

Fluid Statics

between fluids & solids  
 Division is not always clear

→ e.g., glass is a fluid

Church windows are thicker on bottom than on top.

Plastic is also intermediate

→ e.g. clay.

Can change state by changing temperature or

pressure → e.g., ~~saline~~ rock layers

→ rock flows at high pressure

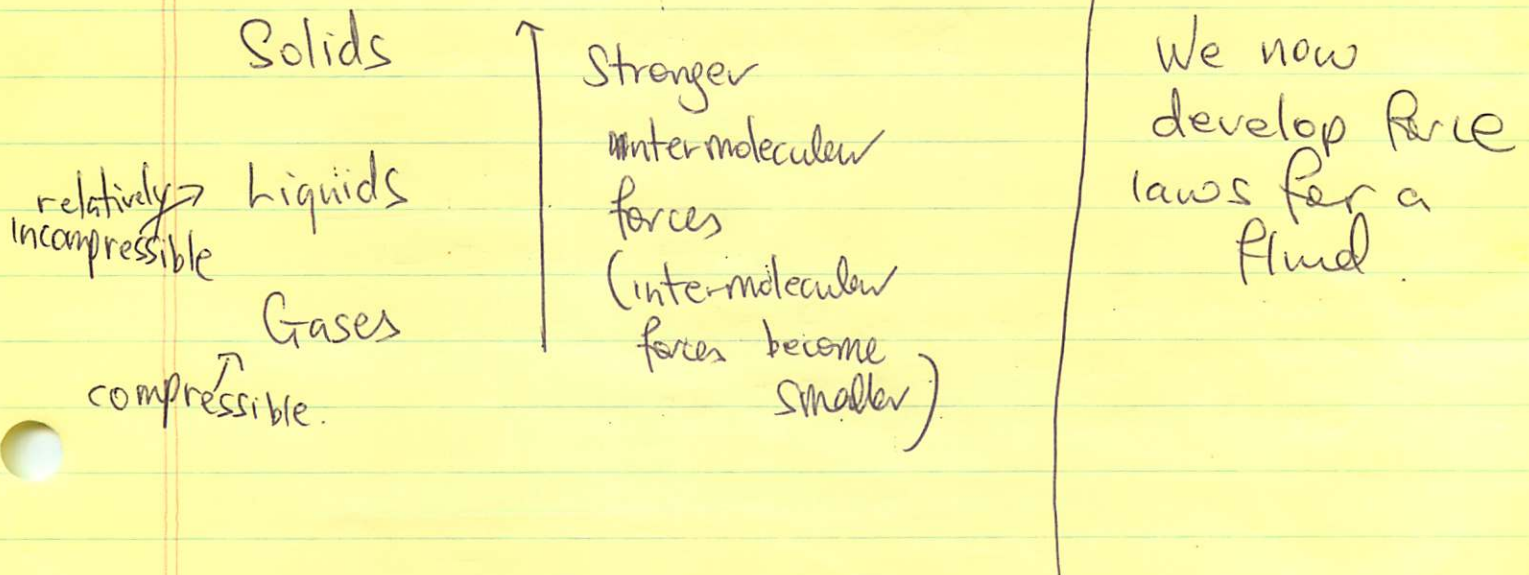
→ aluminium can be made into wires.

Plasmas ionized gas w. very different behaviour

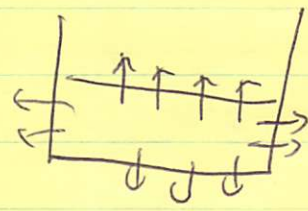
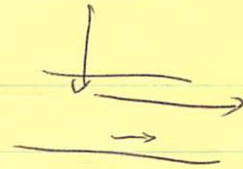
fr. ordinary gas.

Fluorescent light.

Sun.



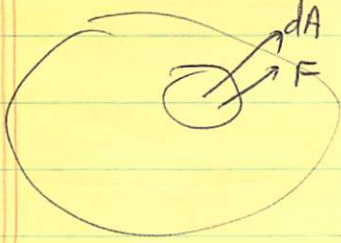
# Pressure



Fluid can flow - unable to sustain shear stress.  
 All forces between interior & exterior are at right angles to fluid boundary.

Pressure is a scalar, no direction.  $P = \frac{F}{A}$

Result of collisions of molecules with surface, imparts momentum ~~& vice versa~~



$$\Delta \vec{F} = p \Delta \vec{A}$$

$$\Rightarrow p = \lim_{\Delta \vec{A} \rightarrow 0} \frac{\Delta \vec{F}}{\Delta \vec{A}}$$

so that independent of size of element

Pressure can vary fr. point to point on surface.

But  $\nabla P$  is a vector (direction of force exerted by pressure gradient).

~~$\vec{P}$~~

$$1 \text{ Pa} = 1 \text{ Nm}^{-2}$$

$$1 \text{ atm} = 10^5 \text{ Pa}$$

$$\text{normal Blood pressure} = 1.6 \times 10^4 \text{ Pa}$$

$$\text{Center of Earth} \approx 4 \times 10^{11} \text{ Pa}$$

$$\text{" " Sun} \approx 2 \times 10^{16} \text{ Pa}$$

Density  $\rho = \lim_{\Delta V \rightarrow 0} \frac{\Delta m}{\Delta V}$  a scalar.

[If density is uniform, then  $\rho = \frac{m}{V}$ ]

Some examples  
Density can depend on  
pressure & T.  
(varies little for solids +  
liquids, lots for gas)

$$\rho (\text{interstellar space}) \approx 10^{-20} \text{ kg m}^{-3}$$

$$\rho (\text{water}) \approx 10^3 \text{ kg m}^{-3}$$

$$\rho (\text{air, at 1 atm}) \approx 1 \text{ kg m}^{-3}$$

$$\rho (\text{black hole}) \approx 10^{19} \text{ kg m}^{-3}$$

$B = - \frac{\Delta p}{\Delta V/V}$  Bulk modulus [has same units  
as p, since  $\frac{\Delta V}{V}$  is  
dimensionless].

stress  $\Delta p$   
strain  $\Delta V/V$   
makes B +ve (since  $\Delta p$  &  $\Delta V$  have  
opposite sign).

$$\Delta p > 0 \rightarrow \Delta V < 0$$

B is large.

Then large  $\Delta p \Rightarrow$  small  $\frac{\Delta V}{V}$   
resists compression

Water has  $B \approx 2.2 \times 10^9 \text{ N m}^{-2}$   
(nearly incompressible).

At bottom of Pacific ocean,  $p = 400 \text{ atm}$   
 $\approx 10^7 \text{ Pa}$

$$\Rightarrow \frac{\Delta V}{V} = - \frac{\Delta p}{\Delta B} \approx \frac{10^7}{2.2 \times 10^9} \approx 0.5\%$$

B large incompressible

B small compressible

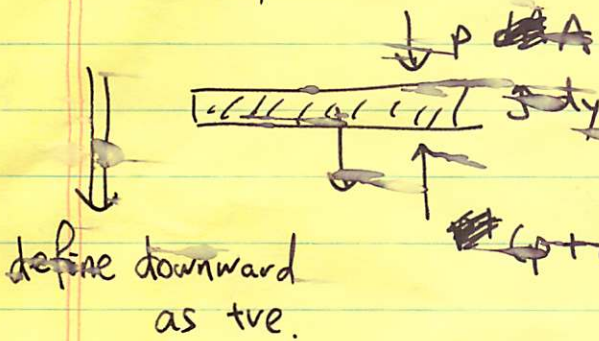
Gas have  $B \approx 10^5 \text{ Nm}^{-2}$

$$\Delta p \approx 0.1 \text{ atm} \approx 10^4 \text{ Pa}$$

$$\Rightarrow \frac{\Delta V}{V} = \frac{\Delta p}{B} \approx 0.1 = 10\%$$

### Variation of Pressure in Fluid at Rest

For fluid at rest, net force + net torque on every fluid element = 0.



All forces balance:

$$\begin{aligned}
 F_{\text{gravity}} &= \cancel{g A dy} \\
 &= \rho dV g \\
 &= \rho A g dy
 \end{aligned}$$

$$F_{\text{pressure}} = pA - (p + dp)A$$

$$F_{\text{gravity}} = F_{\text{pressure}}$$

$$\Rightarrow \rho A g dy = pA - (p + dp)A = -dpA$$

$$\Rightarrow \left( \frac{dp}{dy} = -\rho g \right) \text{ equation of hydrostatic equilibrium.}$$

Can integrate this:

$$\int_{p_1}^{p_2} dp = - \int_{y_1}^{y_2} \rho g dy$$

can take this out if  $g = \text{const}$  e.g. water

$$\Rightarrow (p_2 - p_1) = -\rho g (y_2 - y_1)$$



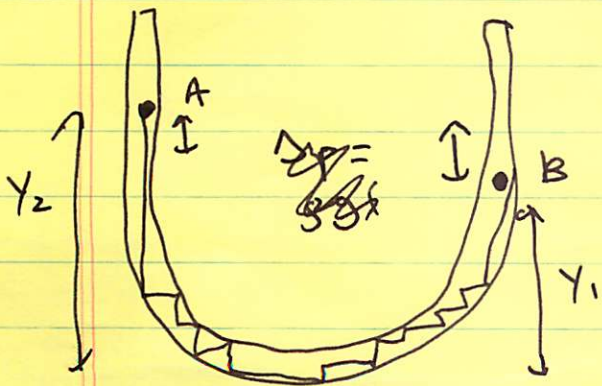
If fluid has free surface, then let that be the zero point

$$(p - p_0) = -\rho g (-h - 0)$$

$$\Rightarrow p = p_0 + \rho g h$$

Pressure increases with depth, but is the same for all points at the same depth.

True regardless of the shape of container!

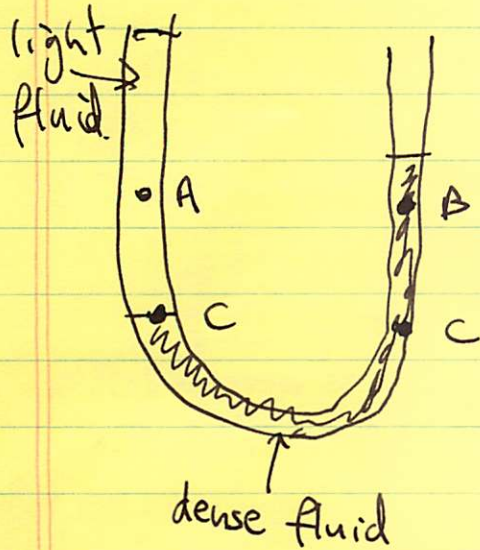


$$\Delta p = \rho g (y_2 - y_1)$$

Any 2 points can be connected by mix of horizontal & vertical steps, pressure doesn't change on

horizontal steps, pressure changes by  $\Delta p = -\rho g \Delta y$  on vertical steps.

(6)



If 2 immiscible fluids have different densities, pressure can be different at same level. on different sides.

Ask class

Where is the pressure greater, A or B?

Liquid between C on left & C on right in equilibrium  $\Rightarrow$  forces at C are same.

But pressure falls less between C & A than C & B, since A is less dense

$\Rightarrow$  (A) is at higher pressure.

### Variation of Pressure in Atmosphere

$\Delta p = -\rho g \Delta y$  assumed that density  $\rho$  was constant.

This is no longer true if  $\Delta y$  is large.

Ideal gas law  $P = \frac{\rho}{n m_p} k_B T$

$$\propto \rho T$$

$\propto \rho$  if  $T$  is const  
(not exactly true, only approximate).

(7)

Then write  $P = P_0 \left( \frac{\rho}{\rho_0} \right)$

$$\frac{dP}{dy} = -\rho g$$

$$\Rightarrow \frac{P_0}{\rho_0} \frac{d\rho}{dy} = -\rho g$$

$$\Rightarrow \frac{d\rho}{\rho} = -\frac{\rho_0}{P_0} g dy$$

$$\Rightarrow \ln \rho = -\frac{\rho_0}{P_0} g y + C$$

At  $y=0$ ,  $\rho = \rho_0$

$$\Rightarrow \rho = \rho_0 e^{-\frac{\rho_0}{P_0} g y}$$

$$P \propto \rho \Rightarrow P = P_0 e^{-\frac{\rho_0}{P_0} g y}$$

$$\frac{g \rho_0}{P_0} = \frac{10 \cdot 1}{10^5} \approx 10^{-4} \text{ m}^{-1} \quad (\text{must have units of } L^{-1})$$

$$\approx \text{~~10~~ } 0.1 \text{ km}^{-1} \quad (L^{-1})$$

or  $a \approx \left( \frac{g \rho_0}{P_0} \right)^{-1} \approx 10 \text{ km.}$

$$P = P_0 e^{-y/a}$$

(more careful

$$a = 8.55 \text{ km})$$

a is the <sup>pressure scale</sup> ~~e-folding~~ height: pressure drops by  $e^{-1}$  everytime go 1 a.

Hence drop by a factor 10 when

$$e^{-\frac{y}{a}} = \cancel{10} 0.1$$

$$\begin{aligned} \Rightarrow -\frac{y}{a} &= \ln 0.1 \Rightarrow y = a \ln 10 \\ &= 2.3a \\ &= \underline{20 \text{ km}} \end{aligned}$$

So drop by factor  $\sim 10$  every 20 km;  
at 40 km,  $P = 0.01 \text{ atm}$ .

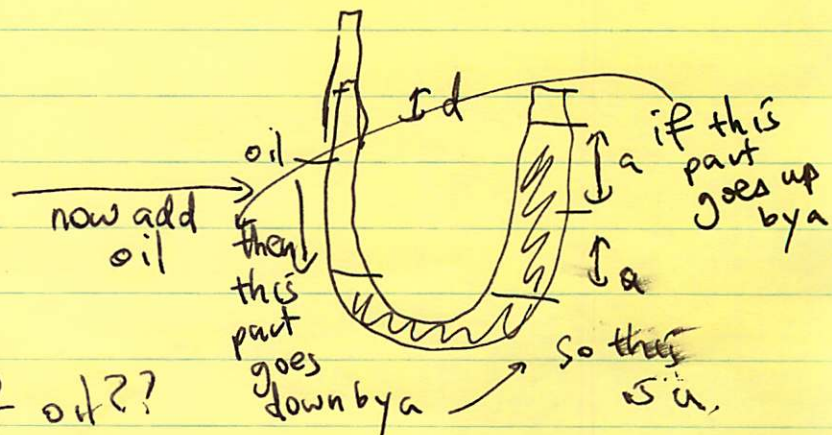
[Actual pressure drops a little faster than this, bcos temperature falls with height as well].

Note: pressure changes linearly with height for liquid but exponentially for gas. Liquids are incompressible, but gas is compressible  $\rightarrow$  density & hence weight of fluid above changes much faster w/ height.

### Example



U tube initially filled with water



What's the density of oil??



For U-shaped portion to be static, pressure has to be same on both sides.

Pressure on exposed ends has to be the same (atmospheric pressure).

$$\text{So } \rho_{\text{oil}}(2a + d) = \rho_{\text{water}} 2a$$

$$\Rightarrow \rho_{\text{oil}} = \left( \frac{2a}{2a + d} \right) \rho_{\text{water}}$$

\* This device is used to find relative densities (or specific gravity).

## Pascal's Principle & Archimedes Principle

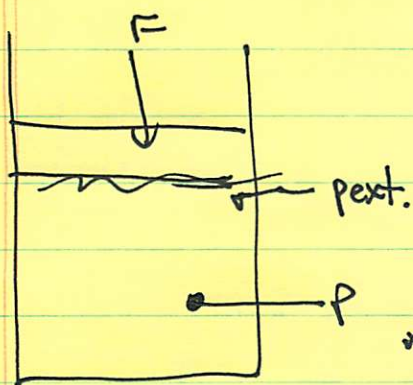
Pressure applied to an enclosed fluid is transmitted undiminished to every portion of the fluid and to the walls of the containing vessel.

↳ so if increase pressure at one location by  $\Delta P$ , that increase is felt every where!

↳ Examples

- squeeze toothpaste
- hydraulics [brakes, flaps on aircraft, dentist chair]
  - enables one to amplify a small force, as well as transmit forces over large distances.

(10)



Consider an incompressible fluid

$$p = p_{ext} + \rho gh.$$

↖ consider pressure here.

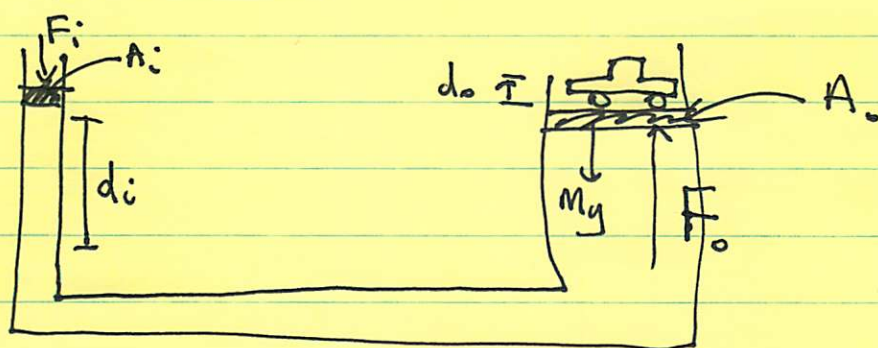
If add pressure  $\Delta p_{ext}$

$$\Delta p = \Delta p_{ext} + \Delta(\rho gh) \rightarrow 0 \text{ since incompressible}$$

(actually, is also true for compressible fluids  $\rightarrow$  due to change in density get initial sound waves which die out.)

E.g. when inflate a balloon w/ air, it expands evenly in all directions.

### Example Hydraulic Lever



Using hydraulic level, can apply a small force to lift a heavy object like a car!

$$p_i = p_o$$

$$\Rightarrow \frac{F_i}{A_i} = \frac{F_o}{A_o} \Rightarrow F_i = F_o \frac{A_i}{A_o} = M_g \frac{A_i}{A_o}$$

So if  $\left(\frac{A_i}{A_o}\right) \ll 1$ , can use small force to lift a heavy car.

~~Are we getting something for nothing??~~  
For incompressible fluids, volume of liquid displaced

$$\begin{aligned} \text{Vol of liquid displaced downwards} \\ = \text{Vol of liq moved upwards} \end{aligned}$$

$$\Rightarrow A_i d_i = A_o d_o$$

$$\Rightarrow d_o = d_i \left(\frac{A_i}{A_o}\right) \ll d_i \text{ if } \left(\frac{A_i}{A_o}\right) \ll 1$$

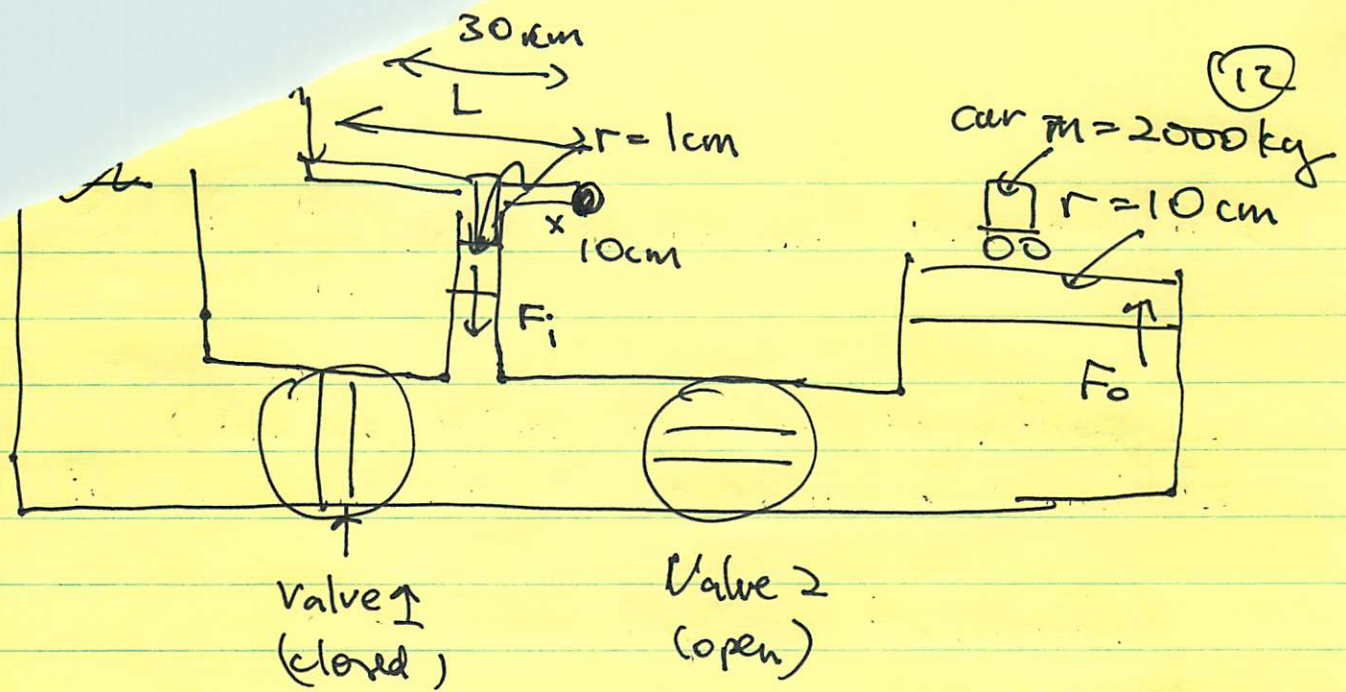
Work done

$$\therefore F_i d_i = F_o \left(\frac{A_i}{A_o}\right) d_o \left(\frac{A_o}{A_i}\right) = F_o d_o$$

→ energy is conserved (ignoring frictional losses) → not getting a free lunch.

### Example 2

Let's consider a hydraulic jack used to raise an automobile. How is it possible that our hands can raise such a heavy object?



- ~~F<sub>i</sub>~~ ① What force do we need to apply?  
 ② How far will the car move?

$$F_i = Mg \frac{A_i}{A_o} = 2000 (10) \frac{\pi (1)^2}{\pi (10)^2}$$

$$= 200 \text{ N}$$

There's also a lever on the pump.

Assuming angular acceleration of pump handle is negligible,

$$\sum \tau = F_h L - F_i x = 0$$

$$\Rightarrow F_h = F_i \left( \frac{x}{L} \right) = F_i \left( \frac{10 \text{ cm}}{30 \text{ cm}} \right)$$

$$\approx 67 \text{ N}$$

This is small! Like lifting a 6.7 kg weight.

~~d<sub>i</sub>~~ Suppose hand moves through distance  $h \approx 30 \text{ cm}$   
 Small piston moves through distance

$$d_i \approx h \left( \frac{x}{L} \right) \approx \frac{h}{3} \approx 10 \text{ cm.}$$

Large piston moves through distance:

$$d_o = d_i \left( \frac{A_i}{A_o} \right) = 10 \left( \frac{1}{100} \right) = 0.1 \text{ cm} \\ = 1 \text{ mm!}$$

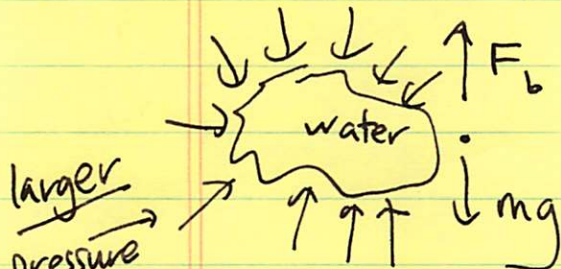
So need to apply many strokes to raise car.

Use valves to transfer fluid fr. reservoir to pistons:

|            |            |            |
|------------|------------|------------|
| downstroke | 1 (closed) | 2 (open)   |
| upstroke   | 1 (open)   | 2 (closed) |

To let car down  $\rightarrow$  both valves open, fluid goes back to reservoir.

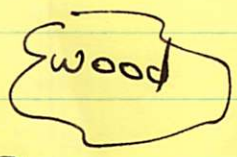
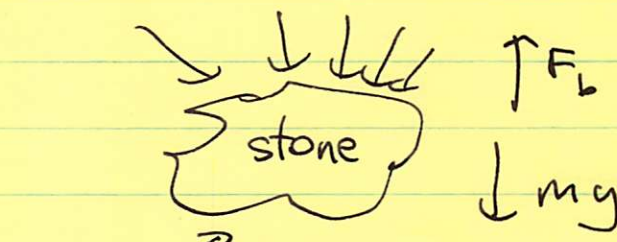
# Archimedes' Principle



larger pressure (since pressure increases w/ depth)

Water doesn't rise or fall

$F_b = mg$  balance  
(vector sum of all forces on surface)



If they are the same shape (same vol) then  $F_b$  must be the same as for the parcel of water.

Will then rise ( $F_b > mg$ )  
or sink ( $F_b < mg$ )

depending on  $mg \rightarrow M = \rho V$  since  $\rho$  fixed, depends on relative density.

Hence

buoyancy force depends on pressure difference between top & bottom of fluid element.

So we have Archimedes' principle:

A body wholly or partially immersed in a fluid is buoyed up by a force equal in magnitude to the weight of fluid displaced by the body.

Denser than water  $\rightarrow$  sinks  
less dense than water  $\rightarrow$  rises.

Also: things weigh less in water, due to the buoyant

force. So, for example:  $\rightarrow$  astronauts practice in underwater tanks, to simulate weightless conditions in space.

Object less dense than water floats on surface.

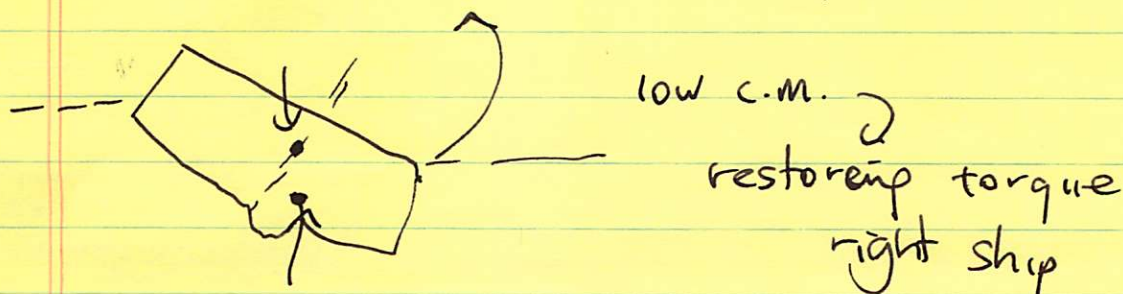
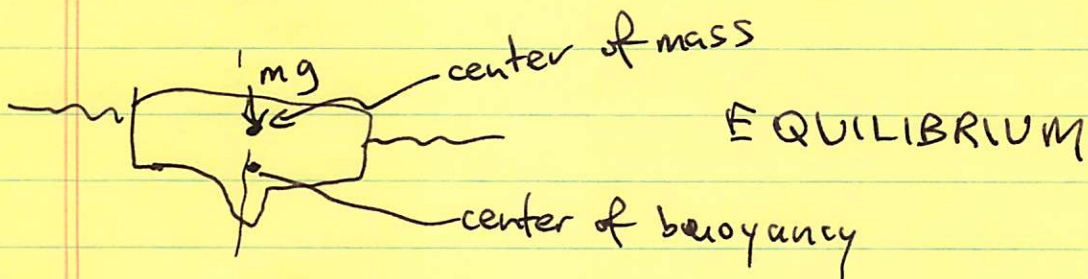
Submerged portion has ~~the~~ volume  $V$   
s.t.

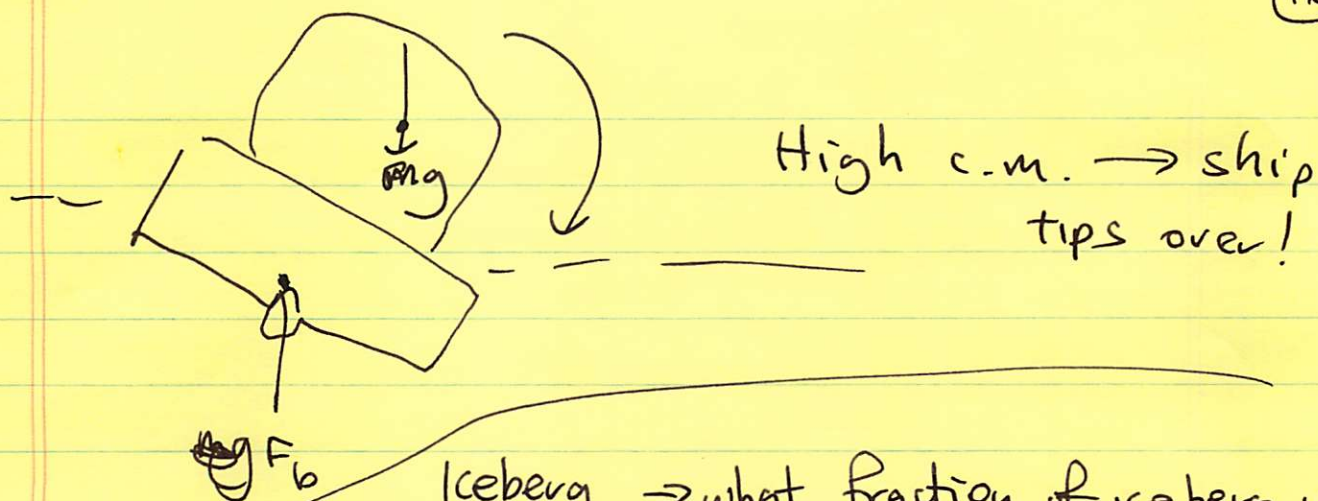
$$\rho_{\text{water}} V = m_{\text{obj}} g \quad \leftarrow \text{equilibrium condition}$$

Buoyant force acts on center of buoyancy = center of gravity of displaced fluid.

Gravity acts on center of mass.

As ship tips, center of buoyancy moves  $\rightarrow$  depending on relative position of c. buoyancy & c. mass, torques develop which restore balance or tip the ship over.





Iceberg  $\rightarrow$  what fraction of iceberg is exposed, given that

$$\rho_{\text{water}} = 1024 \text{ kg m}^{-3}$$

$$\rho_{\text{ice}} = 917 \text{ kg m}^{-3}$$

weight  $W_i = \rho_i V_i g$

buoyant force  $F_b = \rho_w V_w g$   $\leftarrow$  vol. of water displaced

$$W_i = F_b$$

$$\rho_w V_w g = \rho_i V_i g$$

$$\Rightarrow \frac{V_w}{V_i} = \frac{\rho_i}{\rho_w} = \frac{917 \text{ kg m}^{-3}}{1024 \text{ kg m}^{-3}} = 89.6\%$$

So only  $\sim 10\%$  of iceberg is above water!

## Measurement of Pressure

How to do it? — static (discuss now)

— dynamic (based on speed of moving fluid, we discuss in chap 18)

Use atmospheric pressure as reference level;

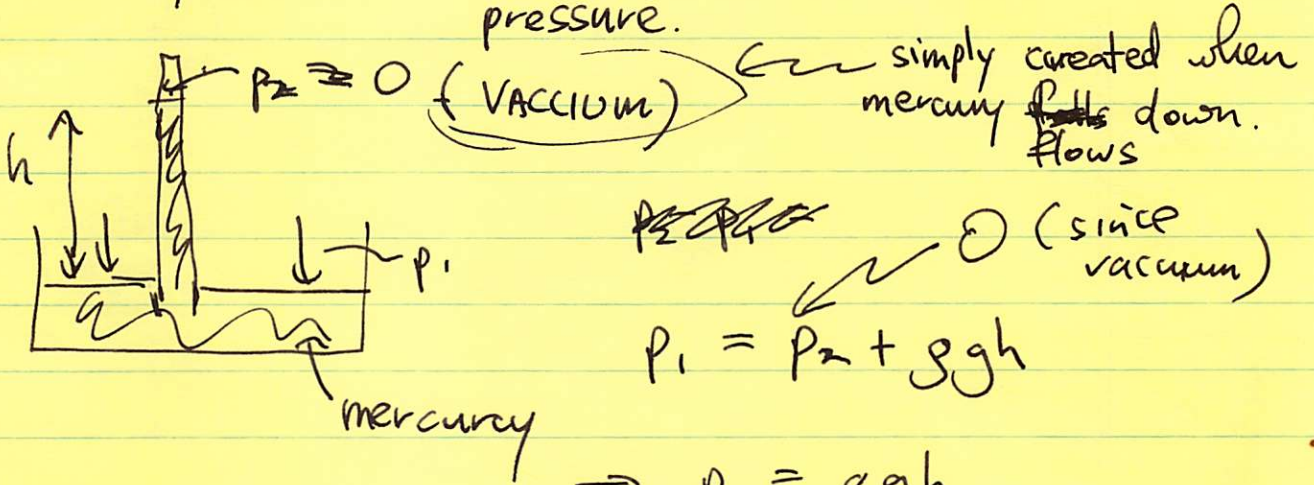
measure gauge pressure  $\Delta p = p_{\text{actual}} - p_{\text{atm}}$

$\leftarrow$  absolute pressure



we use mercury because it is dense, so required height of fluid is less.

Mercury barometer → used to measure atmospheric pressure.



← simply created when mercury falls down. flows

~~$P_2 = P_1$~~  (since vacuum)

$$P_1 = P_2 + \rho gh$$

$$\Rightarrow P_1 = \rho gh.$$

$$\rho_{\text{mercury}} = 1.36 \times 10^4 \text{ kg m}^{-3}$$

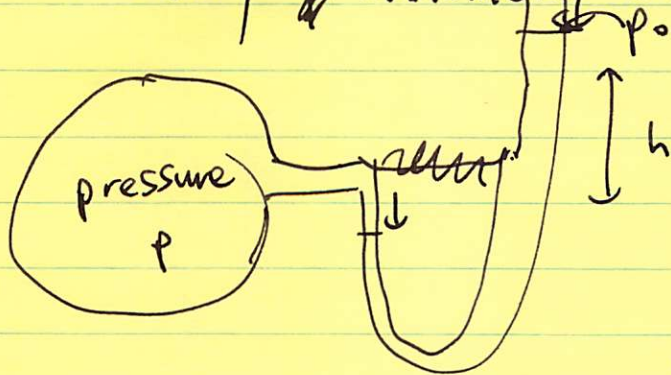
$$h = 760 \text{ mm}$$

$$1 \text{ bar} = 10^5 \text{ Pa}$$

$$P_1 = (1.36 \times 10^4)(9.8)(0.76) \checkmark$$
$$= 1.013 \times 10^5 \text{ N m}^{-2} = 1.013 \text{ bar}$$

This is the weight of column density of air from bottom to top of atmosphere.

Mercury ~~of~~ manometer can be used to measure pressure:



$$P - P_0 = \rho gh$$

(if gas is under low pressure, can use less dense liquid, like water)

Historically = Pascal ~~the~~ got brother-in-law to measure height of mercury on a mountaintop  $\rightarrow$  8cm lower!  
Created a sensation.

## Surface Tension

Leaves & insects float ~~on~~ not bcos of Archimedes' principle, but bcos of surface tension  $\rightarrow$  attractive force between molecules of water.

$\hookrightarrow$  also causes droplets of water to have spherical shape, minimizing  $\frac{\text{Area}}{\text{Vol.}}$

$\hookrightarrow$  can show it by floating a steel needle or razor blade, which should sink by Archimedes principle (indeed, if put it in water, will sink).

$\hookrightarrow$  if put in detergent, which reduces surface tension, will sink!