

ENTANGLEMENT AND ENTROPY

0. CONTENTS

1. A First Example
2. Visualizing Entanglement
3. Density Matrices I
4. Density Matrices II
5. Entanglement Entropy

→ Two key tools: tensor products \otimes "AND" \neq superposition $+$ "OR"

1. A FIRST EXAMPLE

Let's take a quantum-mechanical exam.
What strategies do we have?

1. Will it rain tomorrow?

A Yes.

B No.

Options

1. A
2. A

1. B
2. B

2. Will you wear a raincoat?

A Yes.

B No.

1. A or B
2. A or B

1. A or 1. B
2. A or 2. B

① Classical guess, $|\psi_A\rangle = |A_1\rangle \otimes |A_2\rangle$, or

$$|\psi_B\rangle = |B_1\rangle \otimes |B_2\rangle.$$

② Superposition I, not entangled:

$$|\psi_s\rangle = \frac{1}{\sqrt{2}} \left[|A_1\rangle + |B_1\rangle \right] \otimes \frac{1}{\sqrt{2}} \left[|A_2\rangle + |B_2\rangle \right].$$

③ Superposition II, entangled:

$$|\psi_e\rangle = \frac{1}{\sqrt{2}} \left[\left(|A_1\rangle \otimes |A_2\rangle \right) + \left(|B_1\rangle \otimes |B_2\rangle \right) \right].$$

Notation:

- $|\psi_A\rangle = |A_1\rangle |A_2\rangle = |A_1 A_2\rangle$,
- $|\psi_B\rangle = |B_1\rangle |B_2\rangle = |B_1 B_2\rangle$;

NOT a product state!

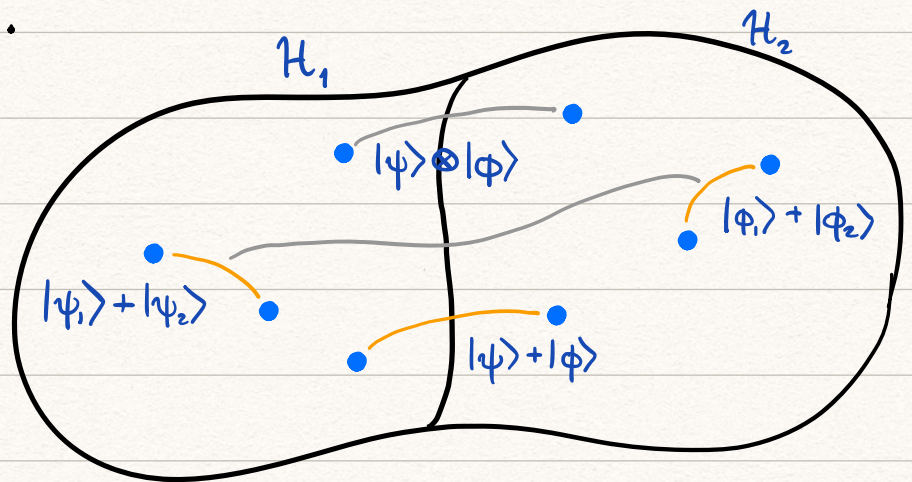
• $|\psi_s\rangle \propto |A_1 + B_1\rangle |A_2 + B_2\rangle$,

• $|\psi_e\rangle \propto |A_1 A_2\rangle + |B_1 B_2\rangle$.

2. VISUALIZING ENTANGLEMENT

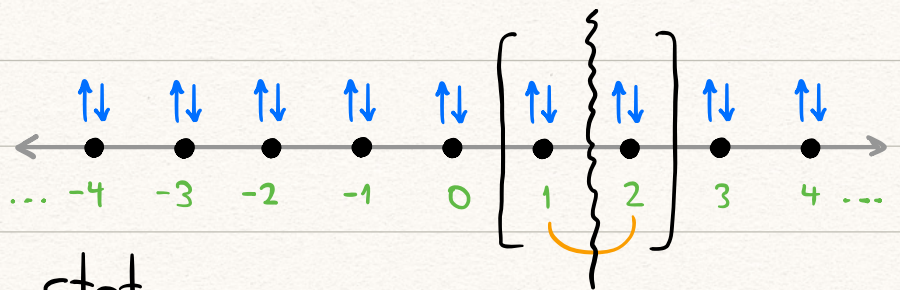
Let $\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2$.

E.g. Two identical particles.



E.g. Spin chain!
[1D lattice]

Q: How do we specify the full state of the system?



A: 3 options:

- * Specify \uparrow or \downarrow on each site
- * Build local superpositions of \uparrow & \downarrow
- * Superpose states nonlocally, across sites!

E.g. $|\psi_e\rangle = \frac{1}{\sqrt{2}} [|\uparrow\rangle_1 |\downarrow\rangle_2 - |\downarrow\rangle_1 |\uparrow\rangle_2]$ (EPR state)

Q: How entangled can a state be?

3. DENSITY MATRICES I [Measuring entanglement: motivation & ideas]

Big idea: nonlocal correlations are stored in the "threads" across tensor factors H_i .
(lattice sites i)

Break these threads, then count broken connections.
That is, forget one of the tensor factors.

E.g. $|\psi_e\rangle \rightsquigarrow \frac{1}{\sqrt{2}} [|\uparrow\rangle_1 (?)_2 - |\downarrow\rangle_1 (?)_2] \equiv \text{Red}_1(|\psi_e\rangle)$.

(Note: Arrows point from the $(?)_2$ terms to $|\uparrow\rangle_2$ or $|\downarrow\rangle_2$)

The resulting object is not a state:
it's a probability distribution of states!

Two options, \uparrow or \downarrow , for each of the $(?)$'s.

Each occurs w/ equal probability $\frac{1}{2}$.

Capture all 4 possibilities in a matrix:

$\rho_1 \stackrel{?}{=} \begin{pmatrix} \frac{1}{4} & \frac{1}{4} \\ \frac{1}{4} & \frac{1}{4} \end{pmatrix}$ first attempt to describe mixed states

[Reduced density matrix]

WRONG! Do NOT do this.

Better: ① & ② are entangled, so forgetting ② leaves ① in a classical mixture of $|\uparrow\rangle$ & $|\downarrow\rangle$.
Hence, $\rho_1 = \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}$. ← correct!

4. DENSITY MATRICES II [Measuring entanglement: formalism & calculations]

Density matrices: eigenvalues are probabilities.

Def. A density matrix ρ satisfies 3 properties:

- * ρ is Hermitian - ρ has enough eigenvalues
- * ρ is positive-definite - eigenvalues are positive
- * $\text{Tr}(\rho) = \sum_i \rho_{ii} = 1$. - eigenvalues sum to 1.

Pure states have only one nonzero eigenvalue:

$\rho = \begin{pmatrix} 1 & & \\ & 0 & \\ & & \ddots \\ & & & 0 \end{pmatrix}$. This is a projection:

$$\begin{aligned} \rho &= |\psi\rangle\langle\psi|, \text{ and } \rho^2 = (|\psi\rangle\langle\psi|)(|\psi\rangle\langle\psi|) = \\ &= |\psi\rangle\langle\psi|\psi\rangle\langle\psi| = \\ &= |\psi\rangle\langle\psi| = \rho. \end{aligned}$$

Mixed states contain a distribution $\{\lambda_1, \dots, \lambda_n\}$ of pure states $|\psi_i\rangle$ w/ probabilities λ_i .

How much don't we know? How mixed is ρ ?
How wide is the probability distribution of λ_i ?

Answer: $S_{vN}[\rho] \equiv -\sum_{i=1}^n \lambda_i \ln(\lambda_i) = -\text{Tr}[\rho \ln(\rho)]$.

E.g. If $\rho = |\psi\rangle\langle\psi|$ is pure, then $\lambda_i \in \{0, 1\}$, so

$$S_{vN}[\rho] = -\sum_{i=1}^n \lambda_i \ln(\lambda_i) = \begin{cases} 0 \ln(0) = 0 \\ 1 \ln(1) = 0 \end{cases} = 0.$$

E.g. If $\rho = \frac{1}{n} \mathbb{1} = \begin{pmatrix} \frac{1}{n} & & \\ & \ddots & \\ & & \frac{1}{n} \end{pmatrix}$, then all $\lambda_i = \frac{1}{n}$, so

$$S_{vN}[\rho] = -\sum_{i=1}^n \frac{1}{n} \ln\left(\frac{1}{n}\right) = -\frac{n}{n} \ln\left(\frac{1}{n}\right) = \ln(n).$$

5. ENTANGLEMENT ENTROPY

Mixed states arise from (ignorance of) entanglement, so we can use S_{vN} to measure entanglement!

To do so, follow these 3 steps:

① Start with a pure state $\rho = |\psi\rangle\langle\psi|$ in a composite system $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$. It could be entangled, i.e. $|\psi\rangle \neq |\phi_A\rangle \otimes |\phi_B\rangle$.

② Take the partial trace over \mathcal{H}_B :
 $\rho_A = \text{Tr}_B(\rho)$ ← reduced density matrix.

N.B. If $|\psi\rangle = |\phi_A\rangle|\phi_B\rangle$ is a tensor product, then ρ AND ρ_A are pure. Otherwise, ρ_A becomes mixed due to entanglement.

③ The entanglement entropy of ρ on A is the von Neumann entropy of ρ_A :

$$S_A[\rho] \equiv S_{\text{vN}}[\rho_A] = -\text{Tr}_A[\rho_A \ln(\rho_A)].$$

→ My research: computing S_A in QFT, and using AdS/CFT to make dual statements about quantum gravity.