# PHYS 21: Assignment 6 Solutions

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## $\rm K$  + K 10.2



## $K + K$  10.5

The equation of motion for the damped oscillator is:

 $0=\ddot{x} + \gamma \dot{x} + \omega^2 x$ 

For  $\gamma = 2\omega_0$ , we have:

 $0 = \ddot{x} + 2\omega_0 \dot{x} + \omega_0^2 x$ 

Let's use a trial solution of

$$
x(t) = (A + Bt)e^{-\gamma t/2}
$$
  
\n
$$
= (A + Bt)e^{-\omega_0 t}
$$
  
\n
$$
\Rightarrow \dot{x} = -\omega_0 (A + Bt)e^{-\omega_0 t} + Be^{-\omega_0 t}
$$
  
\n
$$
\ddot{x} = \omega_0^2 (A + Bt)e^{-\omega_0 t} - 2B\omega_0 e^{-\omega_0 t}
$$
  
\n
$$
0 = \omega_0^2 (A + Bt)e^{-\omega_0 t} - 2B\omega_0 e^{-\omega_0 t} - 2\omega_0^2 (A + Bt)e^{-\omega_0 t} + 2\omega_0 Be^{-\omega_0 t} + \omega_0^2 (A + Bt)e^{-\omega_0 t}
$$
  
\n
$$
0 = 2\omega_0^2 (A + Bt)e^{-\omega_0 t} - 2\omega_0^2 (A + Bt)e^{-\omega_0 t} - 2B\omega_0 e^{-\omega_0 t} + 2\omega_0 Be^{-\omega_0 t}
$$
  
\n
$$
= 0
$$

Now we apply the initial conditions:

 $x(0) = 0$  $\Rightarrow 0 = A$  $\dot{x}(0) = \frac{I}{m}$  $\Rightarrow$  $\frac{I}{m} = B$  $\Rightarrow$   $x(t) = \frac{It}{m}e^{-\omega_0 t}$ 

We see that  $\dot{x} = 0$  when  $t = \frac{2}{\gamma}$ , and since the second derivative is always negative this is the maximum.

#### $\overline{\rm{K} + \rm{K}$  10.7

We can see from Equation 10.25 of the textbook that  $x(t)$  has the form:

$$
x(t) = A \cos(\omega t + \varphi)
$$

$$
\varphi = \arctan \frac{\gamma \omega}{\omega_0^2 - \omega^2}
$$

So we see that:

 $\dot{x}(t) = -A\omega\sin(\omega t + \varphi)$ 

To be "in phase" with the driving force, we need

$$
-\sin(\omega t + \varphi) = \cos(\omega t)
$$

$$
-\sin(\omega t)\cos(\varphi) - \sin(\varphi)\cos(\omega t) = \cos(\omega t)
$$

This means that  $\sin \varphi = -1$  and  $\cos \varphi = 0$  so  $\varphi = \frac{3\pi}{2}$ . But this forces  $\tan \varphi = 0$ , which occurs when  $\omega = \omega_0$ , or when resonance occurs.

However much power is generated by dropping the weight is dissipated by friction. If the weight has descended a distance *L* in a time *T*, then:

$$
\bar{p} = \frac{\Delta E}{\Delta t} = \frac{MgL}{T}
$$

And we'd like to find the energy lost per radian swung:

$$
\bar{\Delta E} = \frac{\bar{p}}{\omega}
$$

We can also write down the average stored energy

$$
\bar{E} = \frac{m}{2}\bar{v^2} + \frac{mg\ell}{2}\bar{\theta^2}
$$

*On average*, the kinetic and potential energies are equal so if  $\theta_0$  is the angle at the maximum of the swing:

$$
\begin{aligned}\n\bar{E} &= \frac{mg\ell}{2}\theta_0^2\\ \Rightarrow Q &= \frac{\bar{E}}{\bar{\Delta E}}\\ &= \frac{\omega}{\bar{p}} \frac{mg\ell}{2} \theta_0^2\\ &= \frac{\omega T}{MgL} \frac{mg\ell}{2}\n\end{aligned}
$$

We can now plug in the numbers the problem gives us to find:

*Q* = 68

 $\theta_0^2$ 

The amount of energy required to drive the clock for one day is:

$$
E = MgL
$$
  
=0.2kg · g · 2m  
=4J

If the battery only holds 1 joule of power, then it will only make it through a quarter of the day (6 hours).

#### $\rm K$  + K 10.12

The equation of motion for the free damped oscillator is

 $\ddot{x} + \gamma \dot{x} + \omega_0^2 x = 0$ 

The equation of motion for the forced damped oscillator is

$$
\ddot{x} + \gamma \dot{x} + \omega_0^2 x = \frac{F_0}{m} \cos(\omega t)
$$

Let  $x_a(t)$  be a solution for the first equation and let  $x_b(t)$  be a solution for the second. We also say  $x_c(t) = x_a(t) + x_b(t)$ . Then we substitute  $x_c(t)$  into the second DE.

$$
\ddot{x}_c + \gamma \dot{x}_c + \omega_0^2 x_c = \ddot{x}_a + \gamma \dot{x}_a + \omega_0^2 x_a + \ddot{x}_b + \gamma \dot{x}_b + \omega_0^2 x_b
$$

$$
= \frac{F_0}{m} \cos(\omega t)
$$

So we see that  $x_c(t)$  satisfies this differential equation. Now we see that,

$$
x_a(t) = A\cos(\omega t + \varphi)
$$

And that

$$
x_b(t) = Be^{-\gamma t} \cos(\omega_1 t) + Ce^{-\gamma t} \sin(\omega_1 t)
$$

Where  $\omega_1^2 = \omega_0^2 - \gamma^2/4$ . Now the initial condition requires that

$$
x_c(0) = 0 \dot{x}_c(0) = 0
$$

The conditions require that:

$$
A\cos\varphi + B = 0
$$
  

$$
-A\omega\sin\phi - B\gamma + C\omega_1 = 0
$$

At resonance,  $\omega = \omega_0$ . Also,  $A = \frac{F}{m\omega_0 \gamma}$ .  $\varphi = \arctan \infty = \frac{\pi}{2}$ . Then:

$$
B = 0
$$
  

$$
C = \frac{A\omega}{\omega_1} = \frac{F_0}{m\omega_1\gamma}
$$
  

$$
\gamma \ll 2\omega_0 \Rightarrow C \approx \frac{F}{m\omega_0\gamma}
$$