

# SOUND WAVES

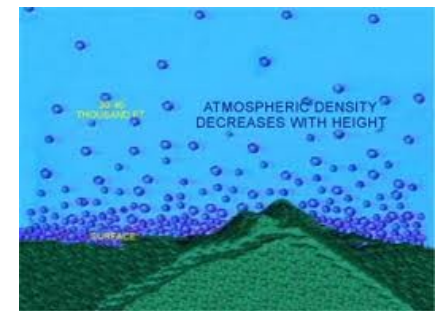
# Mathematical Models of Fluids

- Fluids → molecules roam and collide – no “springs”
  - Collisions cause **pressure** in fluid (Units: Pascal – Pa = N/m<sup>2</sup>)
- 2 mathematical models for fluid motion:
  - 1) “Bulk” properties – pressure/density at each point
  - 2) Statistical treatment of molecules – velocity/temperature

$$p(x, y, z, t) = p_0$$

- Equilibrium state: uniform pressure/density:  $\rho(x, y, z, t) = \rho_0$

- External forces can cause “**pressure gradient**”
- Example: uniform gravity acting on air
- Assume  $p = C \rho \rightarrow$  calculate  $p(z)$  and  $\rho(z)$

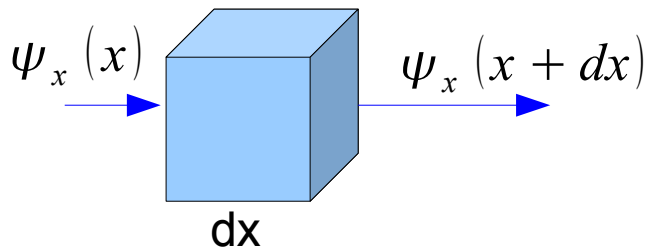


- In general, force due to non-uniform pressure:

$$d \vec{F}_p = -(\vec{\nabla} p) dV$$

# Pressure and Displacement in Fluids

- Force on fluid  $\rightarrow$  displacement and pressure gradient
  - Example: firecracker – air is pushed outward for an instant
  - Sketch graphs of  $\psi_r$  vs.  $r$  and  $p$  vs.  $r$  just after explosion
- Recall definition of Bulk Modulus:  $B = -V \frac{\Delta p}{\Delta V}$



This parcel of fluid is “stretched” ( $dV > 0$ ) due to uneven  $\Psi$

Note:  $dV$  for a “parcel” of fluid depends on the **difference** in  $\Psi$  from one side of parcel to other

$$\Delta V = (\Delta \psi_x) dy dz + (\Delta \psi_y) dx dz + (\Delta \psi_z) dx dy$$

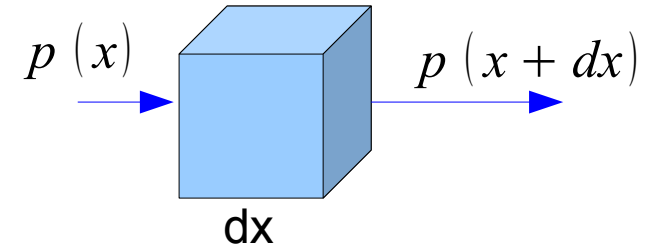
$$\Delta V = \left( \frac{\partial \psi_x}{\partial x} + \frac{\partial \psi_y}{\partial y} + \frac{\partial \psi_z}{\partial z} \right) dx dy dz$$

$$\frac{\Delta V}{V} = \frac{\Delta V}{dx dy dz} = \vec{\nabla} \cdot \vec{\psi}$$

$$\Delta p = -B (\vec{\nabla} \cdot \vec{\psi})$$

# Sound Waves in Fluids

- Newton's 2<sup>nd</sup> Law:  $d \vec{F} = dm \vec{a}$



$$-(\vec{\nabla} p) dV = (\rho dV) \left( \frac{\partial^2 \vec{\psi}}{\partial t^2} \right)$$

Cancel dV and take  
divergence of both sides:

$$\vec{\nabla} \cdot (-\vec{\nabla} p) = \vec{\nabla} \cdot \left( \rho \frac{\partial^2 \vec{\psi}}{\partial t^2} \right)$$

$$-\nabla^2 p = \rho \frac{\partial^2 (\vec{\nabla} \cdot \vec{\psi})}{\partial t^2} = \rho \frac{\partial^2 p}{\partial t^2} \frac{1}{B}$$

Plug in from the definition  
of bulk modulus:

$$\nabla^2 p = \frac{\rho}{B} \frac{\partial^2 p}{\partial t^2}$$

This is the **3-D wave equation**  
using **pressure** as the variable

$$v = \sqrt{\frac{B}{\rho}}$$

- Pressure of neighboring parcels of fluid
  - Has similar physical effect to “springs” – restoring force
  - Leads to waves:  $P = p_0 + p(k_x x + k_y y + k_z z \pm \omega t)$

# Transforming Derivatives

- In Cartesian coordinates:  $\nabla^2 p = \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) p$ 
  - Commonly, sound waves expand from a **point** or **line** source
  - How to express this in **spherical** or **cylindrical** coordinates?

$$dp = \frac{\partial p}{\partial x} dx + \frac{\partial p}{\partial y} dy + \frac{\partial p}{\partial z} dz = \frac{\partial p}{\partial r} dr + \frac{\partial p}{\partial \theta} d\theta + \frac{\partial p}{\partial \phi} d\phi$$

To transform **derivatives** to a different coordinate system:

Must express **dr, dθ, dφ** in terms of **dx, dy, dz**

Then plug in to 3-D wave equation above

Spherical – Partial derivatives expressed in matrix form:

$$\begin{pmatrix} dx \\ dy \\ dz \end{pmatrix} = \begin{pmatrix} \sin \theta \cos \phi & r \cos \theta \cos \phi & -r \sin \theta \sin \phi \\ \sin \theta \sin \phi & r \cos \theta \sin \phi & r \sin \theta \cos \phi \\ \cos \theta & -r \sin \theta & 0 \end{pmatrix} \begin{pmatrix} dr \\ d\theta \\ d\phi \end{pmatrix}$$

$$\begin{pmatrix} dr \\ d\theta \\ d\phi \end{pmatrix} = \begin{pmatrix} \frac{\sin \theta \cos \phi}{r} & \frac{\sin \theta \sin \phi}{r} & \frac{\cos \theta}{r} \\ \frac{-\cos \theta \cos \phi}{r} & \frac{-\cos \theta \sin \phi}{r} & \frac{\sin \theta}{r} \\ \frac{-\sin \phi}{r \sin \theta} & \frac{\cos \phi}{r \sin \theta} & 0 \end{pmatrix} \begin{pmatrix} dx \\ dy \\ dz \end{pmatrix}$$

# 3-D Wave Equation

$$\left( \frac{\partial}{\partial x} \right) p = \left( \sin \theta \cos \phi \frac{\partial}{\partial r} - \frac{\cos \theta \cos \phi}{r} \frac{\partial}{\partial \theta} - \frac{\sin \phi}{r \sin \theta} \frac{\partial}{\partial \phi} \right) p$$

- Repeat for y,z → plug in and use chain rule
  - To get full 3-D wave equation in spherical coordinates
- Common simplifying assumptions:
  - **Spherical symmetry** (i.e. wave is isotropic:  $\frac{\partial p}{\partial \theta} = \frac{\partial p}{\partial \phi} = 0$ )
  - In this case:  $\nabla^2 p = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial p}{\partial r} \right)$
  - **Cylindrical symmetry** (i.e. “line” source:  $\frac{\partial p}{\partial \phi} = \frac{\partial p}{\partial z} = 0$  )
  - In this case:  $\nabla^2 p = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial p}{\partial \rho} \right)$

# Sound Wave Intensity

- Sound wave – at any given instant:
  - PE stored in pressure gradient (i.e. non-uniform pressure)
  - KE stored in motion of fluid
  - Energy travels in a particular direction
  - “Direction of **propagation**” – determined by  $k_x$ ,  $k_y$ ,  $k_z$
- Wave **Intensity** (Units:  $\text{W/m}^2$  )
  - Average **power per unit area** transmitted by wave
  - Area must be measured perpendicular to propagation
  - For linear waves: Intensity is proportional to **(amplitude)<sup>2</sup>**
  - Human ear can detect incredibly wide range of intensities
  - Quietest:  $10^{-12} \text{ W/m}^2$       Loudest (before injury):  $10 \text{ W/m}^2$

# Examples

- Consider a traveling longitudinal “plane wave”:

$$p(x, y, z, t) = p_0 \cos(k_x x - \omega t)$$

- Given the bulk modulus B and density  $\rho$ :
- Calculate the displacement field in terms of B,  $\rho$ ,  $p_0$ ,  $k_x$
- Calculate the wave intensity in terms of B,  $\rho$ ,  $p_0$ ,  $k_x$
- How might one physically create a plane wave?

- Consider an outgoing spherical wave:

$$p(r, t) = p_0 \left( \frac{r_0}{r} \right) \sin(kr - \omega t)$$

- Show that this satisfies wave equation in spherical coordinates
- Calculate the wave intensity as a function of r



# Decibel Scale

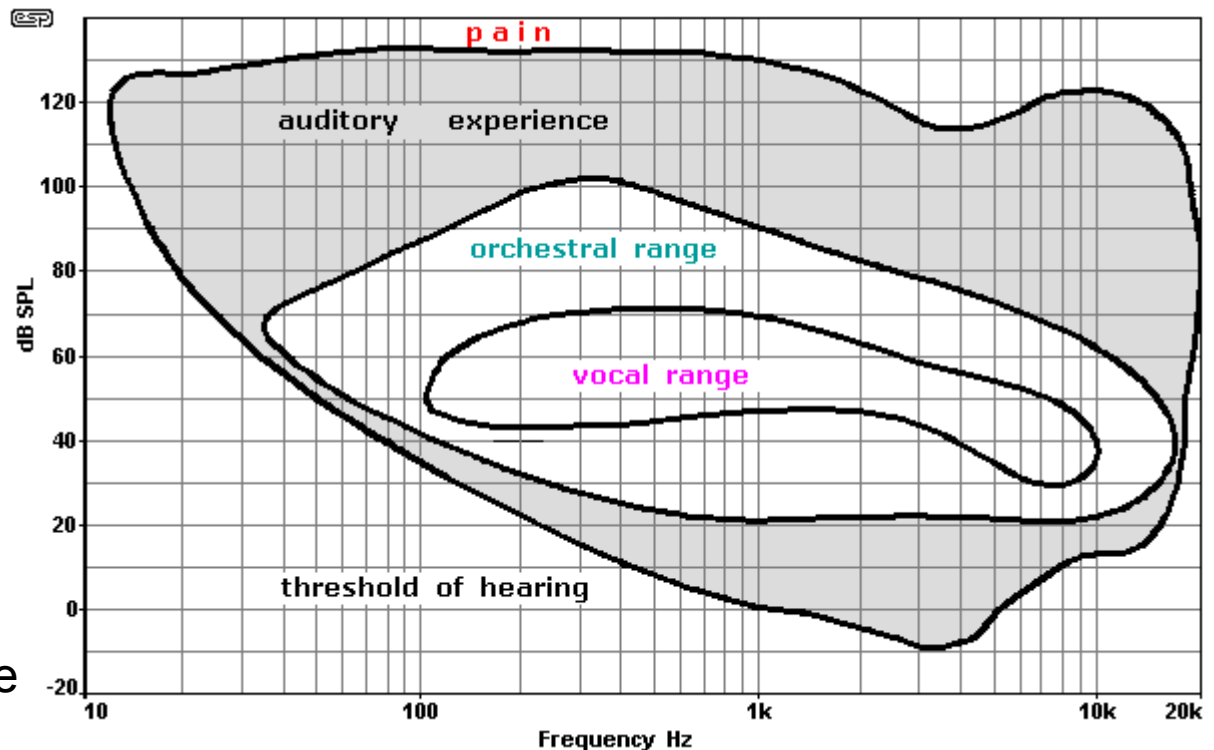
- Intensity / frequency ranges of human hearing: **wide**
  - Loudest vs. quietest sound: 12 orders of magnitude in intensity
  - But **perceived loudness** is not “a trillion times louder”
- Decibel Scale – designed based on **human** perception

$$\beta = 10 \log_{10} \left( \frac{I}{I_0} \right)$$

$$I_0 = 10^{-12} \frac{W}{m^2}$$

Logarithmic scale:

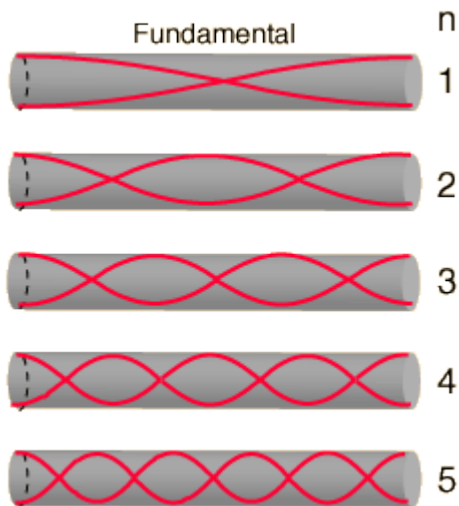
“Compresses” wide range of numbers to more manageable range



# Wind Instruments

- 3-D sound waves can be reduced to 1-D
  - By use of a tube or pipe → standing waves inside
  - These “wind instruments” have similar physics to strings
  - Frequency of standing waves depends on:
    - 1) length of tube      2) whether ends are open or closed

Both ends open:



$$\lambda_n = \frac{2L}{n}$$

$$f_n = n f_1$$

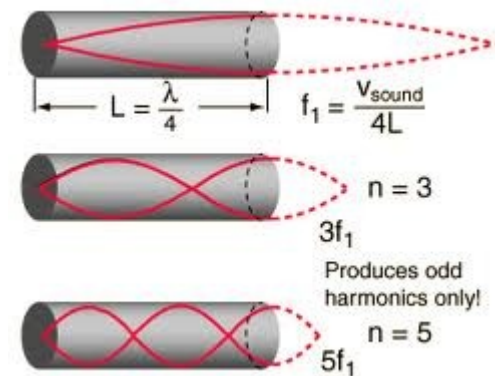
(exactly  
like  
strings)

One end closed:

$$\lambda_n = \frac{2L}{n + \frac{1}{2}}$$

$$f_n = (2n + 1) f_1$$

Examples: reed-based instruments,  
some organs



# Doppler Effect

- Relative velocity of sound source/listener:
  - Perceived frequency/wavelength **differs** from source
  - Relative motion toward each other  $\rightarrow f_L > f_S$

$$f_L = \left( \frac{v \pm v_L}{v \pm v_S} \right) f_S$$

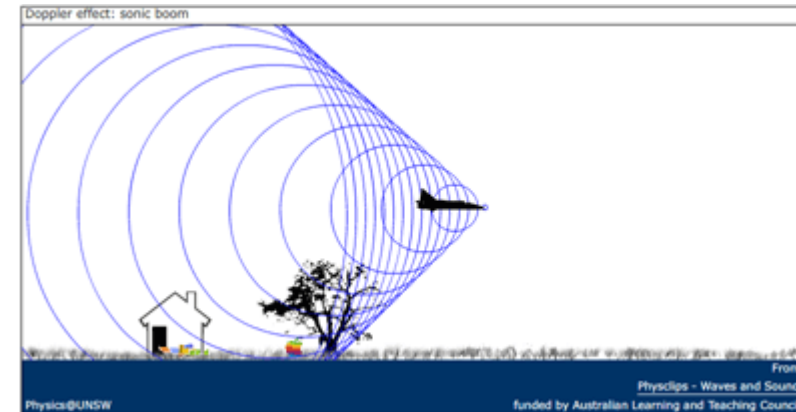


- Note:  $v_L$  and  $v_S$  must be measured relative to medium
  - i.e. there is a “preferred” reference frame for sound waves
- Example: bat chasing moth – uses “sonar”
    - Given  $v_{\text{bat}}$  and  $v_{\text{moth}}$  relative to air:
    - Calculate ratio  $f_{\text{perceived}} / f_{\text{emitted}}$
    - Make sure to pick the correct signs!



# Shock Waves

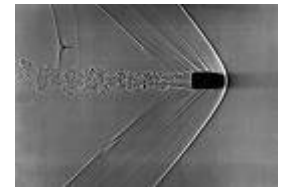
- Problem with Doppler shift formula:
  - What happens as  $v_s$  approaches speed of sound?
  - Sound waves do not “escape” source quickly
  - Many waves build up in small region in front of source:



- Shock wave (nonlinear)
  - Small area of very high intensity
  - Produced when linear waves can not move energy fast enough

- Examples:

- Supersonic jet, bullet → conical shock wave → sonic boom
- Atomic bomb → spherical shock wave → huge impact force



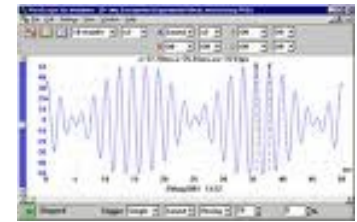
# Beats

- Consider the combination of 2 pure tones:

$$p(t) = p_0 \cos(\omega_1 t) + p_0 \cos(\omega_2 t)$$

- Plug in the trig identity  $\cos(A) + \cos(B) = 2 \cos\left(\frac{A+B}{2}\right) \cos\left(\frac{A-B}{2}\right)$

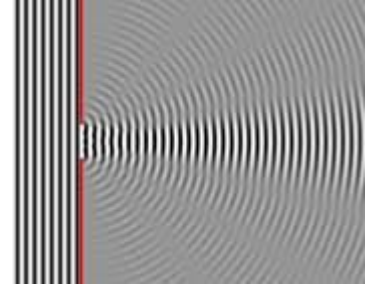
$$p(t) = \left[ 2 p_0 \cos\left(\left(\frac{\omega_1 - \omega_2}{2}\right) t\right) \right] \cos\left(\left(\frac{\omega_1 + \omega_2}{2}\right) t\right)$$



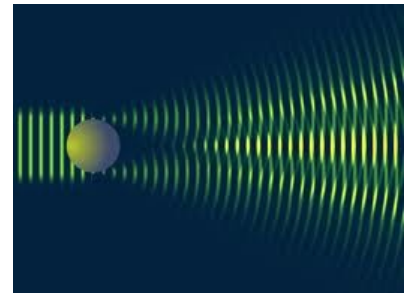
- Single tone  $\rightarrow$  frequency = average of 2 pure tones
  - **Amplitude** varies  $\rightarrow$  frequency = half difference of 2 pure tones
  - Perceived sound: single tone with “wah-wah-wah”
- “**Beat frequency**” defined as  $|f_1 - f_2| \rightarrow$  uses:
  - **Tuning** musical instruments  $\rightarrow$  adjust until beats vanish
  - **Detection** of low speeds using Doppler reflection

# Diffraction

- Consider waves at the “edge” of a boundary:
  - Can a “plane wave” be confined to a small area?



- Wave equation does not allow for abrupt wave “edges”
  - If  $\nabla^2 p \neq 0$  then  $\frac{\partial^2 p}{\partial t^2} \neq 0 \rightarrow$  wave “spreads out”
  - Also occurs when waves encounter obstacles



- **Diffraction** – highly dependent on wavelength
  - Long wavelengths “bend” more easily
  - Short wavelengths  $\rightarrow$  more “directional”
  - Example: Bass sound tends to “fill” a room

