi<sub>2</sub>Ge<sub>2</sub>

 Believed to be "close" to an AF QCP at ambient pressure

